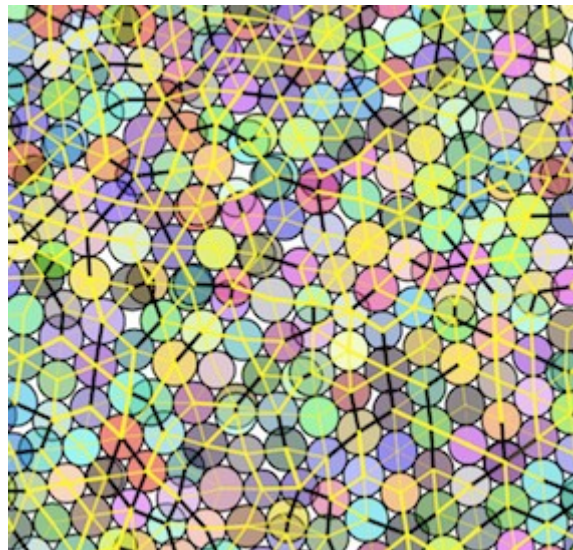


Introduction to the statistical physics of phase transitions and critical phenomena



Emmanuel Trizac
LPTMS / University Paris-Saclay

Disorder in Complex Systems @ IPa

Un Pascal → not only a unit of pressure!

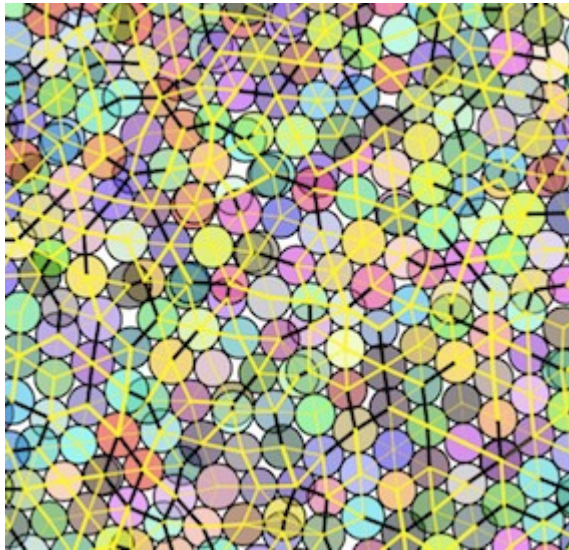


Blaise Pascal



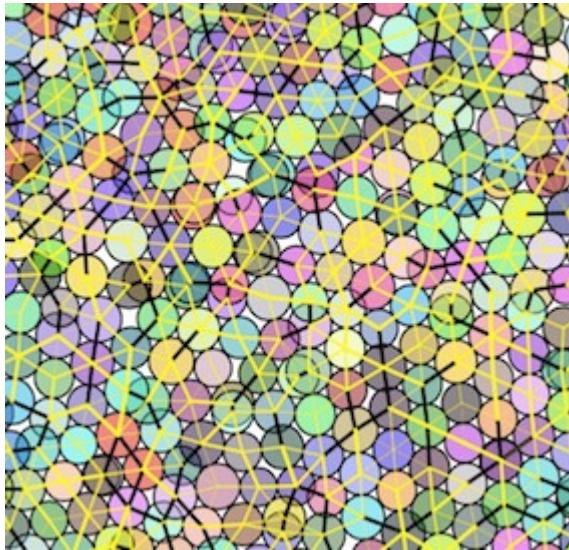
Port-Royal des Champs
15 km from IPa
8 km from St Rémy
(→ *chemin Racine*)

You registered for a school on disorder...



You registered for a school on disorder...

DISCLAIMER



→ course deals with emergence of *order*
through cooperative phenomena

J'écrirai ici mes pensées sans ordre [...]

Je ferais trop d'honneur à mon sujet, si je le traitais avec ordre

Blaise Pascal, Pensées

Statistical physics of phase transitions and critical phenomena

I Introduction

→ Classification, universality, effect of dimension, broken symmetry

II First order phase transitions

→ Unstable isotherms, double-tangent and Maxwell construction, spinodal and binodal

III Critical phenomena : qualitative approaches → mean-field

→ Variational mean-field

→ Statistical field theory / Ginzburg-Landau approach

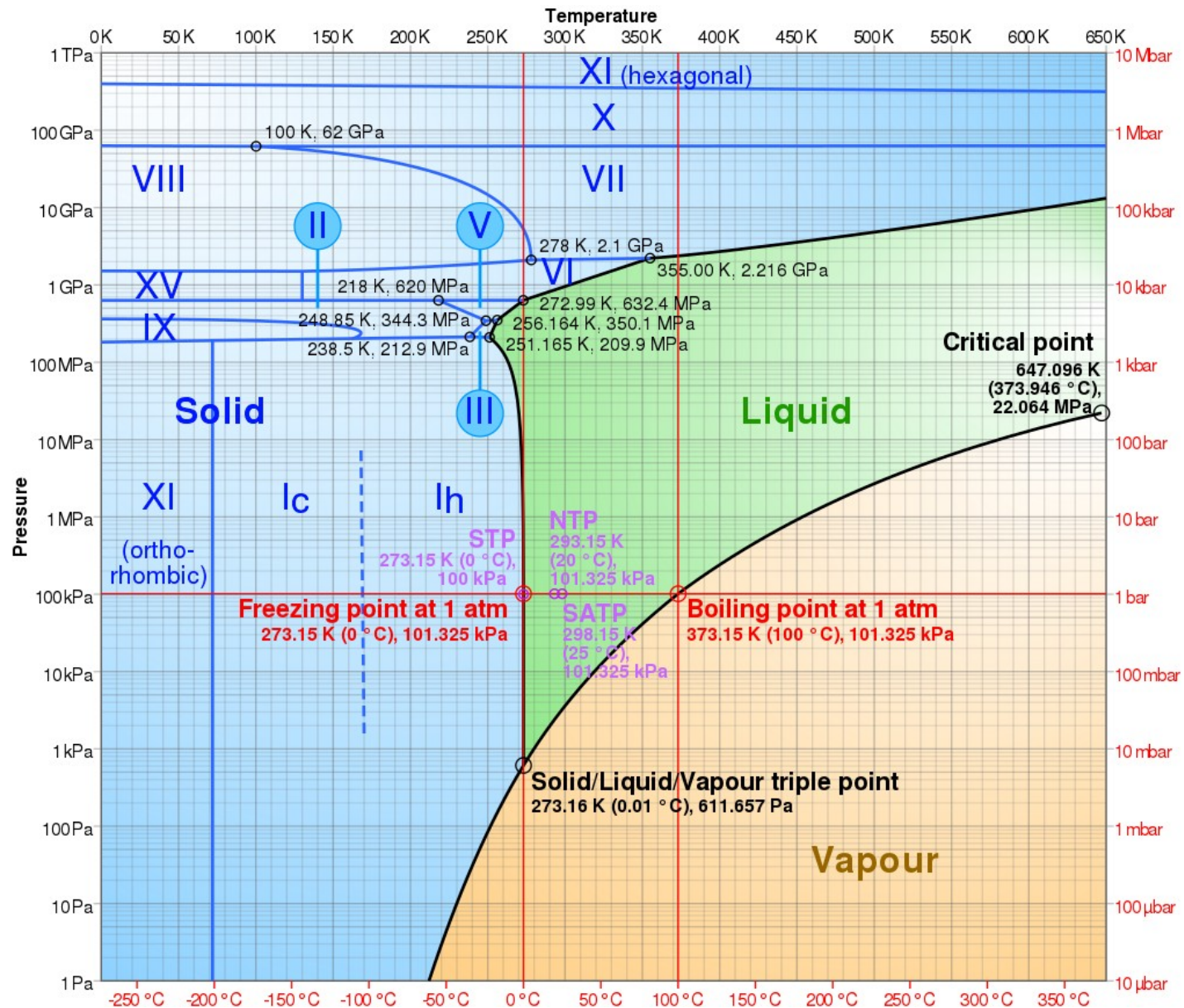
IV Beyond mean-field: fluctuations and scaling

→ Upper critical dimension

V Renormalization group ideas

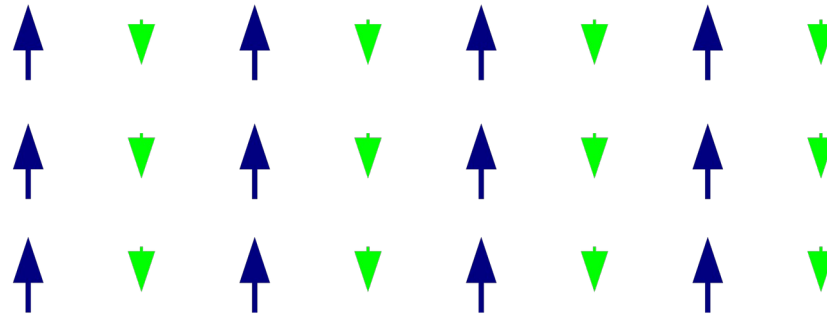
→ Take advantage from the large value of the correlation length ξ

Phase diagram of water



Ce-Sb (Cerium Antimonide)

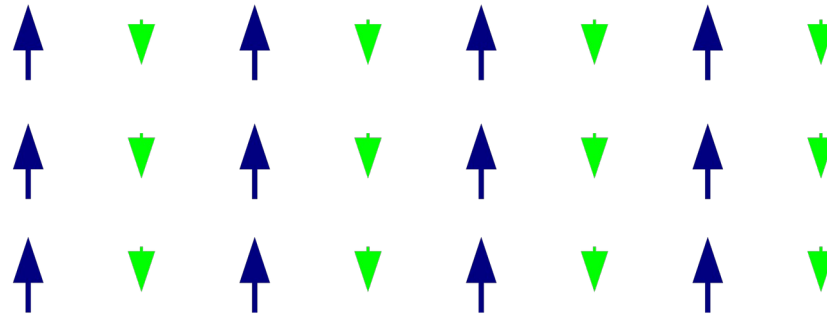
- 14 different phases (ferrimagnetic), revealed by neutron scattering



- Why the strange name *antimoine/antimonium* (against the monk)?
 - of pigs and monks?
 - unlikely: antimony used as a pill until XIXth century
eternal pill...

Ce-Sb (Cerium Antimonide)

- 14 different phases (ferrimagnetic), revealed by neutron scattering



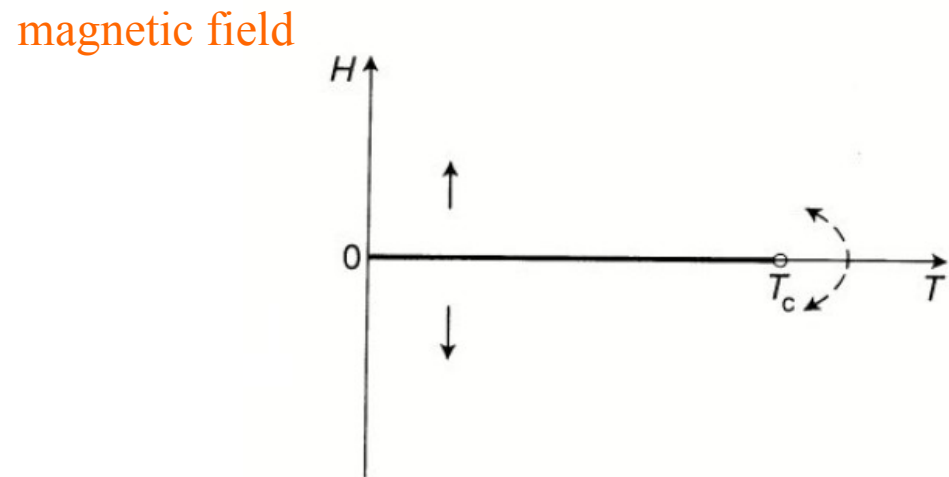
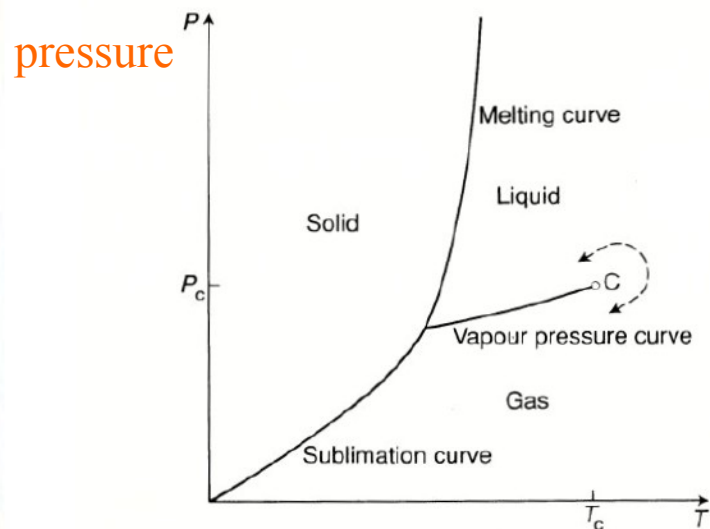
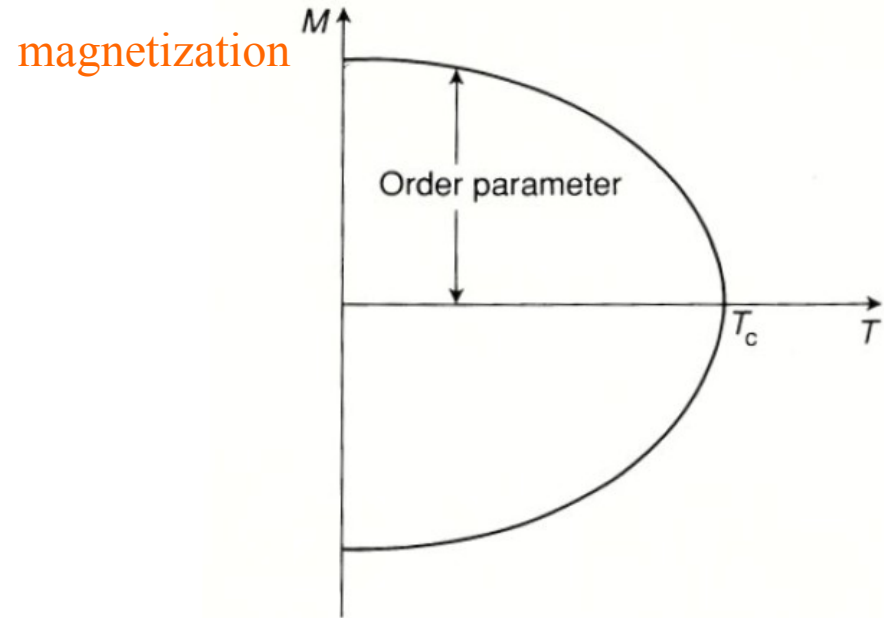
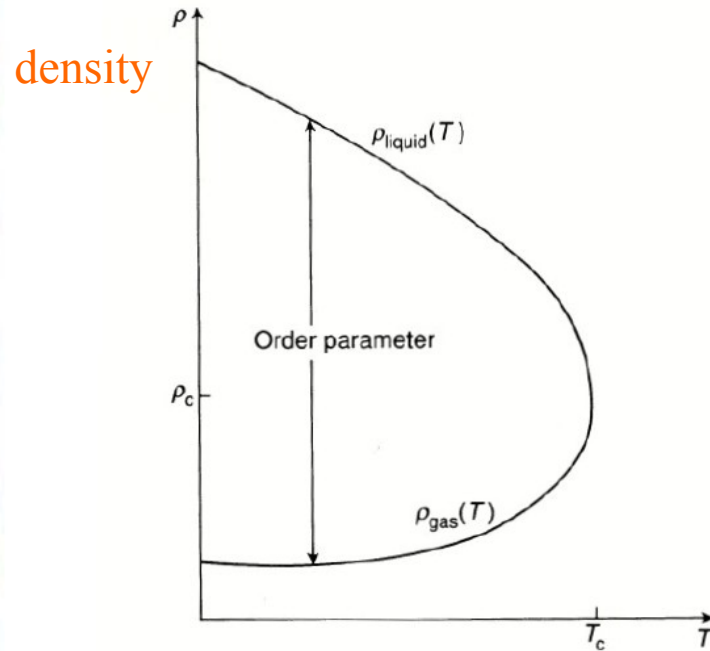
- Why the strange name *antimoine/antimonium* (against the monk)?
 - of pigs and monks?
 - unlikely: antimony used as a pill until XIXth century
eternal pill...
although laxative!

Complexity / variety of phase changes

Transition	Example	Order parameter
ferromagnetic ^a	Fe	magnetization
antiferromagnetic ^a	MnO	sublattice magnetization
ferrimagnetic ^a	Fe ₃ O ₄	sublattice magnetization
structural ^b	SrTiO ₃	atomic displacements
ferroelectric ^b	BaTiO ₃	electric polarization
order-disorder ^c	CuZn	sublattice atomic concentration
phase separation ^d	CCl ₄ +C ₇ F ₁₆	concentration difference
superfluid ^e	liquid ⁴ He	condensate wavefunction
superconducting ^f	Al, Nb ₃ Sn	ground state wavefunction
liquid crystalline ^g	rod molecules	various

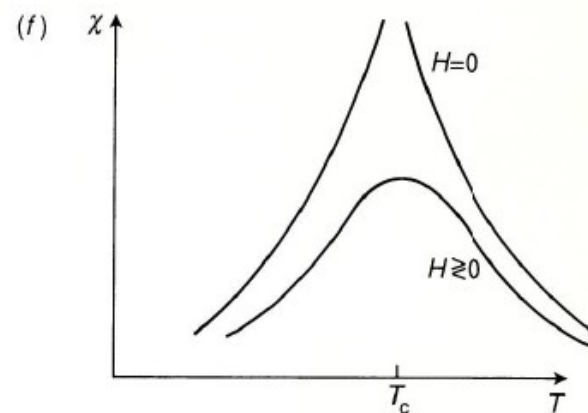
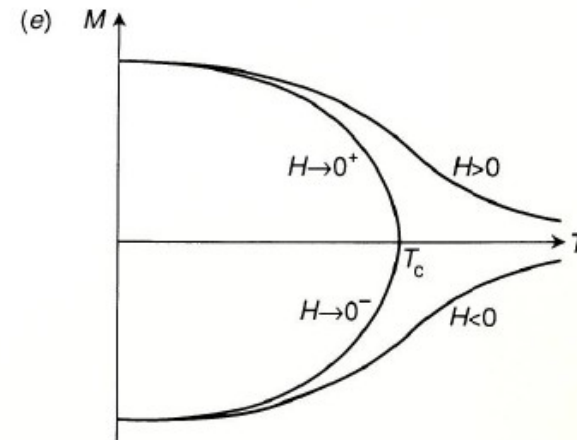
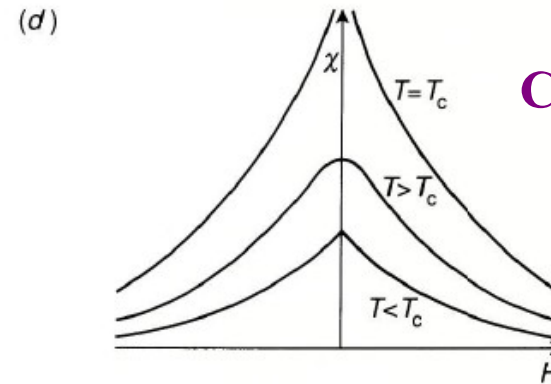
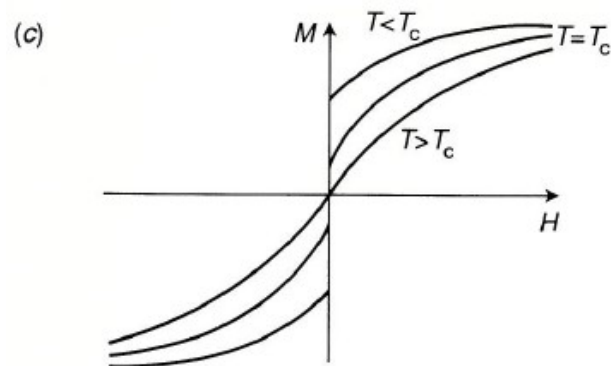
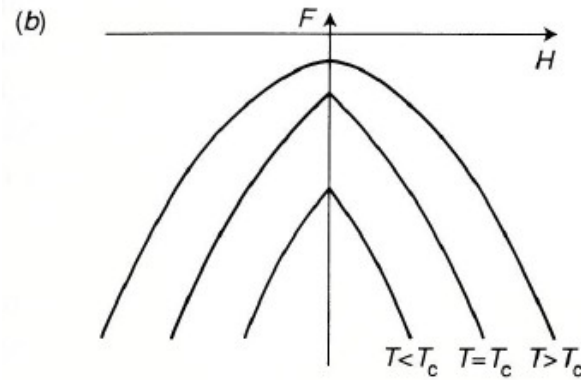
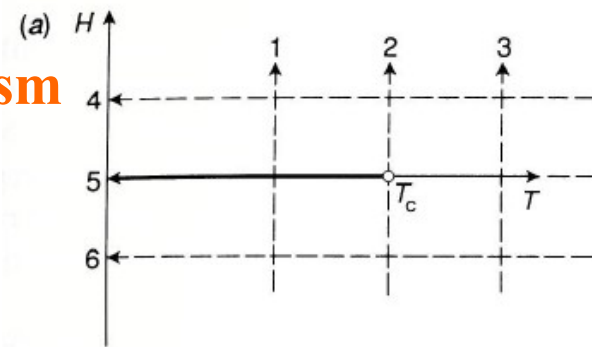
From Yeomans, *Statistical Mechanics of Phase Transitions* (Oxford)

Some systems are equivalent Magnets (\rightleftharpoons liquids) as prototypical examples



MAGNETS

Ferromagnetism



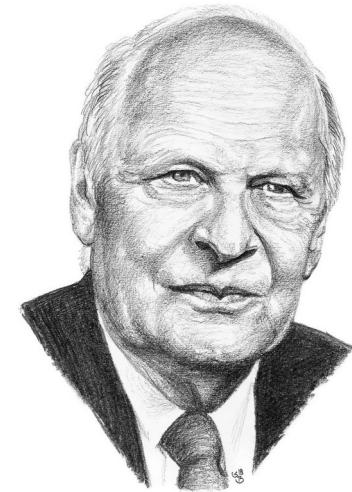
Correlation length diverges

- Abundance/complexity
- Identify an order parameter \leftrightarrow free energy derivative

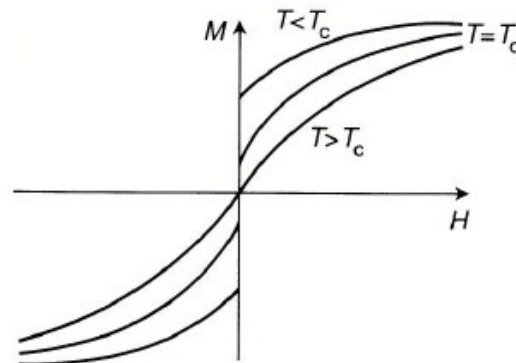
$$M = - \left(\frac{\partial F}{\partial B} \right)_T$$

- Onsager's *singular* legacy

- Singularity absent in a finite system...
→ consider the **thermodynamic limit**



Lars Onsager
(1903-1976)



C h a p t e r

1

Introduction to phase transitions and critical phenomena

- 1- Problems raised by phase transitions, from a stat mech perspective
- 2- **Classification of phase transitions**
- 3- The drosophila of phase transitions
- 4- Order parameter and symmetry breakdown
- 5- Local order and correlation functions : from magnets to liquids

Characterization of 2nd order transitions → critical exponents

fluids

Specific heat at constant volume V_c $C_V \sim |t|^{-\alpha}$

Liquid–gas density difference $(\rho_l - \rho_g) \sim (-t)^\beta$

Isothermal compressibility $\kappa_T \sim |t|^{-\gamma}$

Critical isotherm ($t = 0$)
 $P - P_c \sim |\rho - \rho_c|^\delta \text{sgn}(\rho - \rho_c)$

Correlation length $\xi \sim |t|^{-\nu}$

Pair correlation function at T_c $G(\vec{r}) \sim 1/r^{d-2+\eta}$

magnets

Zero-field specific heat $C_H \sim |t|^{-\alpha}$

Zero-field magnetization $M \sim (-t)^\beta$

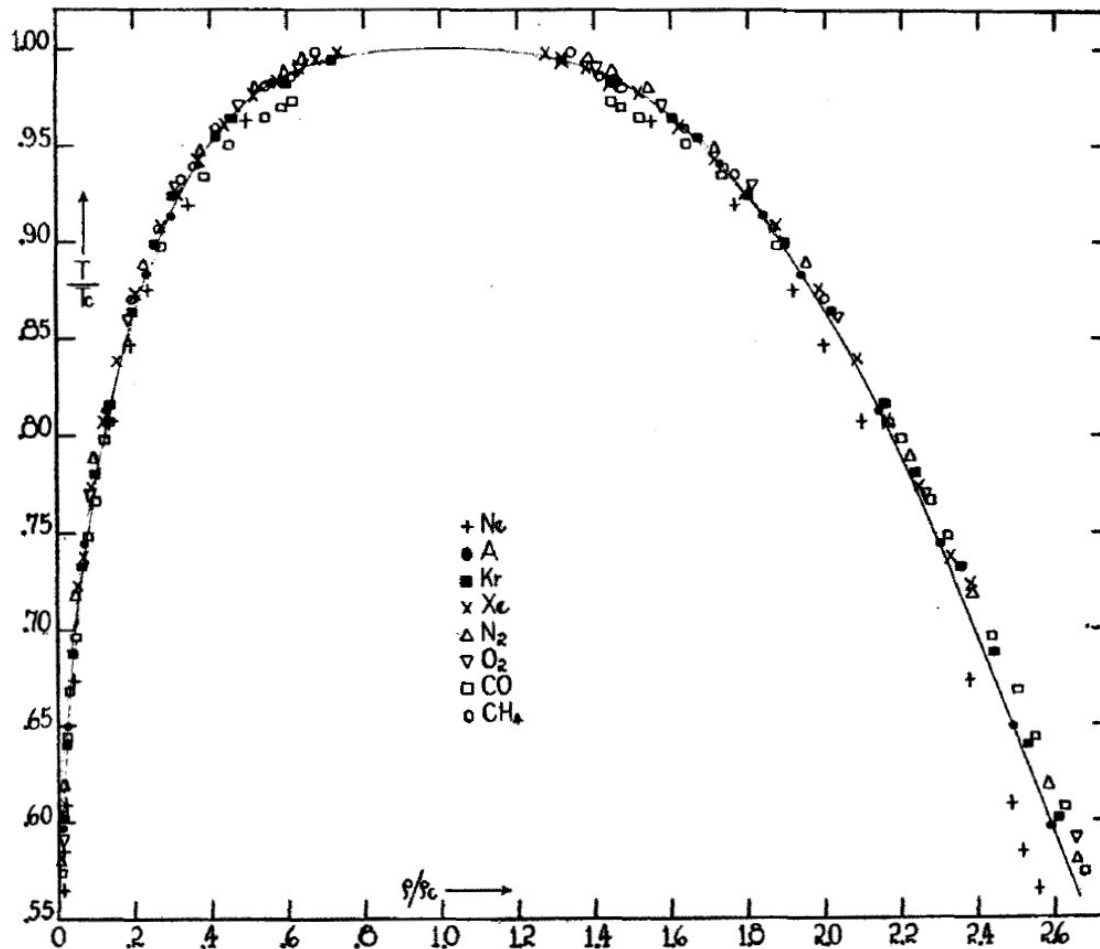
Zero-field isothermal susceptibility $\chi_T \sim |t|^{-\gamma}$

Critical isotherm ($t = 0$)
 $H \sim |M|^\delta \text{sgn}(M)$

Correlation length $\xi \sim |t|^{-\nu}$

Pair correlation function at T_c $G(\vec{r}) \sim 1/r^{d-2+\eta}$

Why are critical exponents interesting ? → several layers of universality



E. A. Guggenheim,
J. Chem. Phys. 13, 253 (1945)

	T_c (K)	P_c (atm)
Ne	45	26
Ar	150	48
Kr	209	54
Xe	290	58
N ₂	126	33
O ₂	154	50
CO	133	34
CH ₄	190	45

Corresponding states for liquid-gas transition

There is more to universality

- Liquid-gas transition: $\beta \sim 0.33$
- Magnets with uniaxial anisotropy (MnF_2): $\beta \sim 0.33$
- Phase separation in binary mixture ($\text{CCl}_4 + \text{C}_7\text{F}_{16}$): $\beta \sim 0.33$
- 3d Ising model on cubic lattice, fcc etc...: $\beta \sim 0.33$
- ...

→ all belong to the same universality class

What matters is space dimension + symmetry of order parameter
e.g. for Ising in 2D : $\beta = 1/8$

Different universality classes

Universality class	Symmetry of order parameter	α	β	γ	δ	ν	η	Physical examples
2-d Ising	2-component scalar	0 (log)	1/8	7/4	15	1	1/4	some adsorbed mono e.g. H on Fe
3-d Ising	2-component scalar	0.10	0.33	1.24	4.8	0.63	0.04	phase separation, flu order-disorder e.g. β
3-d X-Y	2-dimensional vector	0.01	0.34	1.30	4.8	0.66	0.04	superfluids, supercon
3-d Heisenberg	3-dimensional vector	-0.12	0.36	1.39	4.8	0.71	0.04	isotropic magnets
mean-field		0 (dis.)	1/2	1	3	1/2	0	
2-d Potts, $q=3$ $q=4$	q -component scalar	1/3 2/3	1/9 1/12	13/9 7/6	14 15	5/6 2/3	4/15 1/4	some adsorbed mono e.g. Kr on graphite

→ define the $O(n)$ model

→ existence of scaling relations, eg $\nu d = 2 - \alpha$

Chapter

1

Introduction to phase transitions and critical phenomena

- 1- Problems raised by phase transitions, from a stat mech perspective
- 2- Classification of phase transitions
- 3- Ising model : the drosophila of phase transitions
- 4- Order parameter and symmetry breakdown
- 5- Local order and correlation functions : from magnets to liquids



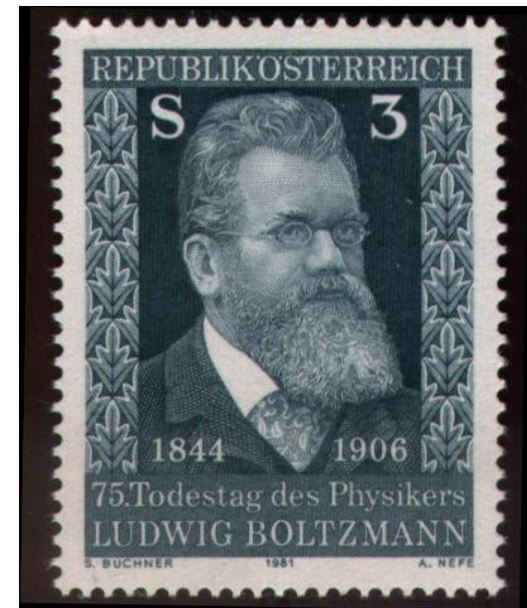
Ernst Ising (1900-1998)

Why a transition ?

energy \leftrightarrow entropy competition...



against





Why a transition ?

Often energy \leftrightarrow entropy competition...

but not always

Non spherical colloids...

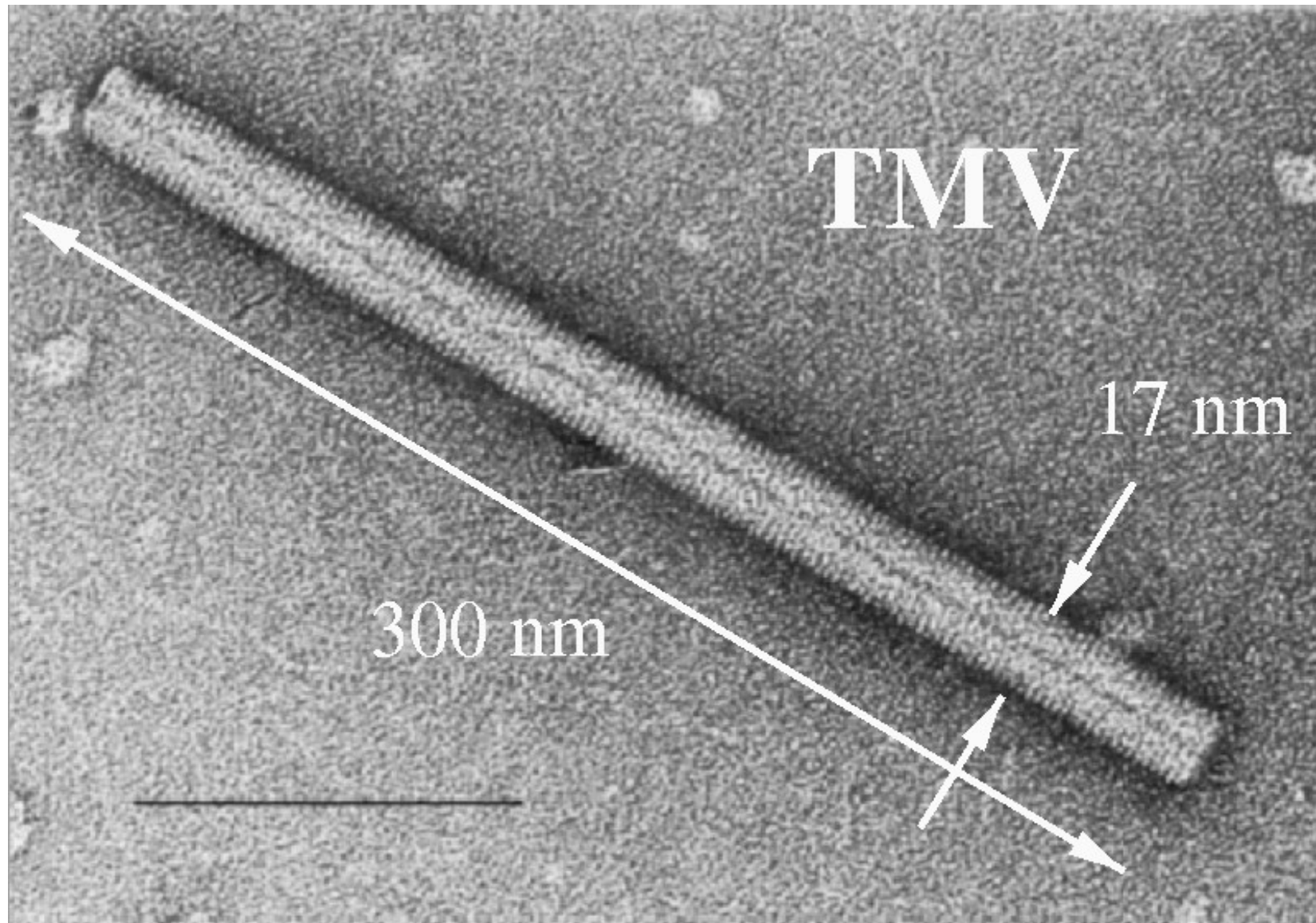
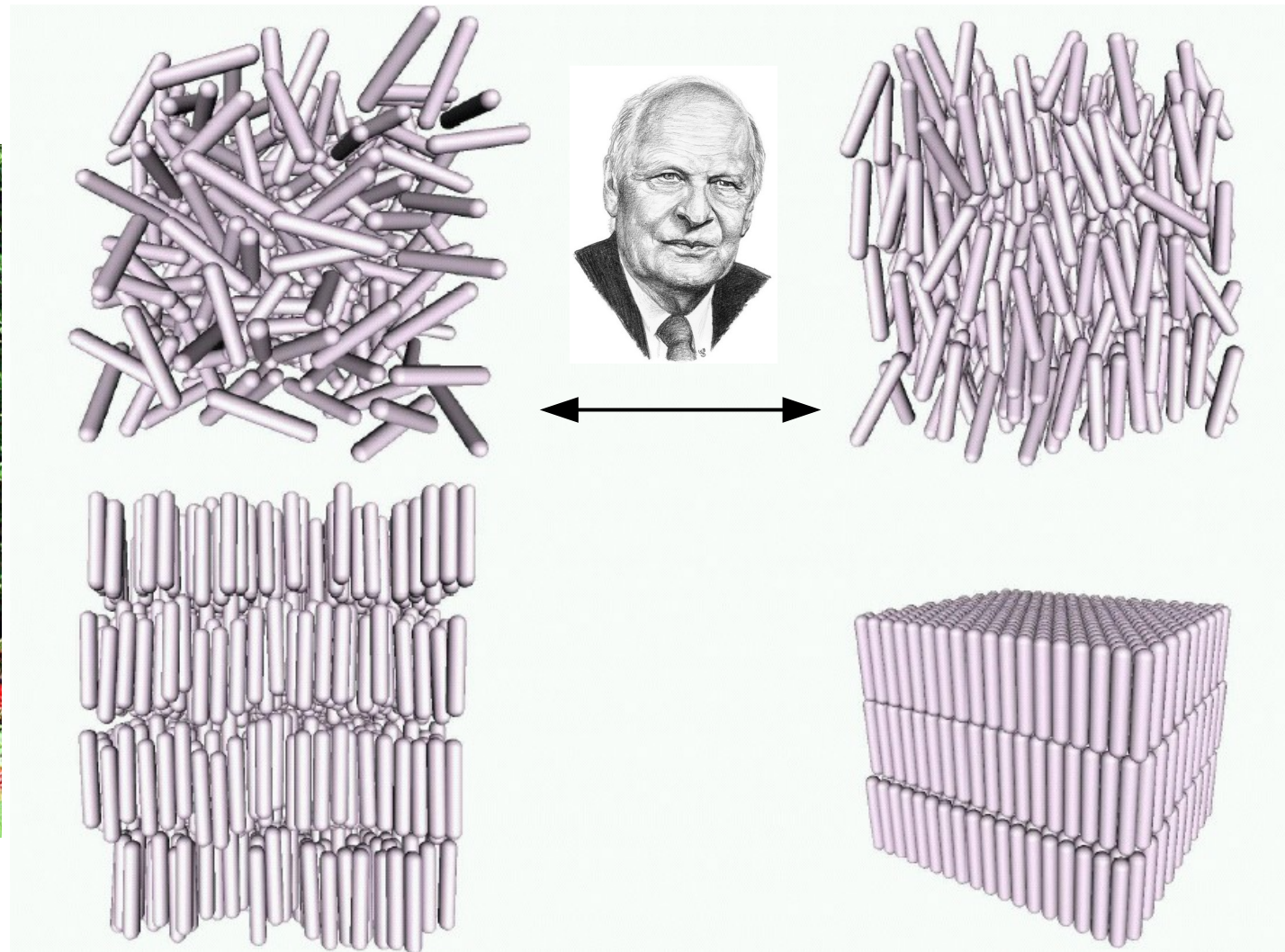


Image from the International Committee on Taxonomy of Viruses database

Phase behaviour of hard spherocylinders



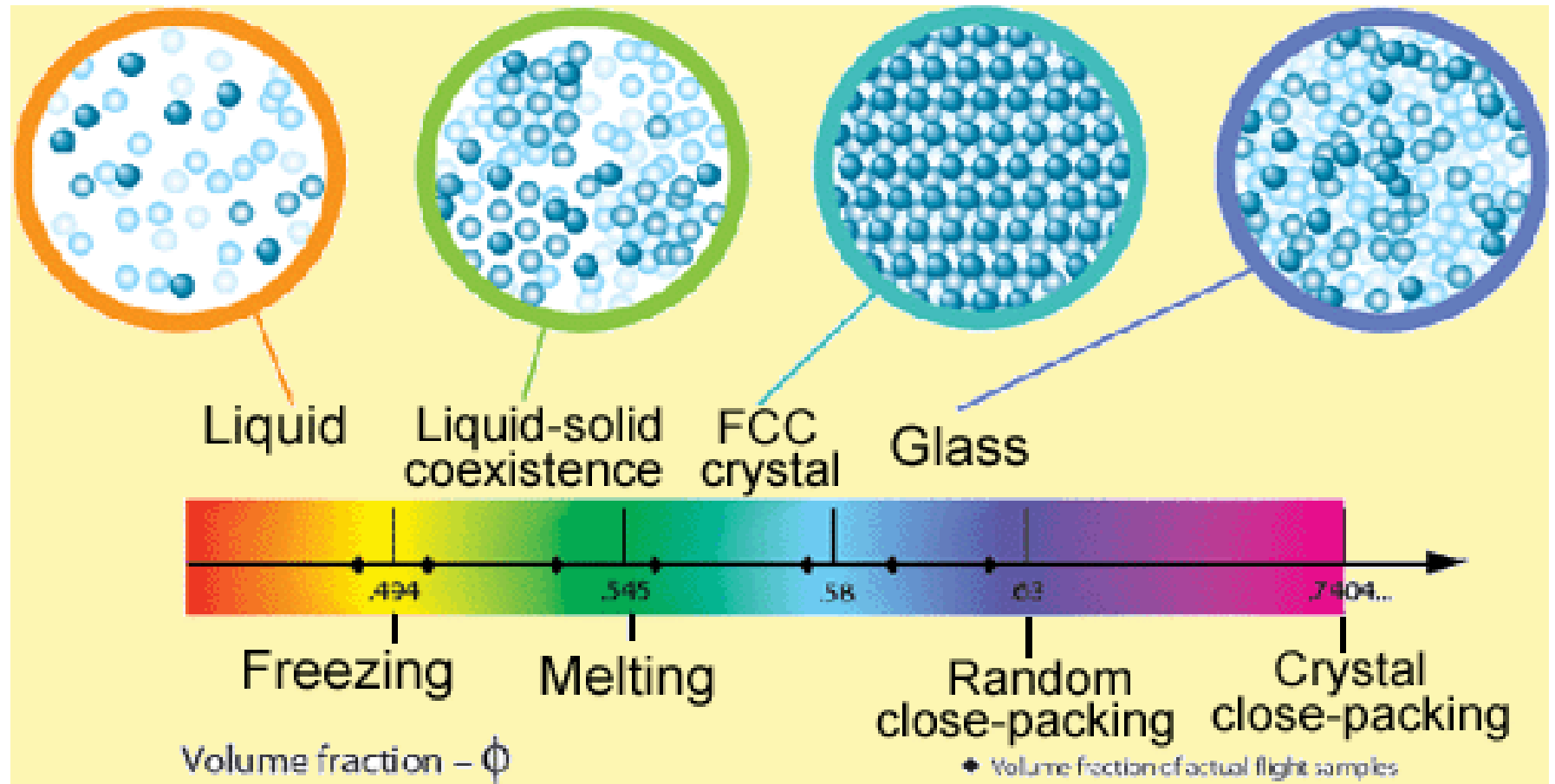
P. Bolhuis, 1996

The dark hand of entropy



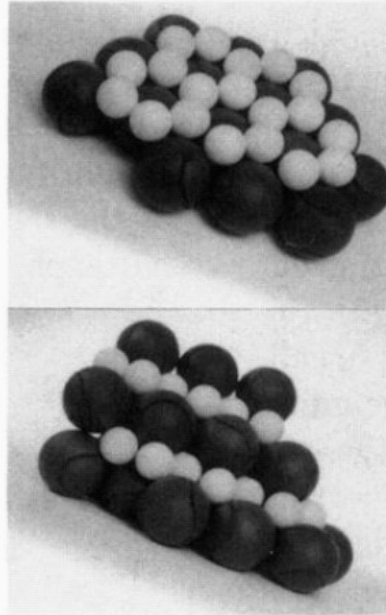
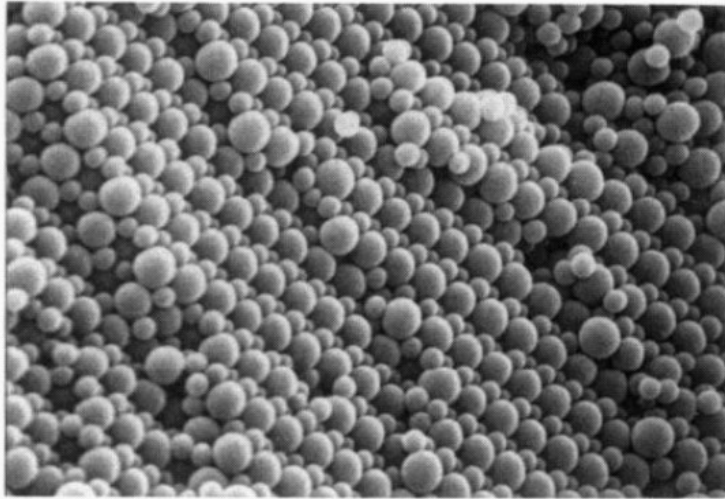
10 fundamental stamps...

The dark hand of entropy



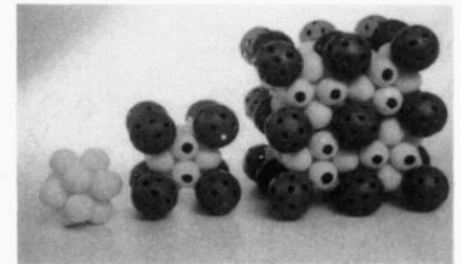
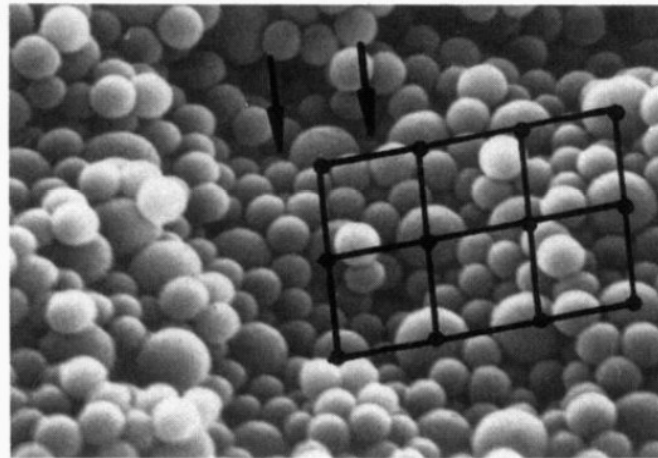
Entropic phase transition for hard spheres (NASA experiment)

Entropic phase transition for binary mixtures...



AB₂ super-structure
also observed in gem opals

AB₁₃ super-structure
also observed in gem opals



System : sterically stabilized PMMA spheres
P. Bartlett, R. Ottewill, P. Pusey, *Phys. Rev. Lett.* **68**, 3801 (1992)

Back to Ising model

- $d=1$: transfer matrix formalism
- $d=2$: Peierls + Onsager's exact solution
- $d=3$: renormalization group treatment



THE RENORMALIZATION GROUP AND CRITICAL PHENOMENA

Nobel lecture, 8 December 1982

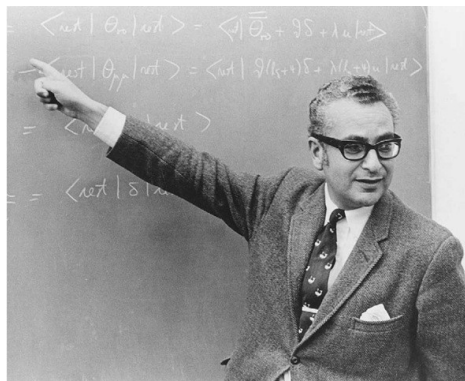
by

KENNETH G. WILSON

Laboratory of Nuclear Studies, Cornell University,
Ithaca, New York 14853



When I entered graduate school, I had carried out the instructions given to me by my father and had knocked on both Murray Gell-Mann's and Feynman's doors, and asked them what they were currently doing. Murray wrote down the partition function for the three dimensional Ising model and said it would be nice if I could solve it (at least that is how I remember the conversation). Feynman's answer was "nothing". Later, Jon Mathews explained some of



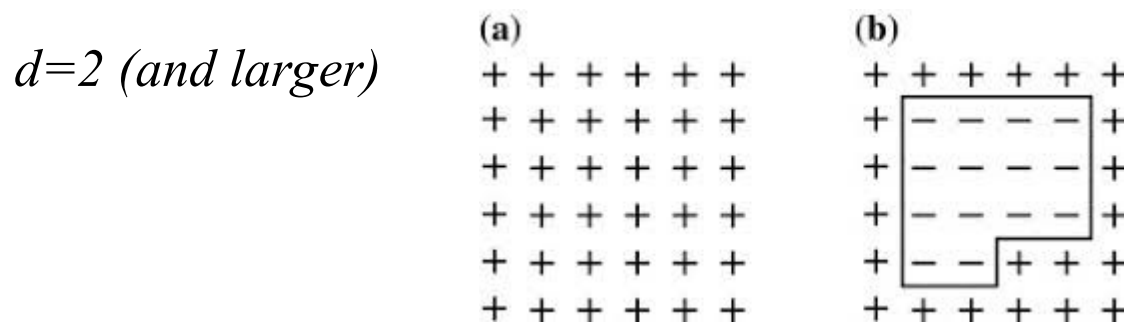
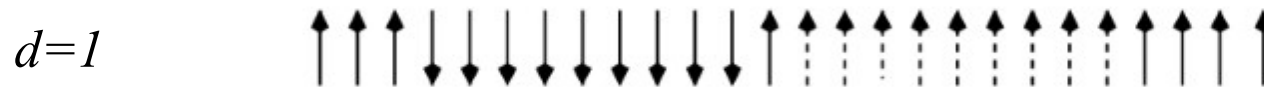
Nobel 1969



Nobel 1965

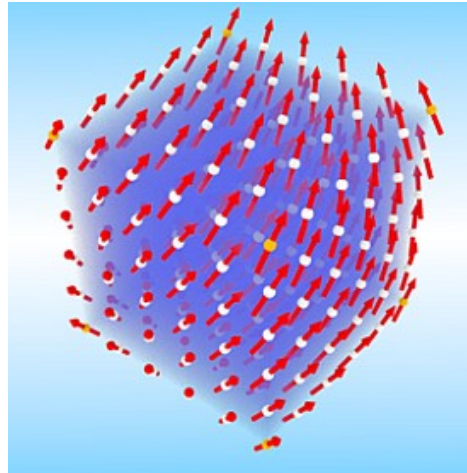
One and higher dimension

Where does the difference come from ?



- In 1d, the energetic cost of fluctuations is too small
Entropy dominates, fluctuations proliferate
- no phase transition
- notion of **lower critical dimension** : $d_{lower} = 1$ for usual Ising

Back to the $O(n)$ model with $n > 1$: the physics of spin waves



The Mermin Wagner theorem

No phase transition with continuous symmetry in
one and two dimensions

Chapter

1

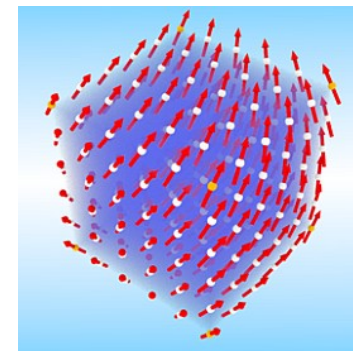
Introduction to phase transitions and critical phenomena

- 1- Problems raised by phase transitions, from a stat mech perspective
- 2- Classification of phase transitions
- 3- Ising model : the drosophila of phase transitions
- 4- Order parameter and symmetry breakdown
- 5- Local order and correlation functions

Conclusion of chapter

- Phase transition \leftrightarrow discontinuity or singularity in free energy derivative
- No phase transition in a finite system
(so... what about computer simulations? \rightarrow see the tutorial on Binder cumulants)
- Phase transition means long range order...
although the correlation function is short range outside T_c
- Correlation length diverges at T_c : microscopic details no longer matter
 \rightarrow universality \rightarrow renormalization group treatment
- Notion of lower critical dimension :
1 for discrete spins, 2 in continuous case (Mermin Wagner)

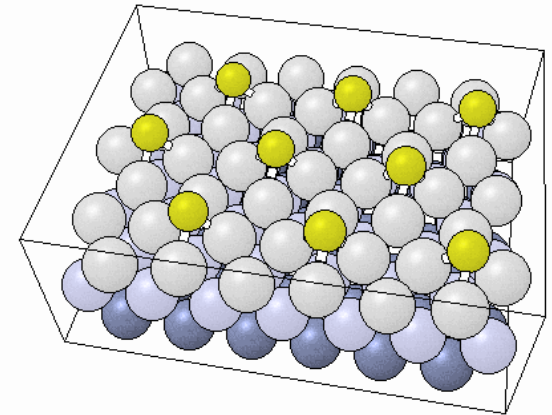
Goldstone modes
destabilize order



Beyond universality, Ising also important through isomorphism of models

- Binary alloys
- Lattice gas model, for liquid/gas transition ; or adsorption of H onto iron

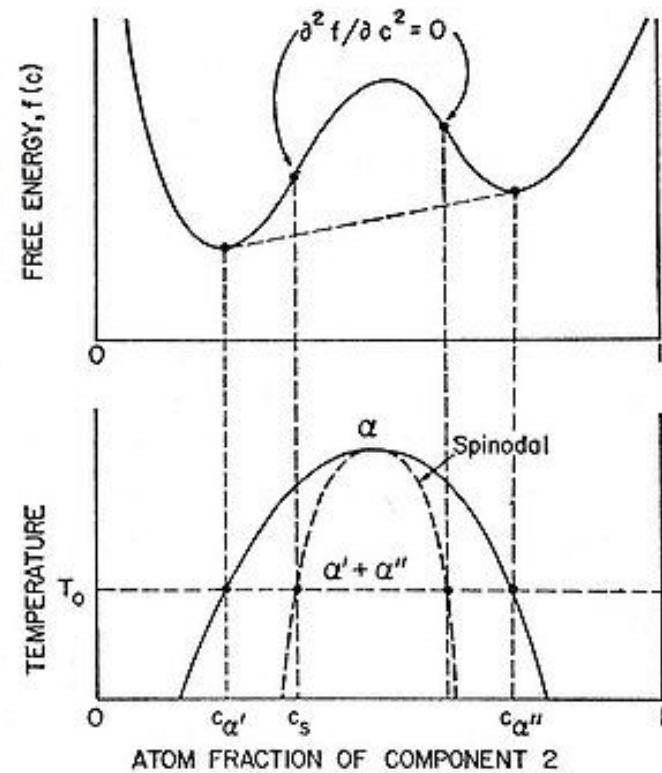
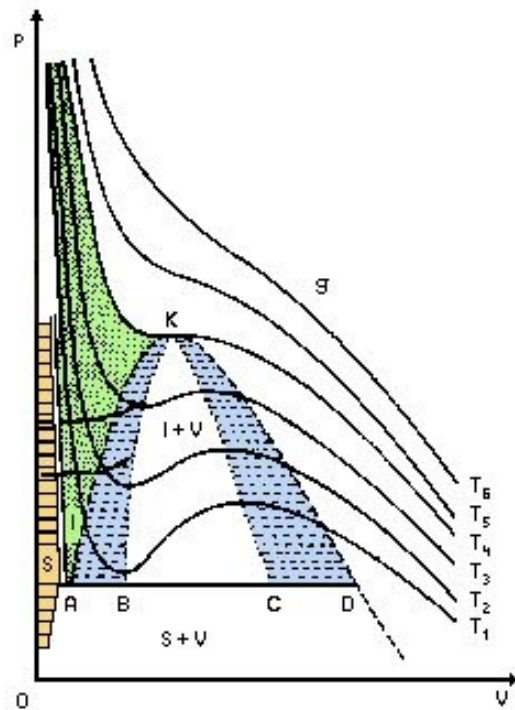
Fe surface (110)

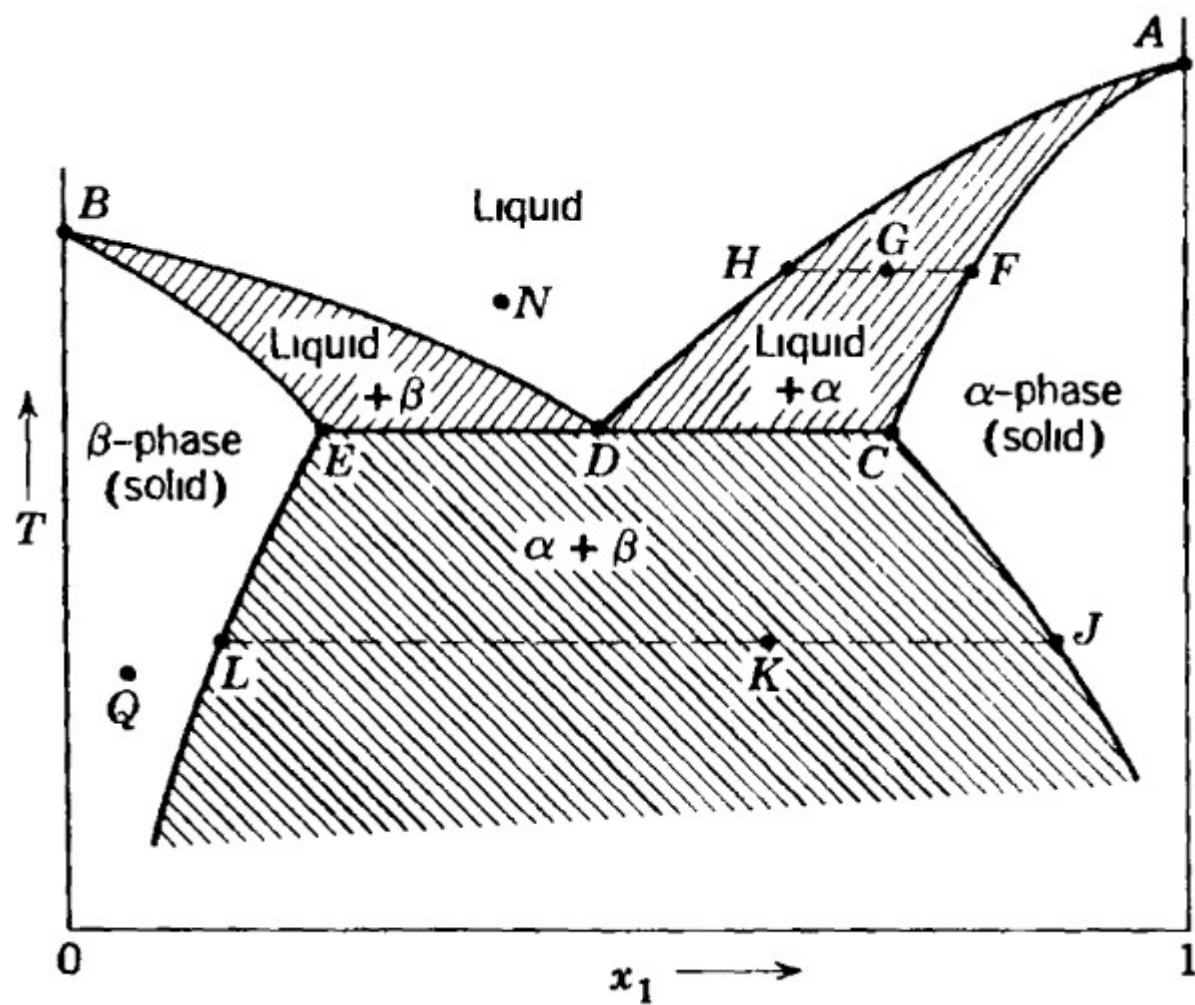


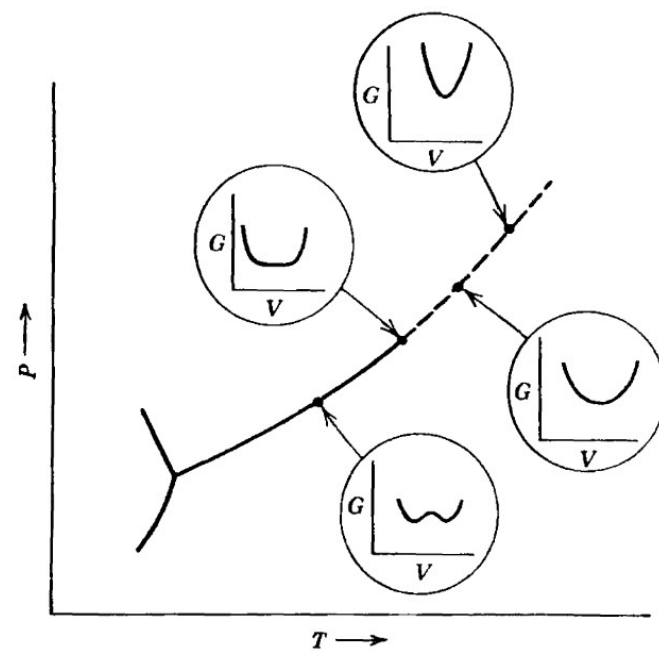
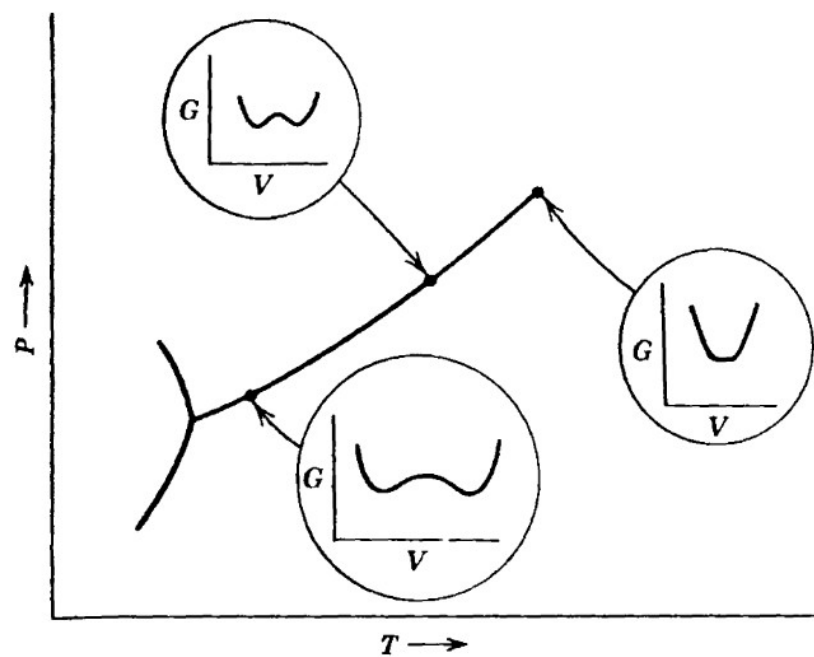
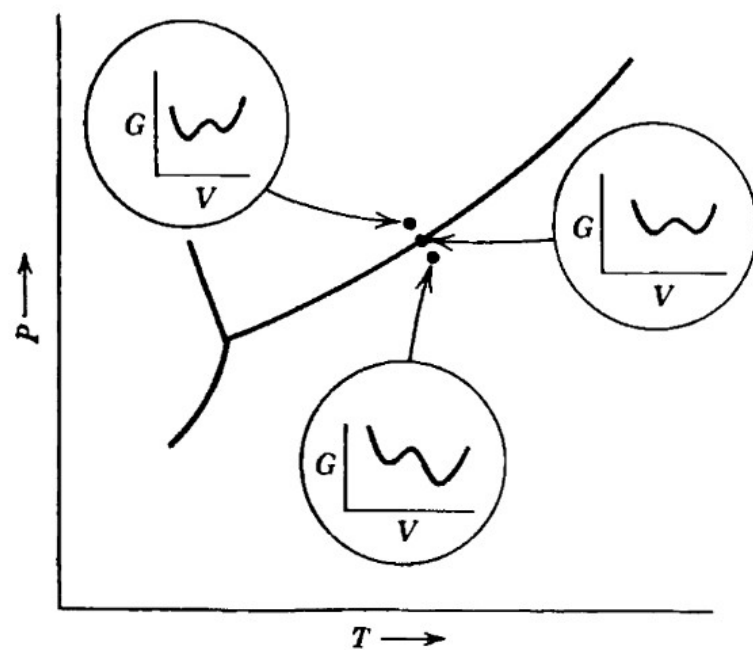
- Fads/herds and hypes (*effets de mode*)
- Neuronal activity
- Protein folding,
- Bird flocking...
and much more



Spinodal and binodal

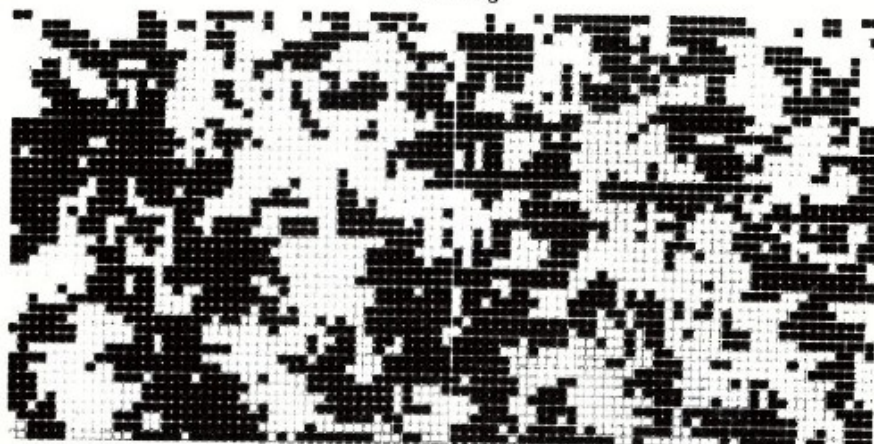




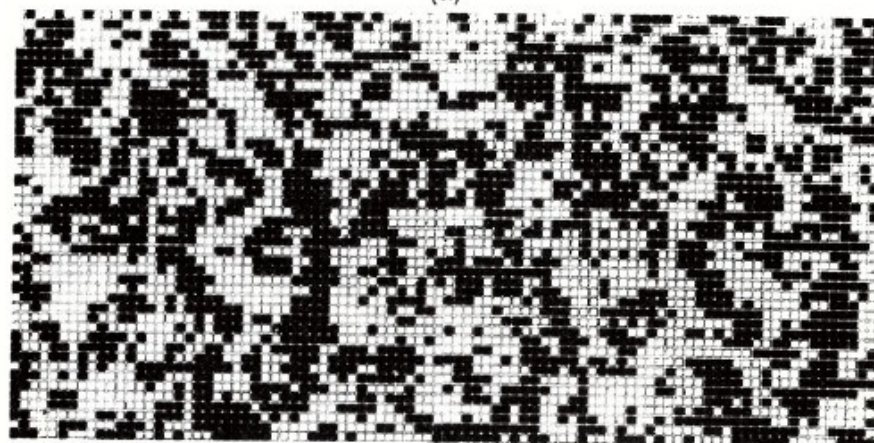




$T=1.22 T_c$



(a)



(b)



(c)

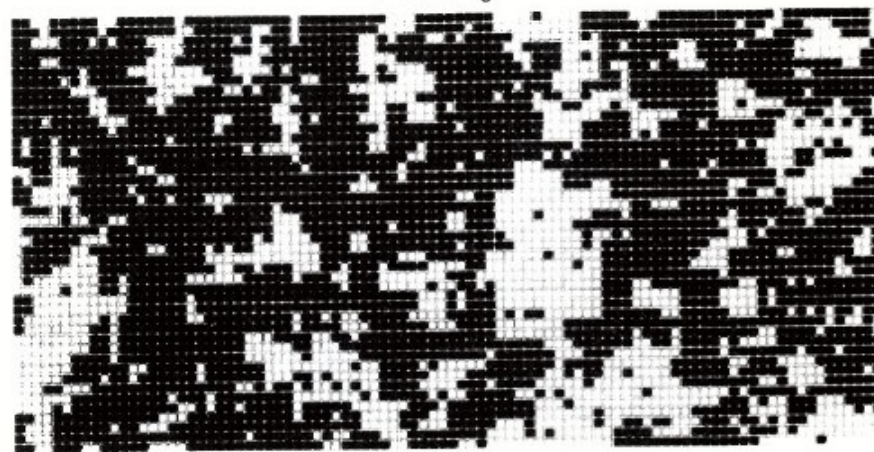


(d)

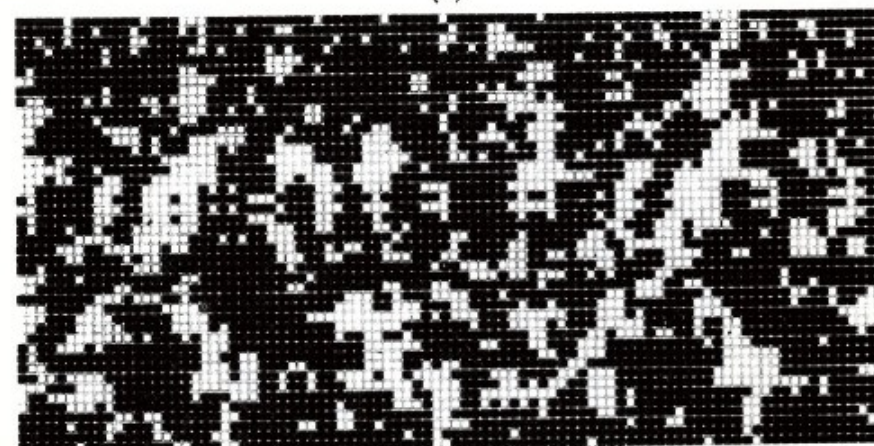


(e)

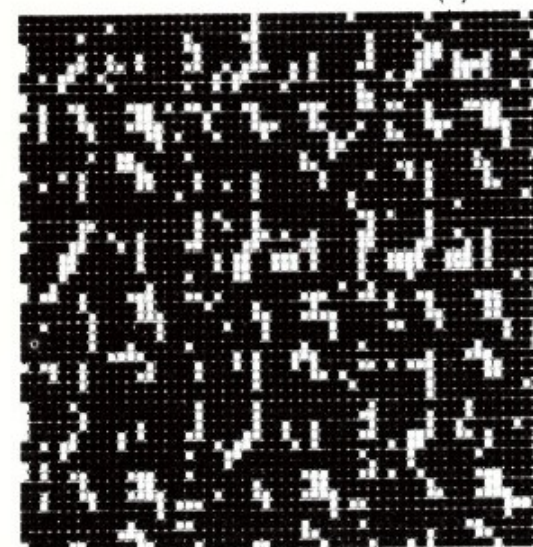
$T=0.99 T_c$



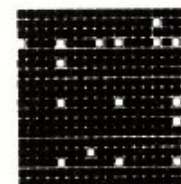
(a)



(b)



(c)



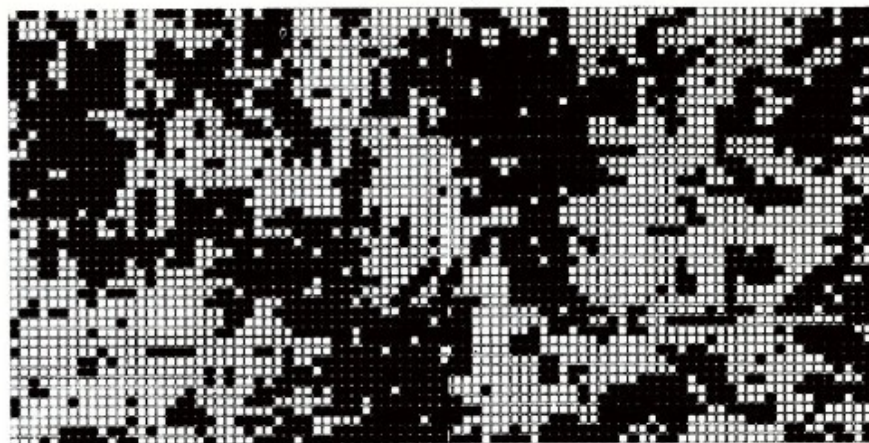
(d)



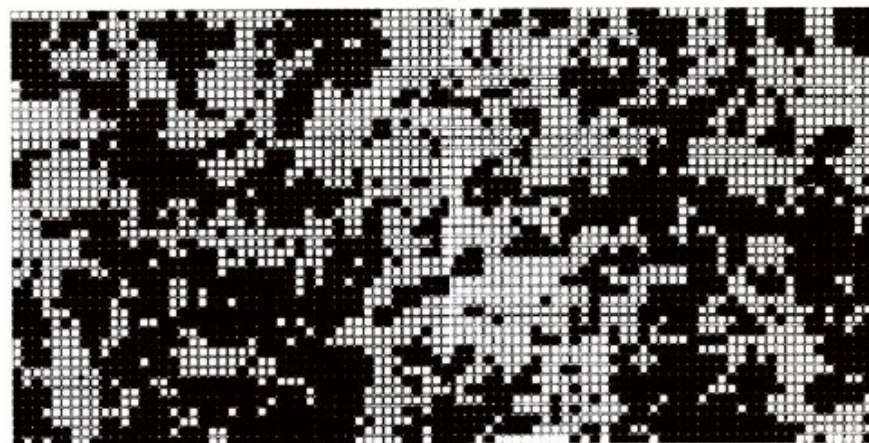
(e)



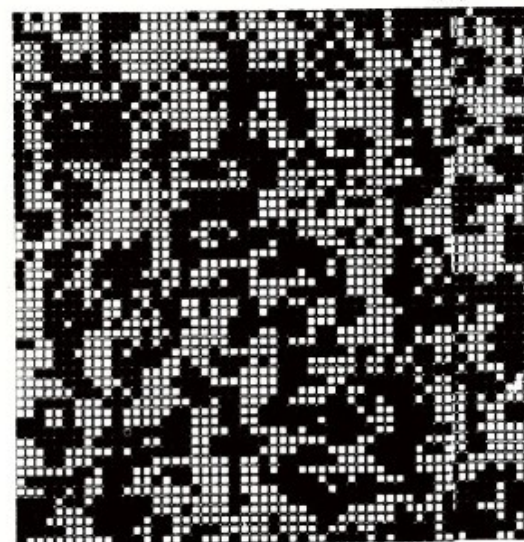
$$T = T_c$$



(a)



(b)



(c)



(d)



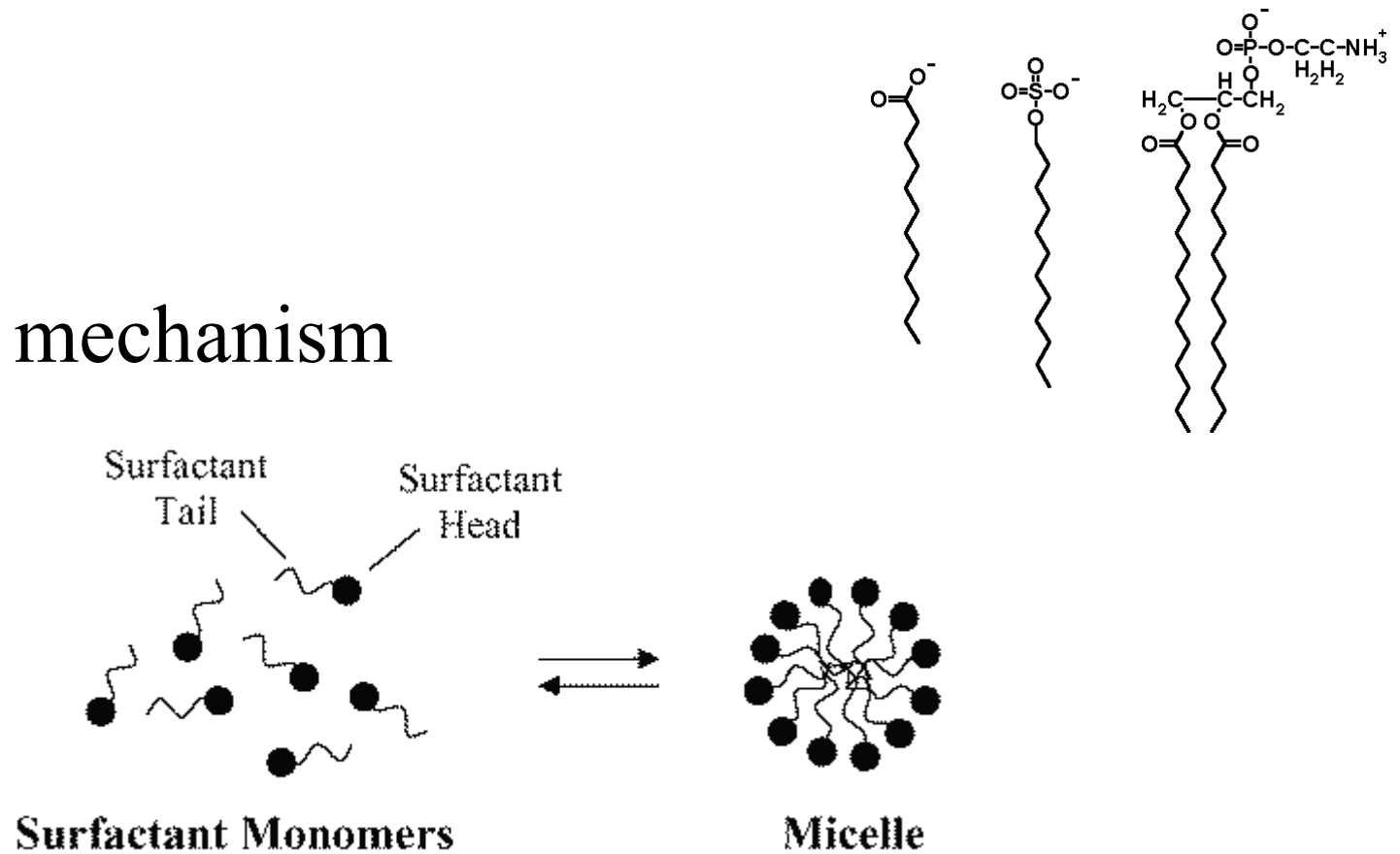
(e)

Backup slides



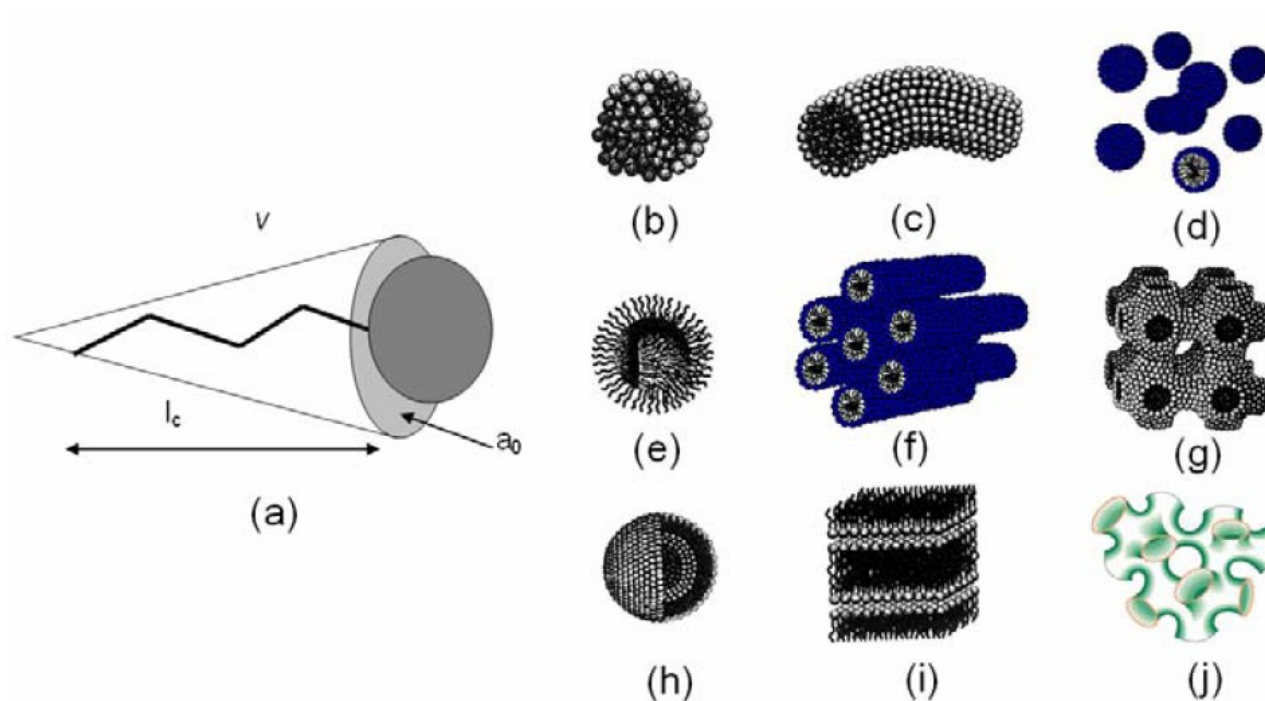
Surfactant molecules and micelle formation

Basic mechanism



Opens for considerable richness

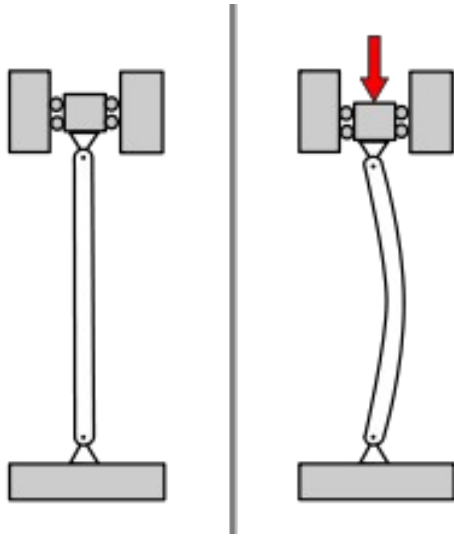
Various surfactant phases



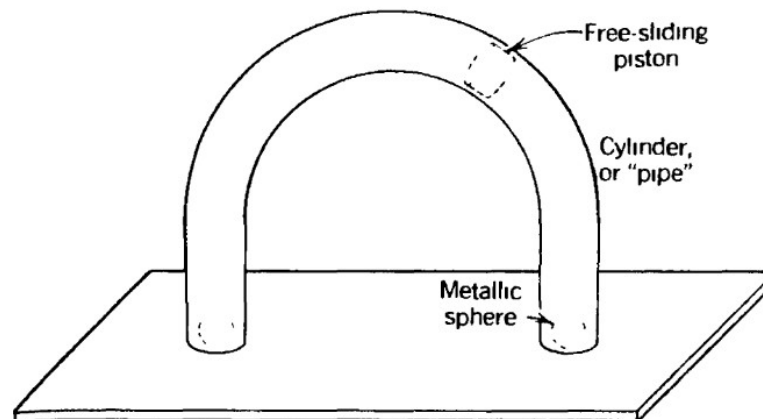
- | | |
|-----------------------------|----------------------------------|
| (a) An amphiphilic molecule | (f) Hexagonal phase |
| (b) Spherical micelle | (g) Bicontinuous cubic structure |
| (c) Cylindrical micelle | (h) Vesicle |
| (d) Cubic phase | (i) Lamellar phase |
| (e) Inverse micelle | (j) Sponge phase |

a whole zoo

Mechanical symmetry breakdown



Buckling



I Introduction to phase transitions and critical phenomena

- 1- The problems raised by phase transitions, from a statistical mechanics perspective
- 2- Classification of phase transitions
- 3- Ising model: why ?
- 4- Order parameter and symmetry breakdown
- 5- Local order and correlation functions

II First order phase transitions (*mostly treated as a tutorial*)

- 1- Unstable isotherms, double-tangent and Maxwell construction
- 2- Spinodal and binodal

III Critical phenomena : qualitative approaches

- 1- Weiss molecular field
- 2- Variational mean-field and critical exponents
- 4- Landau theory
- 5- Correlation functions and Ginsburg-Landau functional

IV Beyond mean-field: fluctuations and scaling

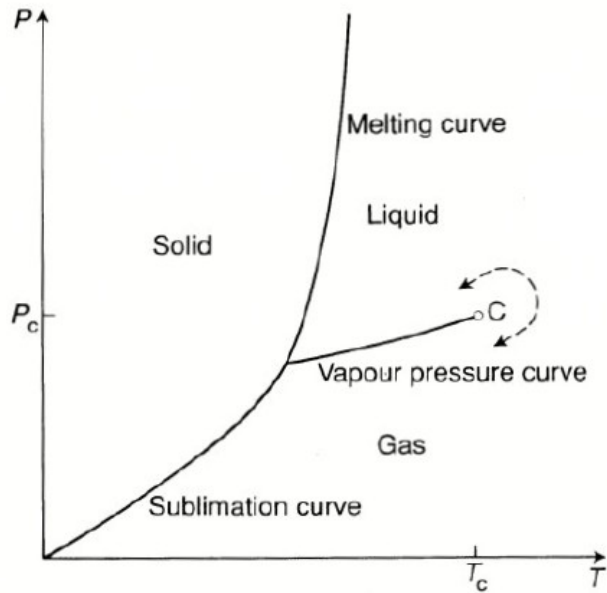
- 1- Fluctuations
 - What do mean-field practitioners really do ?
 - Fluctuation correction to the saddle-point
 - Ginzburg criterion, crossover behaviour
 - Scattering and fluctuations: measure of structure factors
- 2- The scaling hypothesis: life with a large correlation length
 - Homogeneity and scaling relations
 - Finite size scaling: turning a drawback into an advantage

V Renormalisation group ideas

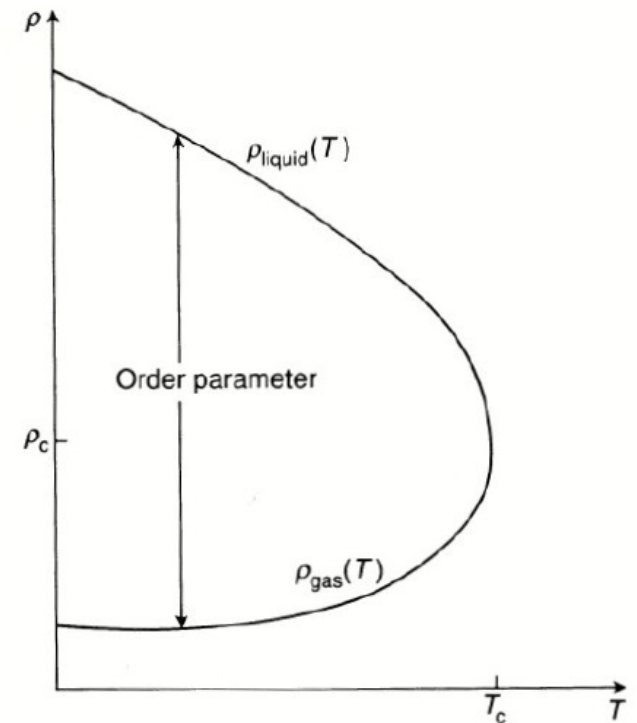
- 1- What are the problems ?
- 2- Definition of a renormalisation group transformation
- 3- Fixed points and universality
- 4- Scale invariance, critical exponents

Identify an ORDER PARAMETER

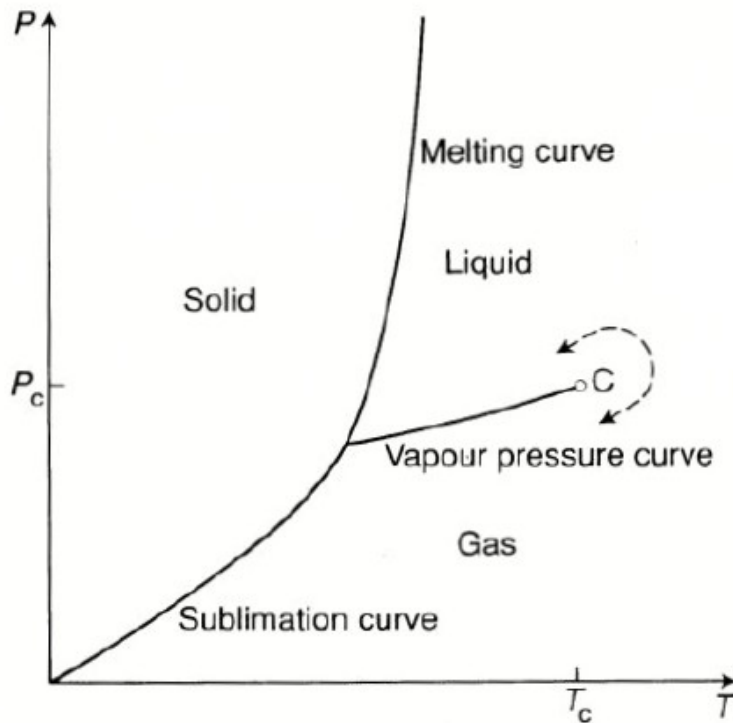
pressure



density



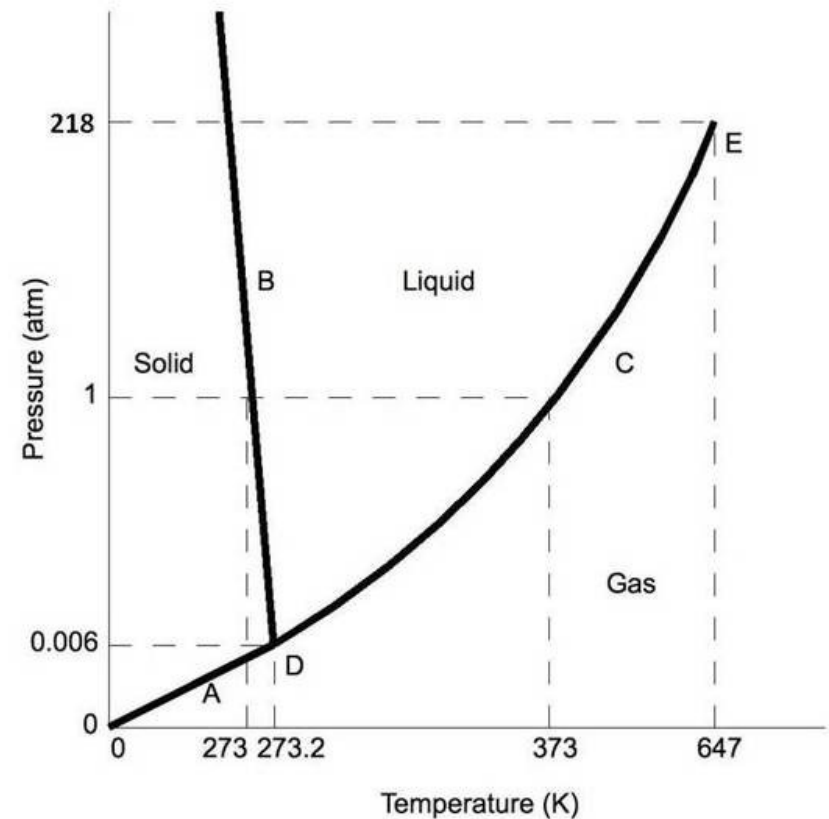
Examples and exceptions



Generic case

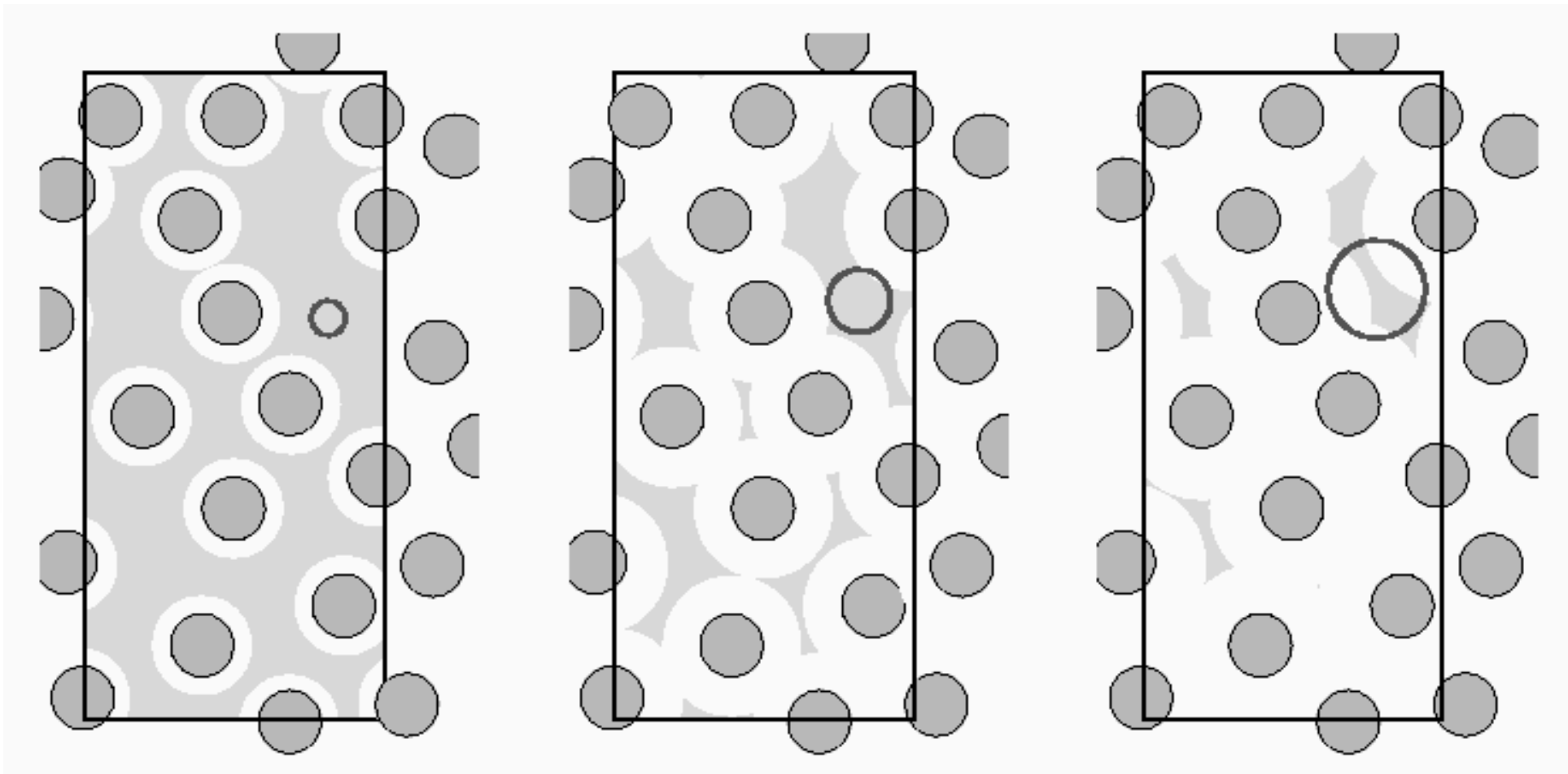
$$\frac{dP}{dT}(\text{coex}) = \frac{1}{T} \frac{L_{\text{liq} \rightarrow \text{gas}}}{v_{\text{gas}} - v_{\text{liq}}}$$

Water



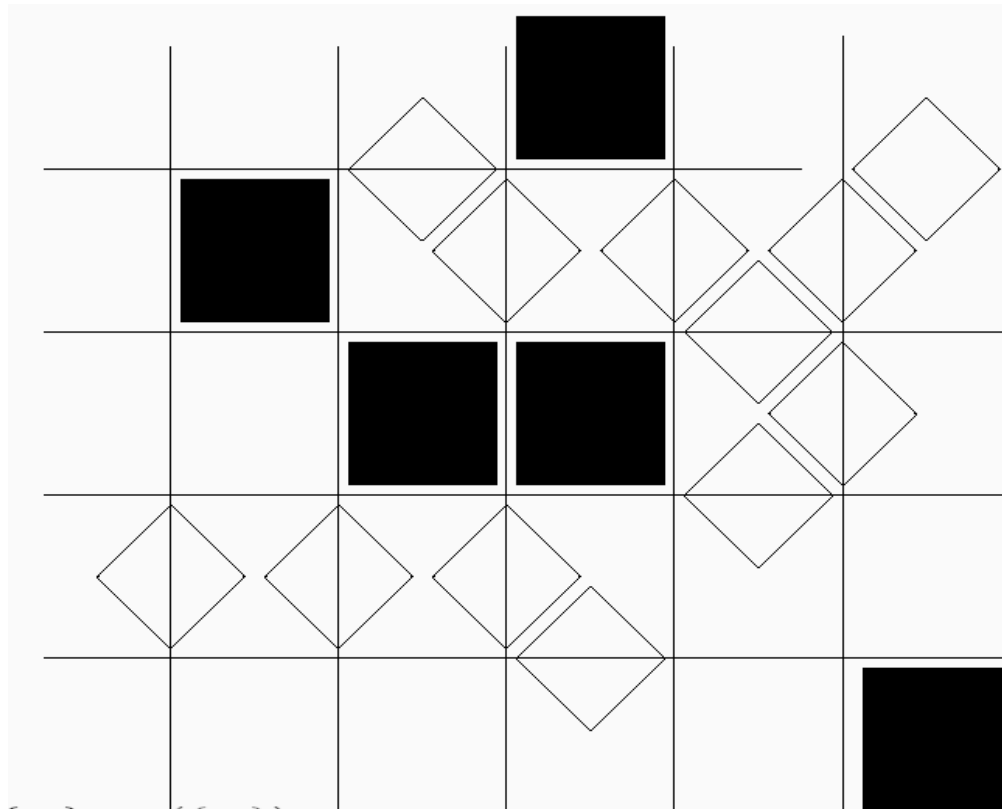
The excluded volume

Increasing the size of a tagged particle



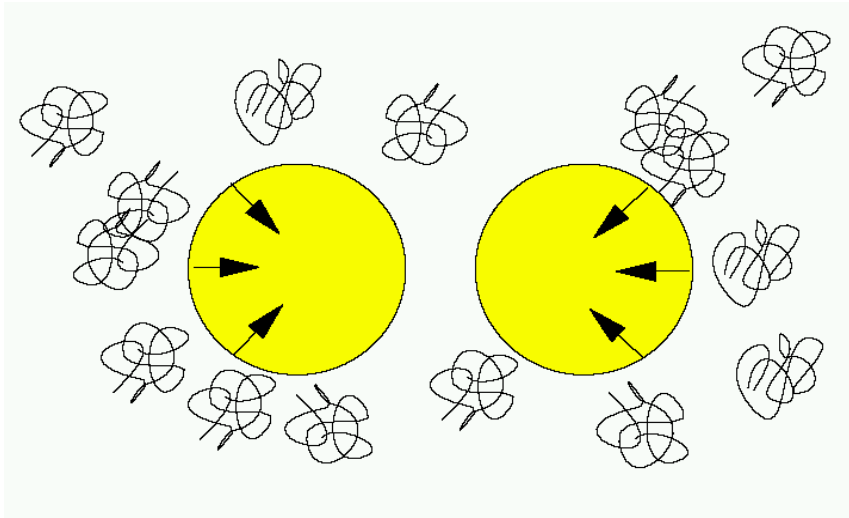
after Barrat/Hansen

Depletion forces: an exactly solvable lattice model



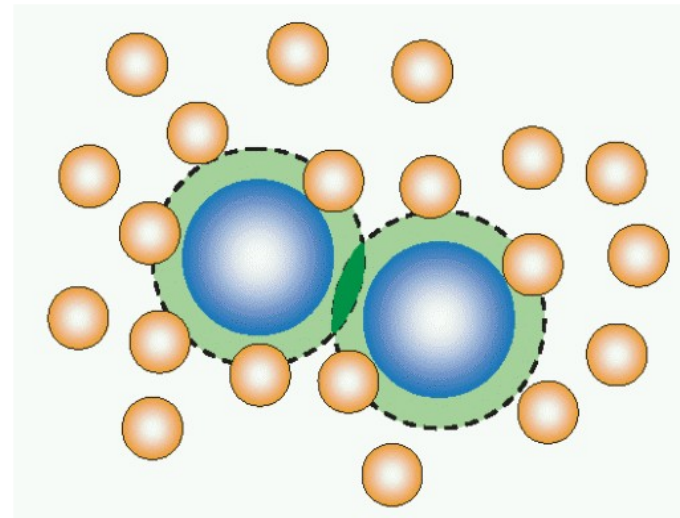
Frenkel and Louis, *Phys. Rev. Lett.* **68**, 3363 (1992)

Sphere-sphere depletion potential

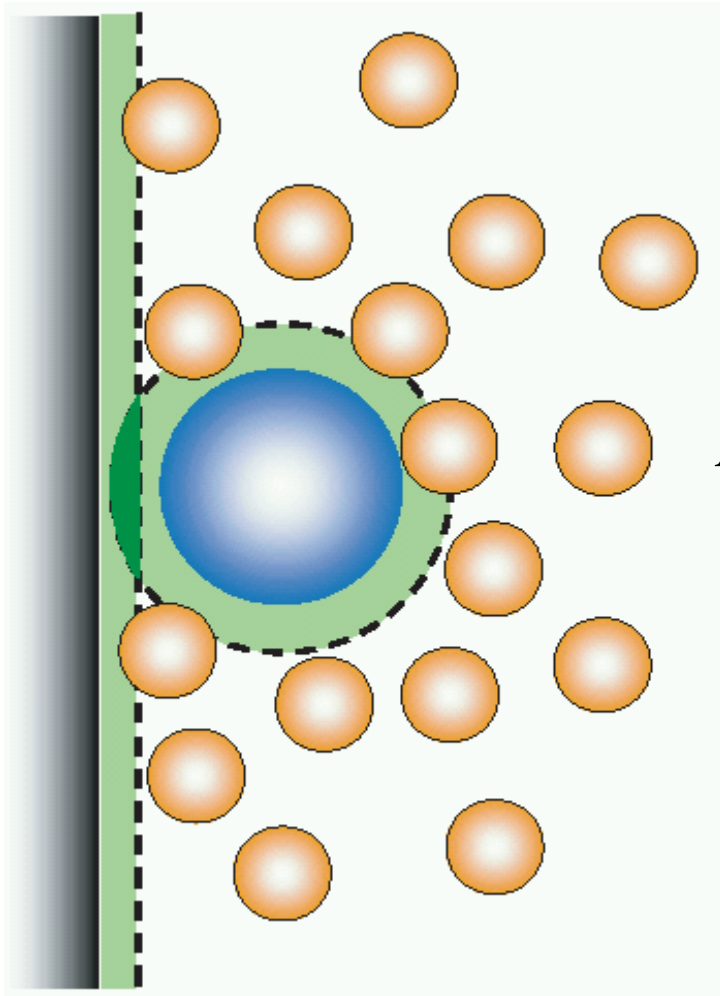


Colloid-polymer mixture

Mechanical and entropic
interpretations fully coincide

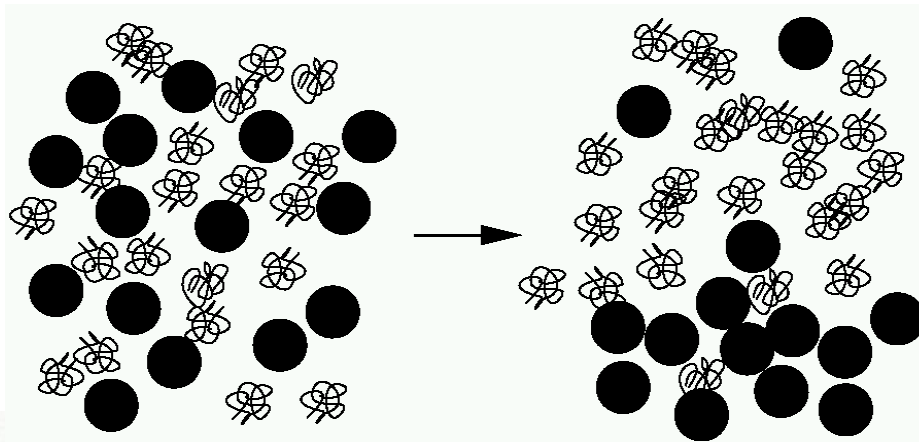


Sphere-plane depletion potential

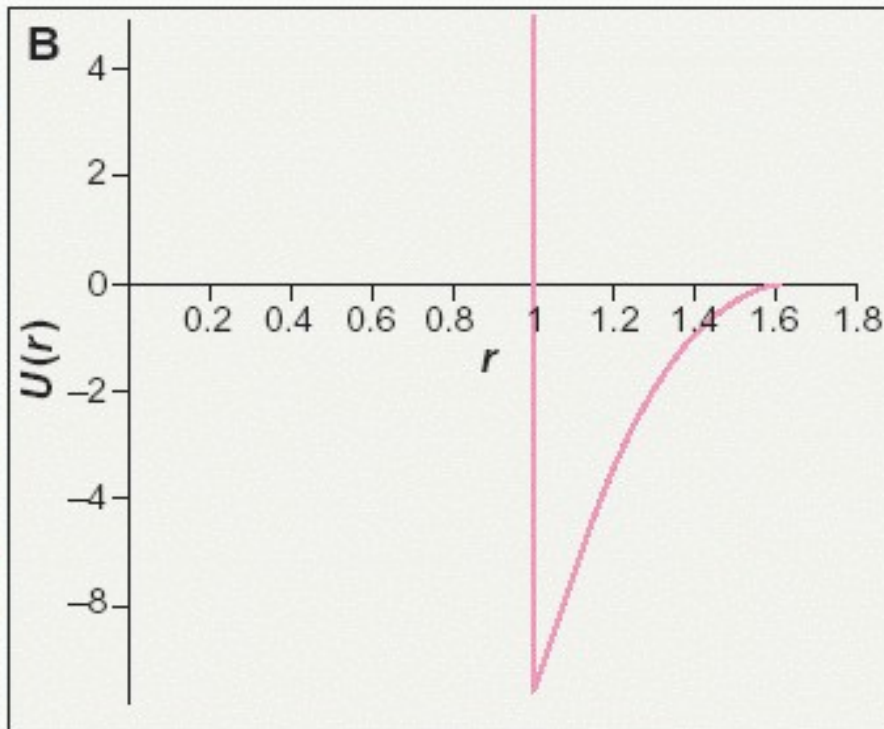
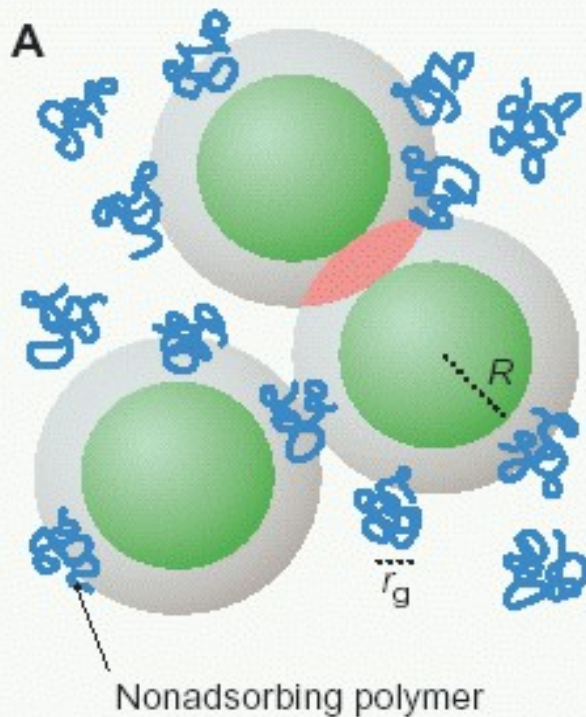


Again... uncompensated pressure
overlap of exclusion zones

Polymer induced flocculation



expected mechanism



PMMA spheres (70 nm) + polystyrene

Chapter

2

Chapter

Chapter

3

Chapter

4

Chapter

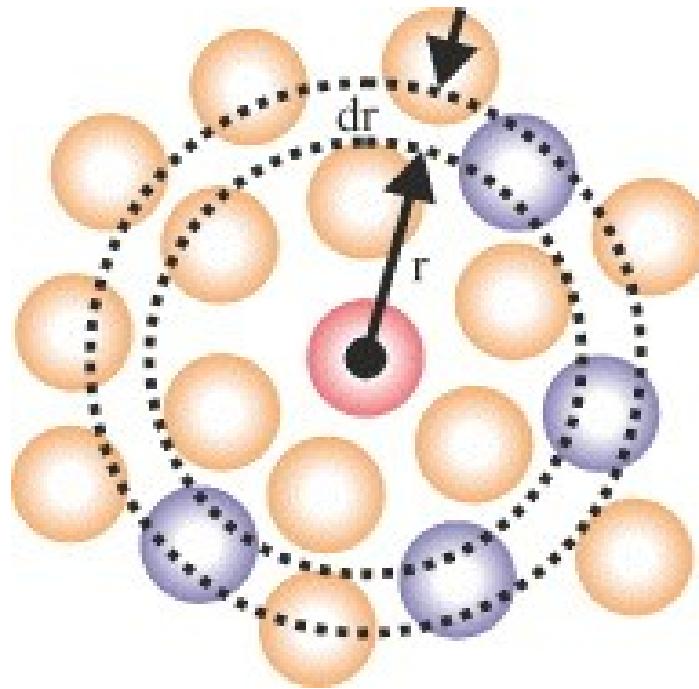
5

Chapter

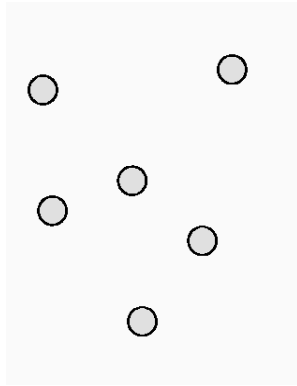
7

6

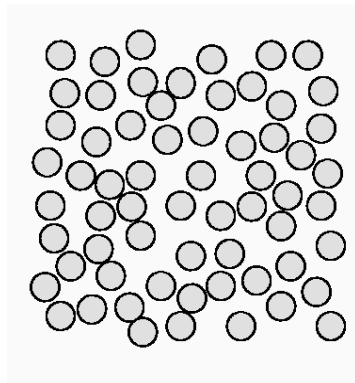
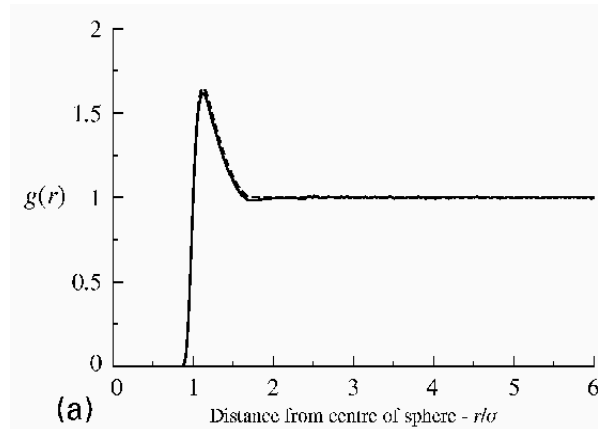
The pair correlation function $g(r)$



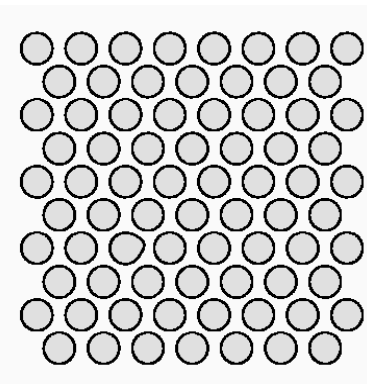
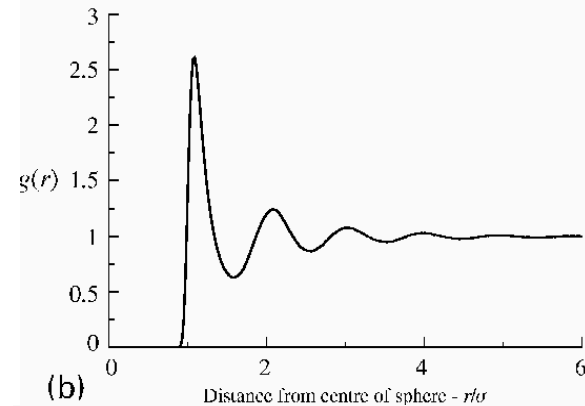
For a Lennard-Jones fluid



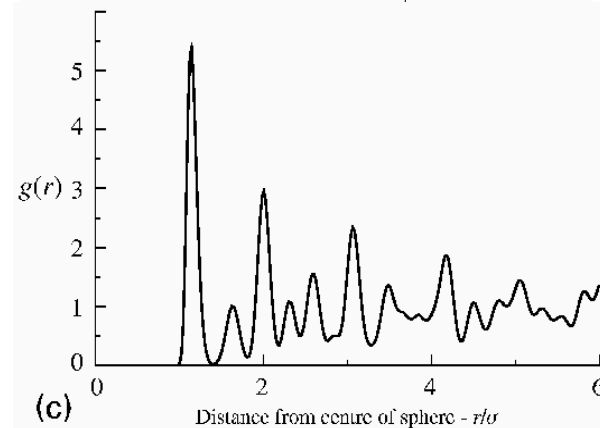
Gas



Liquid

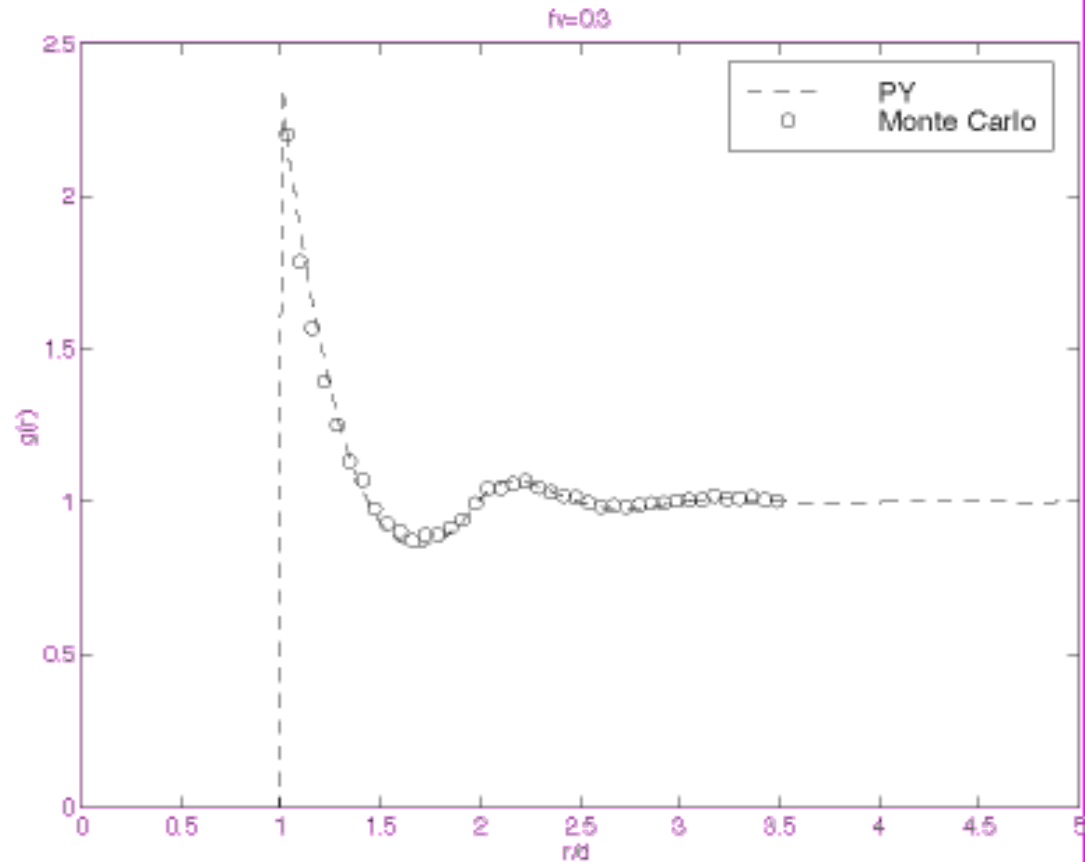


Solid



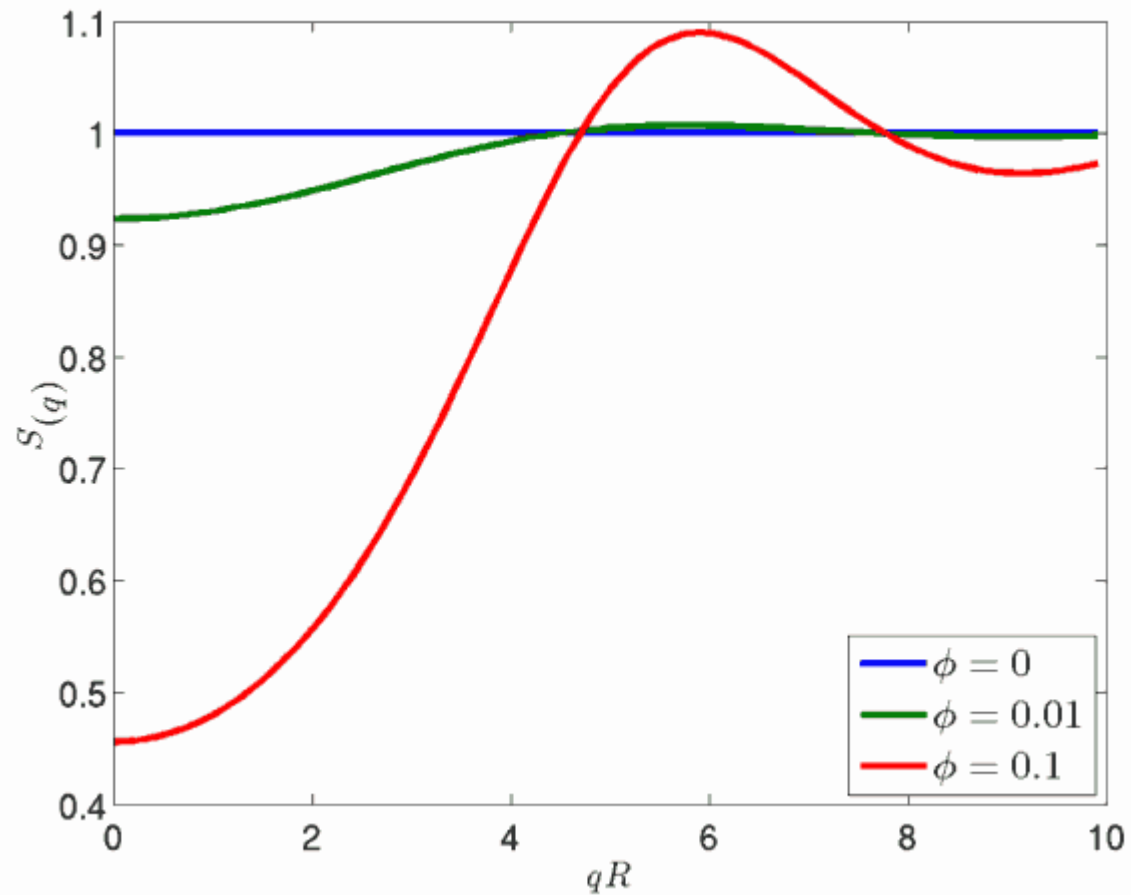
From Barrat/Hansen, *Basic concepts for simple and complex fluids* (2003)

The pair distribution function



Hard spheres ; maximum at contact

Structure factor $S(k)$



Hard spheres