

MID-TERM EXAM OF STATISTICAL PHYSICS

Wednesday 10 march 2021

Duration : **2hours**

*The use of documents, mobile phones, calculators, . . . , is forbidden.*

Recommendations :

Read the text carefully and **write** out your answers as *succinctly* and as *clearly* as possible.

Check your calculations (dimensional analysis, etc.); do not forget to **reread yourself**.

Check the **informations at the end of the text**

## 1 Frenkel defects (~ 30mn)

Consider a crystalline solid made up of  $N$  atoms. In the ground state, atoms occupy the  $N$  sites of the crystal lattice, however, an atom can leave a site in the lattice and position itself at an interstitial site, which has an energy cost of  $\varepsilon > 0$ . We denote by  $N'$  the number of available interstitial sites (in general  $N$  and  $N'$  are of the same order of magnitude). If  $n$  atoms occupy interstitial sites (so  $N - n$  atoms remain on the  $N$  sites of the crystal lattice), the energy of the atoms is  $E = n\varepsilon$ .

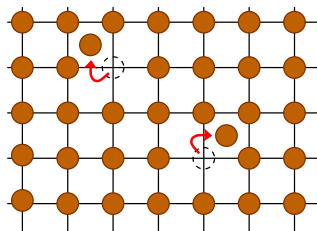


FIGURE 1 : *Crystal of  $N = 35$  atoms with two Frenkel defects.*

- 1/ State the fundamental postulate of statistical physics.
- 2/ Recall the definitions of the microcanonical entropy  $S^*$  and the microcanonical temperature  $T^*$ .
- 3/ How many ways are there to pick the  $n$  atoms leaving the  $N$  sites on the lattice? How many ways are there for these  $n$  atoms to occupy the  $N'$  interstitial sites ?
- 4/ Deduce from the previous question the number  $\Omega$  of microstates corresponding to  $n$  atoms on the interstitial sites. Show that for  $n \ll N, N'$  we have

$$\Omega \simeq \frac{(N'N)^n}{(n!)^2} \quad (1)$$

- 5/ Deduce the microcanonical entropy  $S^*$  atoms of the crystal and give its expression in the limit  $N, N' \gg n \gg 1$ .
- 6/ Deduce the expression of the temperature  $T^*(E)$ . Invert this function to determine how the number  $n$  of atoms on the interstitial sites depends on the temperature  $T^*$ , of  $N$  and of  $N'$ . Does the hypothesis  $n \ll N, N'$  of the question 4/ correspond to high or low temperatures? Qualitatively plot  $n$  as a function of  $T^*$  in this regime and comment.

## 2 Conjugate variables in the microcanonical ensemble (~ 40mn)

Consider a system whose energy depends on a parameter  $\phi$ , a « force », conjugate to an observable  $X$  (for example the magnetic field  $\phi \rightarrow \mathcal{B}$  and the magnetization  $X \rightarrow M$ ). Said differently, the value of the observable in a microstate  $|\ell\rangle$  is related to its energy by

$$X_\ell = -\partial E_\ell / \partial \phi. \quad (2)$$

If the system is isolated, we denote by  $\overline{X}^*$  the microcanonical average of the observable (i.e. the average over the accessible states  $\in [E, E + \delta E]$ ). Let us note  $S^*$  and  $T^*$  the microcanonical entropy and temperature. The objective of the exercise is to show the relation

$$\overline{X}^* = T^* \frac{\partial S^*}{\partial \phi} \quad (3)$$

and to discuss a simple application of this very general formula.

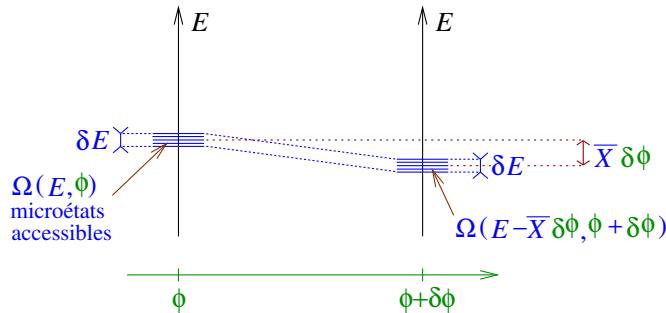


FIGURE 2 : Evolution of the accessible microstates when changing the parameter  $\phi \rightarrow \phi + \delta\phi$ .

We note  $\Omega(E, \phi)$  the number of accessible microstates.

- 1/ If we consider an *adiabatic* transformation driven by a "small" variation of the force,  $\phi \rightarrow \phi + \delta\phi$ , the energy of the system changes like  $E \rightarrow E - \overline{X}^* \delta\phi$ , according to (2). Under such transformation, the number of accessible microstates is conserved (Fig. 2). Show the relation (3).
- 2/ **Application : crystal of 1/2 spins**

We apply these considerations in a very simple case : we consider an (isolated) crystal of  $N$  spins 1/2, under a magnetic field  $\mathcal{B}$ . We note  $n_\pm$  the number of spins in the state  $|\pm\rangle$  of energy  $\varepsilon_\pm = \mp \varepsilon_{\mathcal{B}}$  with  $\varepsilon_{\mathcal{B}} = m_0 \mathcal{B}$  where  $m_0$  is the magnetization of one spin.

- a) Give the expression of the number of accessible microstates  $\Omega$  as a function of  $N$ ,  $n_+$  and  $n_-$ .
- b) Deduce the microcanonical entropy  $S^*$  (assuming  $N$ ,  $n_\pm \gg 1$ ). Justify that it can be written in the form  $S^*(E, N, \mathcal{B}) = N k_B s(E/N \varepsilon_{\mathcal{B}})$  and give the expression of the dimensionless function  $s(x)$ .
- c) Apply the formula (3) for the magnetization (i.e.  $X \rightarrow M$  and  $\phi \rightarrow \mathcal{B}$ ) and interpret the result.
- d) *High temperature limit.*— Show that  $s(x) \simeq \ln 2 - x^2/2$  for  $x \ll 1$ . Deduce an approximate expression of the magnetic temperature  $T^*$ . Give  $\overline{M}^*$  as a function of  $N$ ,  $m_0$ ,  $\mathcal{B}$  and  $T^*$ . Interpret the result.
- e) **Bonus (optional)** : In the more general case of spins  $s > 1/2$ , which part of the analysis would change and which results would be unchanged ?

### 3 Anharmonic oscillator ( $\sim 45\text{mn}$ )

We consider a system in contact with a thermostat at temperature  $T$ .

- 1/ Give the definition of the canonical partition function  $Z$  (we note  $|\ell\rangle$  the microstates and  $E_\ell$  the energies).
- 2/ How do we deduce the average energy  $\overline{E}^c$  of the partition function (recall the demonstration).
- 3/ We consider a one-dimensional system described by its Hamiltonian  $H(x, p) = \frac{p^2}{2m} + V(x)$ . The system is treated classically.
  - a) Show that the partition function factorizes in the form  $Z = Z_{\text{cin}} Z_{\text{pot}}$  where  $Z_{\text{cin}}$  and  $Z_{\text{pot}}$  are respectively associated with the kinetic and potential parts of the energy.
  - b) Calculate  $Z_{\text{cin}}$  and deduce the average kinetic energy  $\overline{E}_{\text{cin}}^c$ .
  - c) We will analyse  $Z_{\text{pot}}$  using an approximate potential  $V(x) = \frac{1}{2}\kappa x^2 + \lambda x^4$ . In that respect, we write  $e^{-\beta V(x)} = e^{-\frac{1}{2}(x/a_2)^2} e^{-(x/a_4)^4}$ . Give the expressions of the two length scales  $a_2$  and  $a_4$  as a function of  $T$ .
  - d) By plotting the shapes of  $e^{-\frac{1}{2}(x/a_2)^2}$  and  $e^{-(x/a_4)^4}$ , justify that the quartic term is negligible within the limit of low temperatures. Deduce  $Z_{\text{pot}}$  and the contribution of the potential energy  $\overline{E}_{\text{pot}}^c$  to the total average energy.
  - e) In the high temperature limit, the quadratic term of  $V(x)$  is negligible. Show that  $Z_{\text{pot}} \propto \beta^{-\theta}$  and specify the value of the exponent  $\theta$ . Infer  $\overline{E}_{\text{pot}}^c$ .
  - f) Calculate the specific heat  $C_V = \partial \overline{E}^c / \partial T$  in the two limits (with  $\overline{E}^c = \overline{E}_{\text{cin}}^c + \overline{E}_{\text{pot}}^c$ ). Identify the temperature scale  $T_{\text{crossover}}$  separating the two regimes, that you will express in terms of  $\kappa$  and  $\lambda$ . Qualitatively plot  $C_V$  and provide a physical interpretation.

### Reread yourself ( $\sim 5\text{mn}$ )

#### Annex

- Stirling formula :  $\ln N! \approx N \ln N - N$  for  $N \gg 1$ .
- $N!/(N-m)! \simeq N^m$  for  $N \gg m$ .
- $\int_{\mathbb{R}} dx e^{-ax^2} = \sqrt{\pi/a}$