LARGE DEVIATIONS IN ACTIVE MATTER

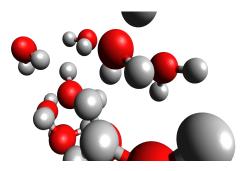
- OVERVIEW AND EXAMPLES -

- OUTLINE -

- Overview
 - What is «Active Matter»?A variety of results
- Example 1: the Active Ornstein–Uhlenbeck Process (AOUP)
 Steady-state and escape rate
 Open questions
- Example 2: an active lattice gas model with non-Gaussian noise
 Fluctuating Hydrodynamics
 Dynamical Phase Transitions

- OVERVIEW - What is «Active Matter»?

What distinguishes the following systems?



molecules in solution



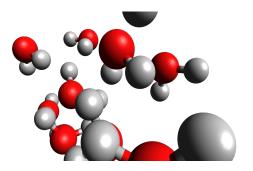
a sandstorm



a murmuration of starlings

- OVERVIEW - What is «<u>Active Matter</u>»?

What distinguishes the following systems?



molecules in solution

Equilibrium



a sandstorm



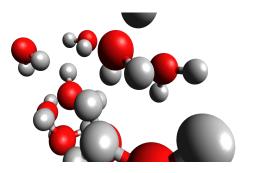
a murmuration of starlings

External drive

<u>(External and) internal drive</u>

- OVERVIEW - What is «<u>Active Matter</u>»?

What distinguishes the following systems?



molecules in solution

a sandstorm



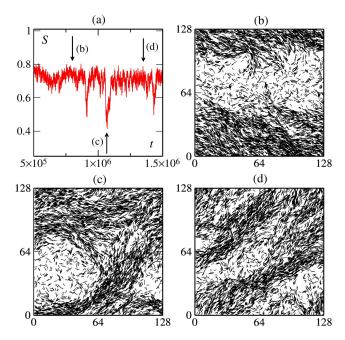
a murmuration of starlings

<u>Recent review</u>: J Tailleur, G Gompper, M C Marchetti, JM Yeomans & C Salomon, 2018 school "Active Matter and Nonequilibrium Statistical Physics", Lecture Notes of the Les Houches Summer School, OUP 2022

Examples and Motivations for studying Large Deviations

· Giant number fluctuations in active fluids

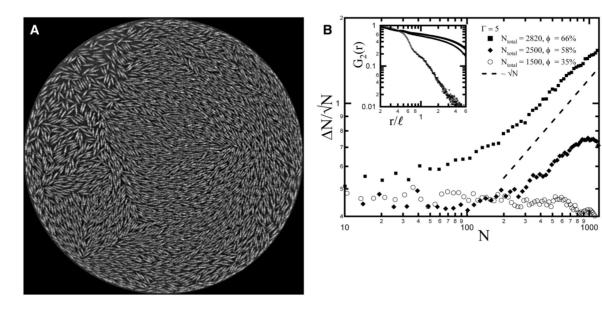
Ramaswamy et al. [EPL 62 196 (2003)], Chaté et al. [PRL 96 180602 (2006)]



EXAMPLES AND MOTIVATIONS FOR STUDYING LARGE DEVIATIONS

· Giant number fluctuations in active fluids

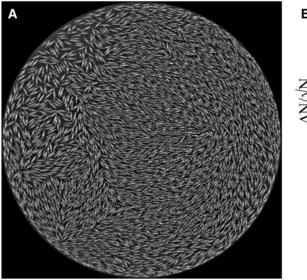
Narayan, Ramaswami & Menon [Science 317 105 (2007)]

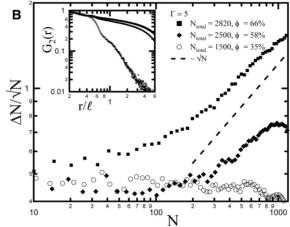


EXAMPLES AND MOTIVATIONS FOR STUDYING LARGE DEVIATIONS

· Giant number fluctuations in active fluids

Narayan, Ramaswami & Menon [Science 317 105 (2007)]





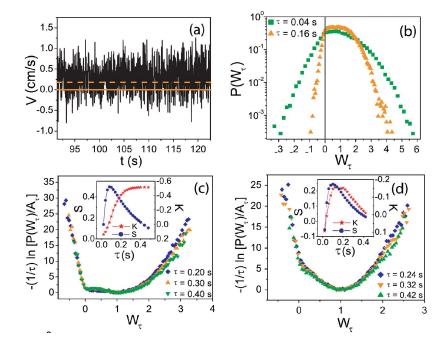
 $\Delta N \equiv [\langle N^2 \rangle - \langle N \rangle^2]^{\frac{1}{2}}$ std dev of #particles

 $\cdot \Delta N \propto N^{\frac{1}{2}} \ll @ eq. \gg$

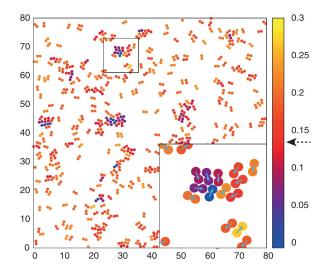
 $\cdot \Delta N \propto N^{\alpha > 1/2}$ active \rightarrow anomalous cumulant scaling

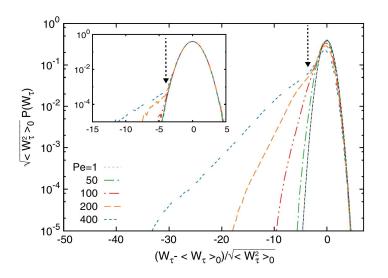
→ anomalous scaling for the cumulant generating function (CGF)?

• <u>An experiment</u>: Distribution of time-averaged velocity for a gas polar granular rods confined in 2D Kumar et al. [PRL **106** 118001 (2011)]



• <u>An example model</u>: Distribution of "active work" for self-propelled dumbbell particles Cagnetta et al. [PRL **119** 158002 (2017)]





• <u>An example model</u>: Active Brownian Particles (APB) $\dot{\mathbf{x}} = -\mu \nabla V + v \mathbf{u}(\theta) + \sqrt{2D} \boldsymbol{\xi}(t)$ Weak-noise description ($D \rightarrow 0$), for a single particle:

$$P(\mathbf{x}_2, t | \mathbf{x}_1, 0) = \int_{\mathbf{x}_1}^{\mathbf{x}_2} \mathcal{D}[\mathbf{x}(t), \theta(t)] e^{-(1/D)\mathcal{A}[\mathbf{x}, \theta]} \mathcal{P}[\theta(t)]$$

$$\mathcal{A}[\mathbf{x},\theta] = \frac{1}{4} \int_0^t \|\dot{\mathbf{x}} + \mu \nabla V(\mathbf{x}) - v\mathbf{u}(\theta)\|^2 dt'$$
$$\mathcal{P}[\theta(t)] \propto e^{-\int_0^t \dot{\theta}^2/(4\alpha) dt'}$$

• <u>An example model</u>: Active Brownian Particles (APB) $\dot{\mathbf{x}} = -\mu \nabla V + v \mathbf{u}(\theta) + \sqrt{2D} \boldsymbol{\xi}(t)$ Weak-noise description @ $D \rightarrow 0$

$$P(\mathbf{x}_2, t | \mathbf{x}_1, 0) = \int_{\mathbf{x}_1}^{\mathbf{x}_2} \mathcal{D}[\mathbf{x}(t), \theta(t)] e^{-(1/D)\mathcal{A}[\mathbf{x}, \theta]} \mathcal{P}[\theta(t)]$$

Minimizing w.r.t. $\theta(t)$ [contraction principle]:

$$\int \mathcal{D}[\theta(t)] e^{-(1/D)\mathcal{A}[\mathbf{x},\theta]} \mathcal{P}[\theta(t)] \underset{D \to 0}{\asymp} e^{-(1/D)\mathcal{A}[\mathbf{x},\tilde{\theta}]} \qquad \mathcal{A}[\mathbf{x},\tilde{\theta}] = \inf_{\theta} \left(\frac{1}{4} \int_{0}^{t} \|\dot{\mathbf{x}} + \mu \nabla V(\mathbf{x}) - v\mathbf{u}(\theta)\|^{2} dt'\right)$$

$$Optimal \text{ angle trajectory:} \quad \mathbf{u}(\tilde{\theta}) = \frac{\dot{\mathbf{x}} + \mu \nabla V(\mathbf{x})}{\|\dot{\mathbf{x}} + \mu \nabla V(\mathbf{x})\|}$$

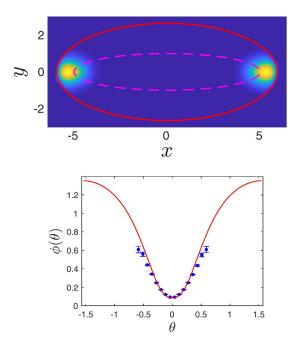
· <u>An example model</u>: Active Brownian Particles (APB) $\dot{\mathbf{x}} = -\mu \nabla V + v \mathbf{u}(\theta) + \sqrt{2D} \boldsymbol{\xi}(t)$ Substituting, we arrive at

$$\mathcal{A}[\mathbf{x}] = \frac{1}{4} \int_{-\infty}^{\infty} (\|\dot{\mathbf{x}} + \mu \nabla V(\mathbf{x})\| - v)^2 dt'$$

Escape problem: mean escape time $\langle T \rangle$ between from domain C_1 to domain C_2 given by

$$\begin{aligned} \langle \tau \rangle & \underset{D \to 0}{\asymp} e^{\frac{\phi}{D}}, \\ \phi &= \inf_{\{\mathbf{x}_1 \in \mathbf{C}_1, \mathbf{x}_2 \in \mathbf{C}_2\}} \inf_{\mathbf{x}(t)} \mathcal{A}[\mathbf{x}(t)] \end{aligned}$$

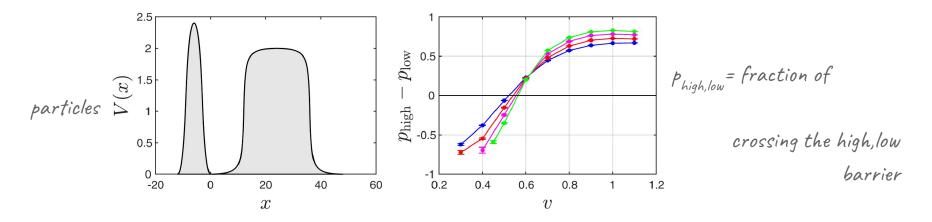
• <u>An example model</u>: Active Brownian Particles (APB) $\dot{\mathbf{x}} = -\mu \nabla V + v \mathbf{u}(\theta) + \sqrt{2D} \boldsymbol{\xi}(t)$ Escape from an elliptic trap in 2D



rate of escape from an elliptic trap $V = \frac{1}{2} \times A \times (\text{pink to purple})$ the escape time would be uniform for a passive particle (@D \rightarrow 0)

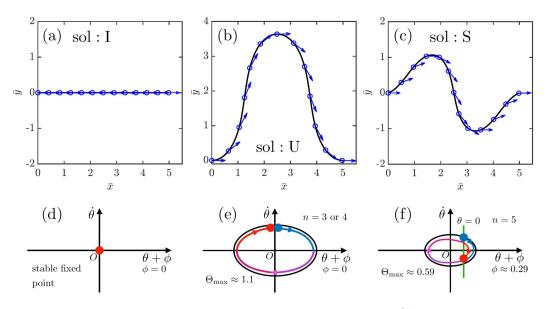
comparison between analytical prediction and numerics

• <u>An example model</u>: Active Brownian Particles (APB) $\dot{\mathbf{x}} = -\mu \nabla V + v \mathbf{u}(\theta) + \sqrt{2D} \boldsymbol{\xi}(t)$ From barrier of potential to barrier of force



At large activity v, the particle prefers to cross the barrier with the smaller maximal force.

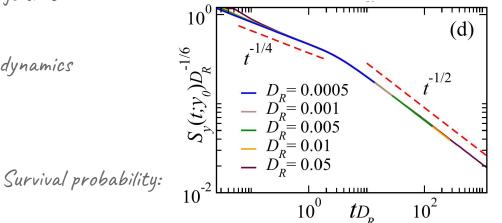
• <u>An example model</u>: Active Brownian Particles (APB) $\dot{\mathbf{x}} = -\mu \nabla V + v \mathbf{u}(\theta) + \sqrt{2D} \boldsymbol{\xi}(t)$ More generic optimal paths in 2D: depending on boundary conditions



[Yasuda & Ishimoto, PRE **106** 064120 (2022)]

• <u>An example model</u>: Active Brownian Particles (APB) $\dot{\mathbf{x}} = -\mu \nabla V + v \mathbf{u}(\theta) + \sqrt{2D} \boldsymbol{\xi}(t)$ 2D case without potential:

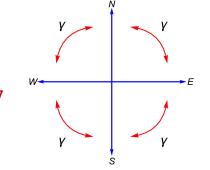
- Anomalous scaling of the marginal distributions at short times
 Anomalous scaling of 1st-passage time at short times
- Fingerprint of the short-time dynamics at large times



[Majumdar & Meerson, PRE 102 022113 (2020)] [Basu, Majumdar, Rosso & Schehr, PRE 98 062121 (2018)]

• <u>An example model</u>: Active Brownian Particles (APB) $\dot{\mathbf{x}} = -\mu \nabla V + v \mathbf{u}(\theta) + \sqrt{2D} \boldsymbol{\xi}(t)$ 2D case with a harmonic potential: Exact solution at all times [Caraglio & Franosch, PRL 129 158001 (2022)]

• <u>Another important model</u>: <u>Run-and-Tumble (RT)</u> 2D case with a harmonic potential: Exact solution at all times <u>[Smith, Le Doussal, Majumdar & Schehr</u> PRE **106** 054133 (2022)]

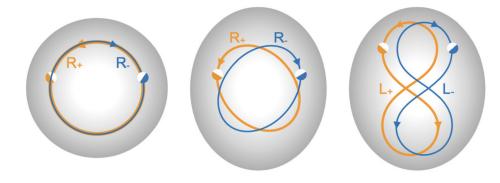


· How to extract the large-time and/or small-noise asymptotics?

Do we learn generic features of active particles from this?

• <u>An example model</u>: Active Brownian Particles (APB) $\dot{\mathbf{x}} = -\mu \nabla V + v \mathbf{u}(\theta) + \sqrt{2D} \boldsymbol{\xi}(t)$ 2D case with a torque:

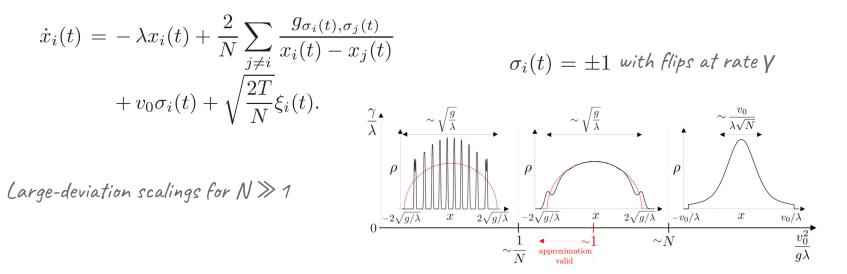
- radially symmetric or elliptic confinement
- different types of optimal trajectories



[Damascena & de Souza Silva, PRE **108** 044605 (2023)]

• <u>An example model</u>: Active Brownian Dyson model

N particles in 1D, non-intersecting, with a logarithmic repulsion



[Touzo, Le Doussal & Schehr, EPL **142** 61004(2023); PRE **109** 014136 (2024)]

· <u>A variety of large-deviation scalings</u>: a run-and-tumble model in a parabolic potential

$$m\ddot{x}(t) + \Gamma\dot{x}(t) + kx(t) = \Sigma(t)$$

 \cdot $\Sigma(t)$ is a telegraphic noise of typical correlation time \mathcal{T}_c

· Large-deviation scaling for the steady-state as $\mathbf{T}_{c} \rightarrow 0$

$$P_{\rm st}(X) \sim e^{-s(X)/\tau_c} \qquad \tau_c = \sqrt{k/m} \, \mathcal{T}_c$$

[Smith & Farago, PRE 106 054118 (2022)]

LARGE DEVIATIONS IN ACTIVE MATTER - EXAMPLE METHODOLOGY

• <u>"Effective" / "auxiliary" forces</u> as a tool

From Hugo Touchette's lecture: if A[trajectory] is extensive in the duration t of the trajectory

 $P[A/t = a] \approx \exp[t I(a)] \iff \langle e^{sA} \rangle \approx \exp[t \Psi(s)]$ "biased" ensemble $(\Psi(s) = \text{scaled CGF})$

modified dynamics with s-dependent <u>"effective forces</u>"

<u>Methodology</u>: if the biased ensemble presents interesting features (clustering, oscillations, etc.), use the effective forces to design protocols to control the system

[Keta et al., PRE **103** 022603 (2021)] [Fodor et al., Ann Rev CondMat Phys **13** 215 (2022)] [Lamtyugina et al., PRL **129** 128002 (2022)] [Piñeros & Fodor, arXiv:2403.16961 (2024)]

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 - Open questions
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THE AOUP - **REFERENCES**

A. J. Bray and A. J. McKane, Instanton Calculation of the Escape Rate for Activation over a Potential Barrier Driven by Colored Noise, Phys. Rev. Lett. **62**(5), 493 (1989), doi:10.1103/PhysRevLett.62.493.

A. J. McKane, *Noise-induced escape rate over a potential barrier: Results for general noise*, Phys. Rev. A **40**(7), 4050 (1989), doi:10.1103/PhysRevA.40.4050.

H. C. Luckock and A. J. McKane, *Path integrals and non-Markov processes*. *III. Calculation of the escape-rate prefactor in the weak-noise limit*, Phys. Rev. A **42**(4), 1982 (1990), doi:10.1103/PhysRevA.42.1982.

M. M. Kłosek-Dygas, B. J. Matkowsky and Z. Schuss, *Colored Noise in Dynamical Systems*, SIAM J. Appl. Math. **48**(2), 425 (1988), doi:10.1137/0148023.

J. F. Luciani and A. D. Verga, *Functional Integral Approach to Bistability in the Presence of Correlated Noise*, EPL **4**(3), 255 (1987), doi:10.1209/0295-5075/4/3/001.

H. S. Wio, P. Colet, M. San Miguel, L. Pesquera and M. A. Rodríguez, *Path-integral formulation for stochastic processes driven by colored noise*, Phys. Rev. A **40**(12), 7312 (1989), doi:10.1103/PhysRevA.40.7312.

Blackboard lecture based on [E Woillez, Y Kafri & VL,] Stat Mech 063204 (2020)]