

Enhancing Superconductivity using Light

Stefan Kaiser

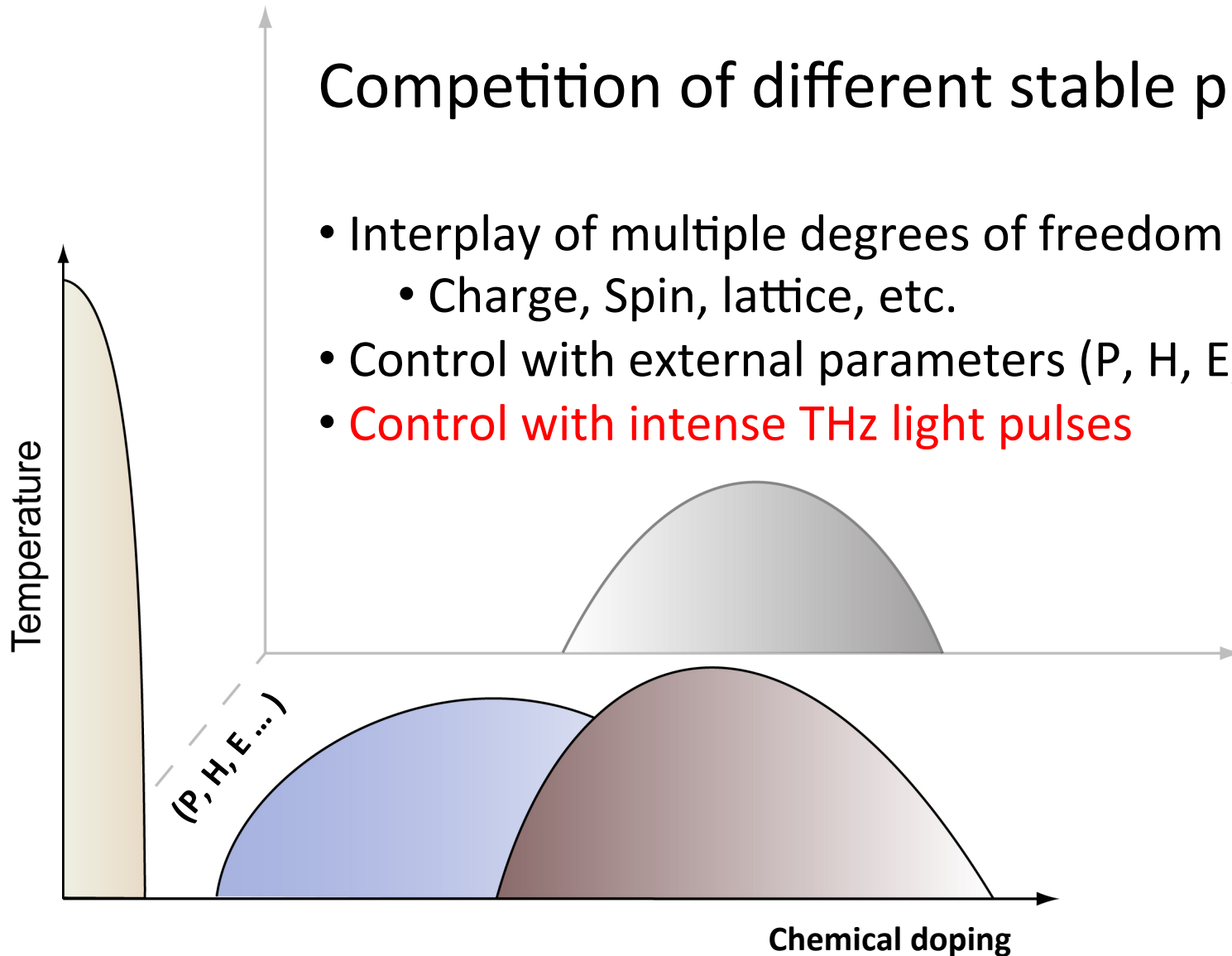
Max Planck Institute for the Structure and
Dynamics of Matter
Hamburg, Germany



Controlling complex matter

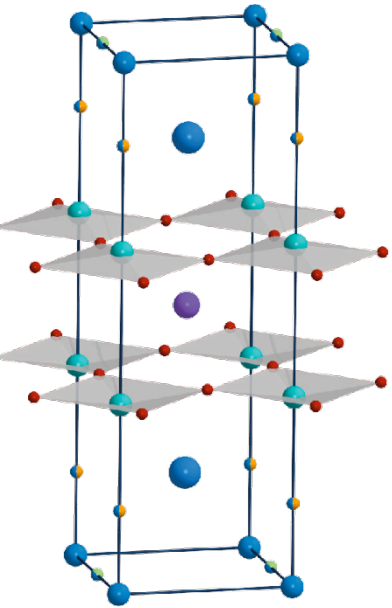
Competition of different stable phases

- Interplay of multiple degrees of freedom
 - Charge, Spin, lattice, etc.
- Control with external parameters (P, H, E, x, ...)
- **Control with intense THz light pulses**

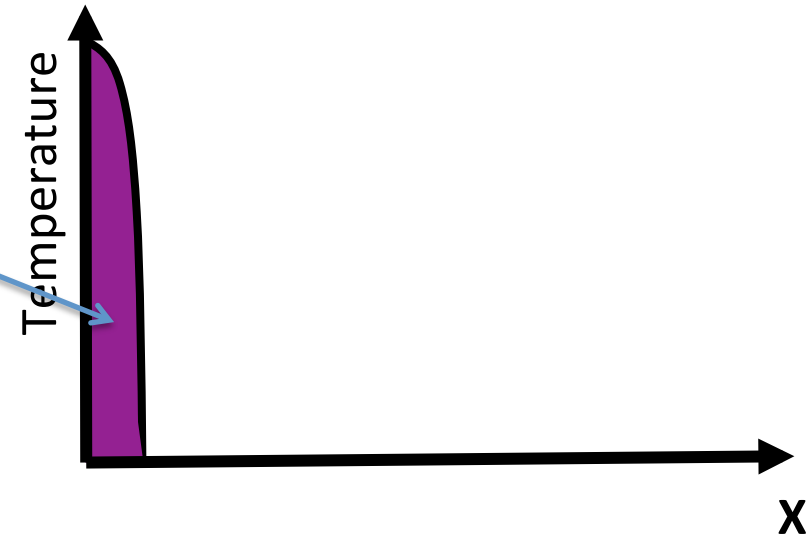
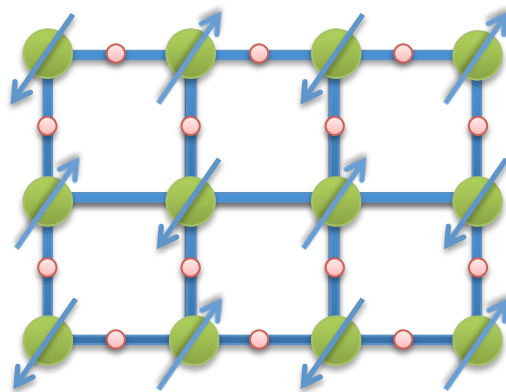


Cuprate-Superconductors

- Layered materials
- 2D-doped Mott insulator

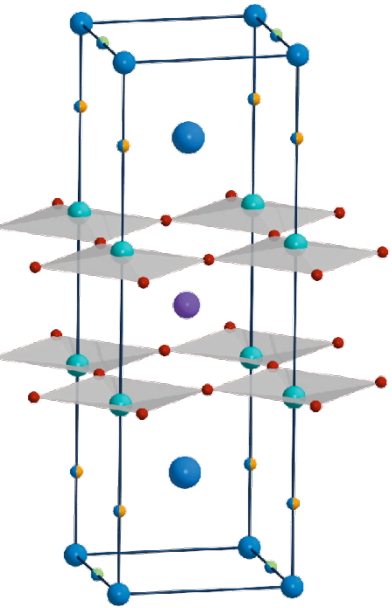


CuO-plane

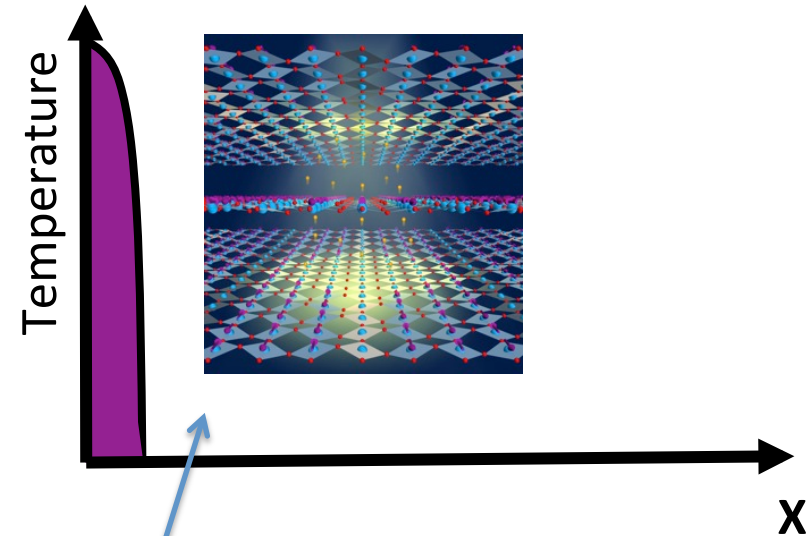
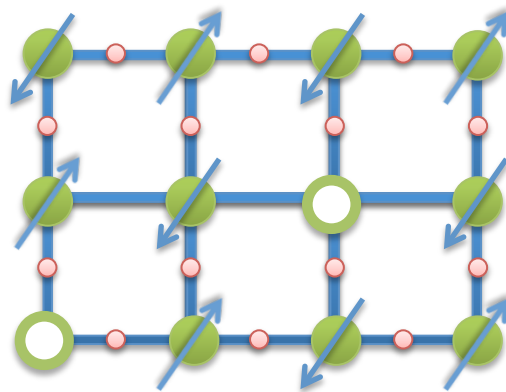


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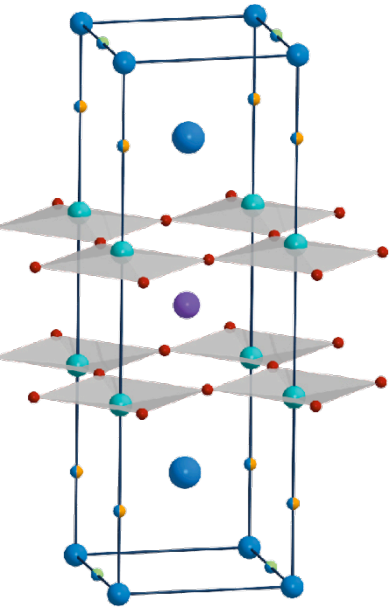
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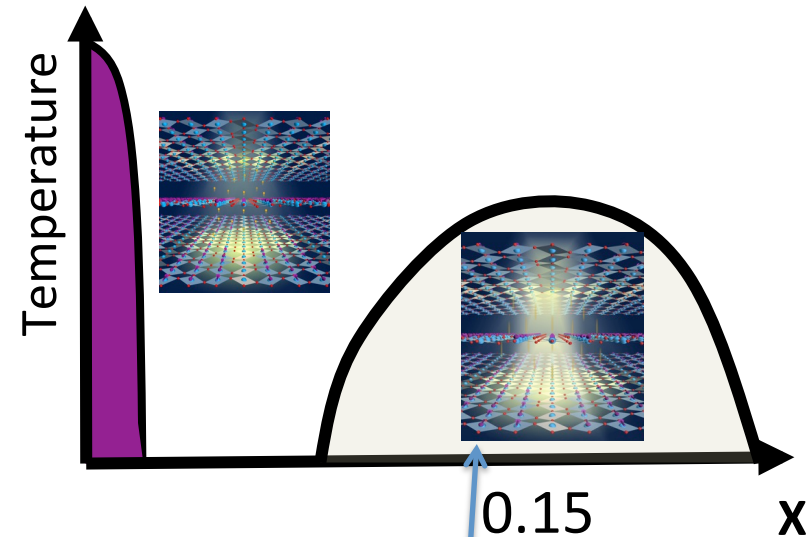
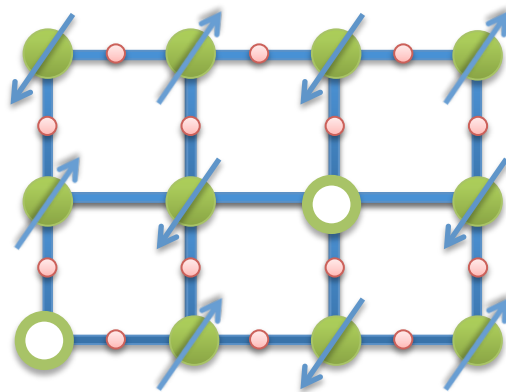
- Unconventional 2D metallic phase

Cuprate-Superconductors

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CuO-plane



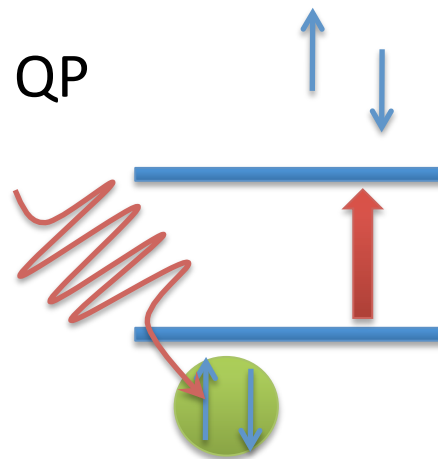
- Unconventional 2D metallic phase

- 3D-Superconductivity

Time resolved pump-probe spectroscopy

- Spectroscopy of non equilibrium QP

1.5 eV ~ 20,000 K



- THz control: Collective excitations, phonons

1 THz ~ 4 meV ~ 50 K

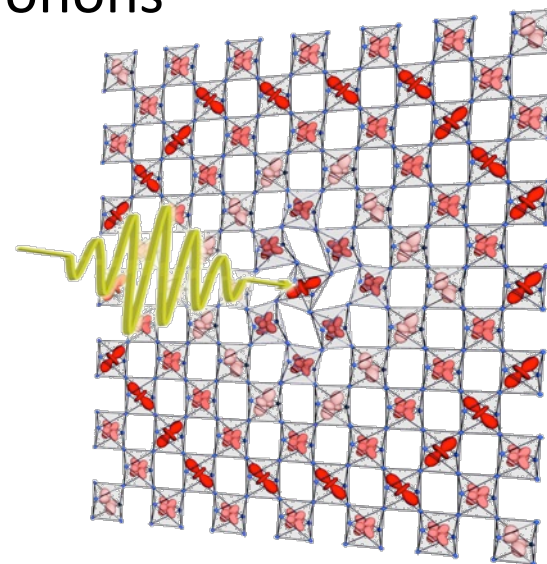


Photo-doping dynamics of cuprates

- Recombination dynamics

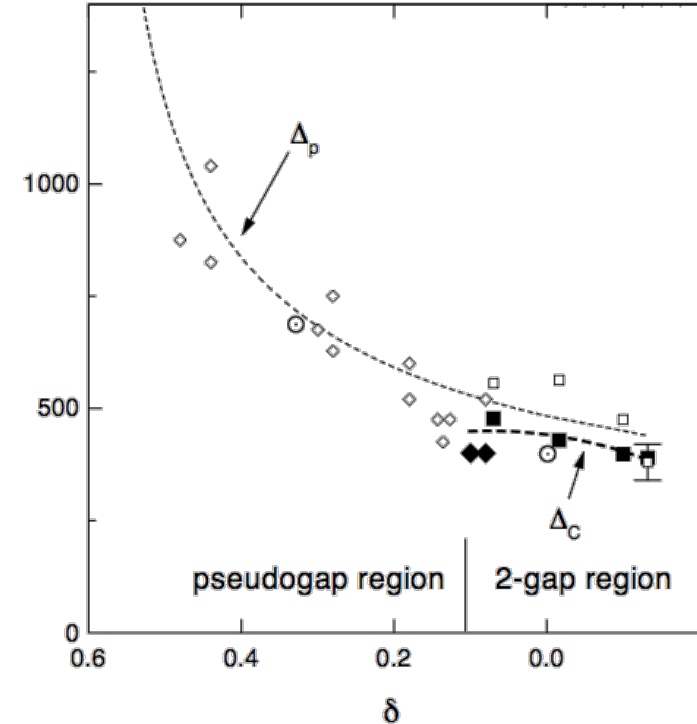
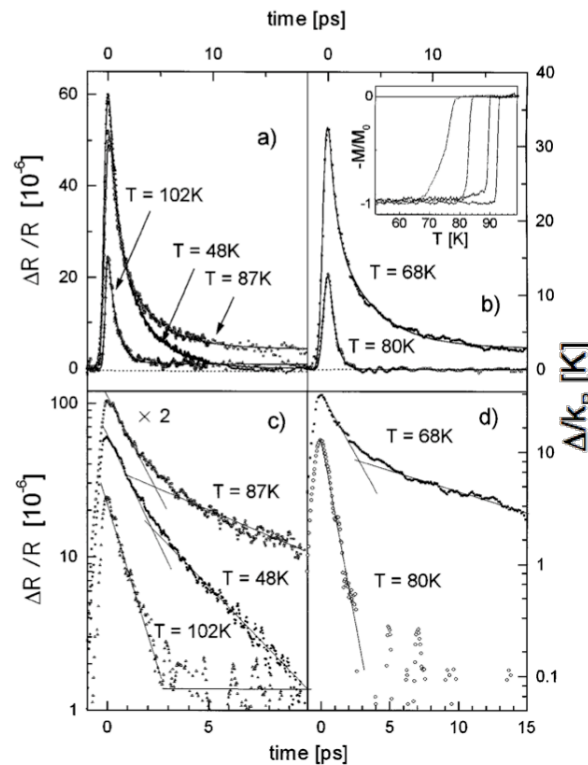
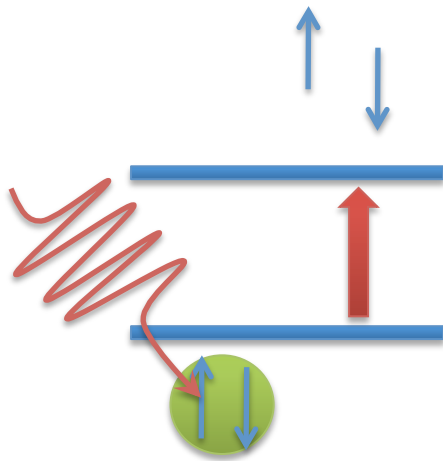
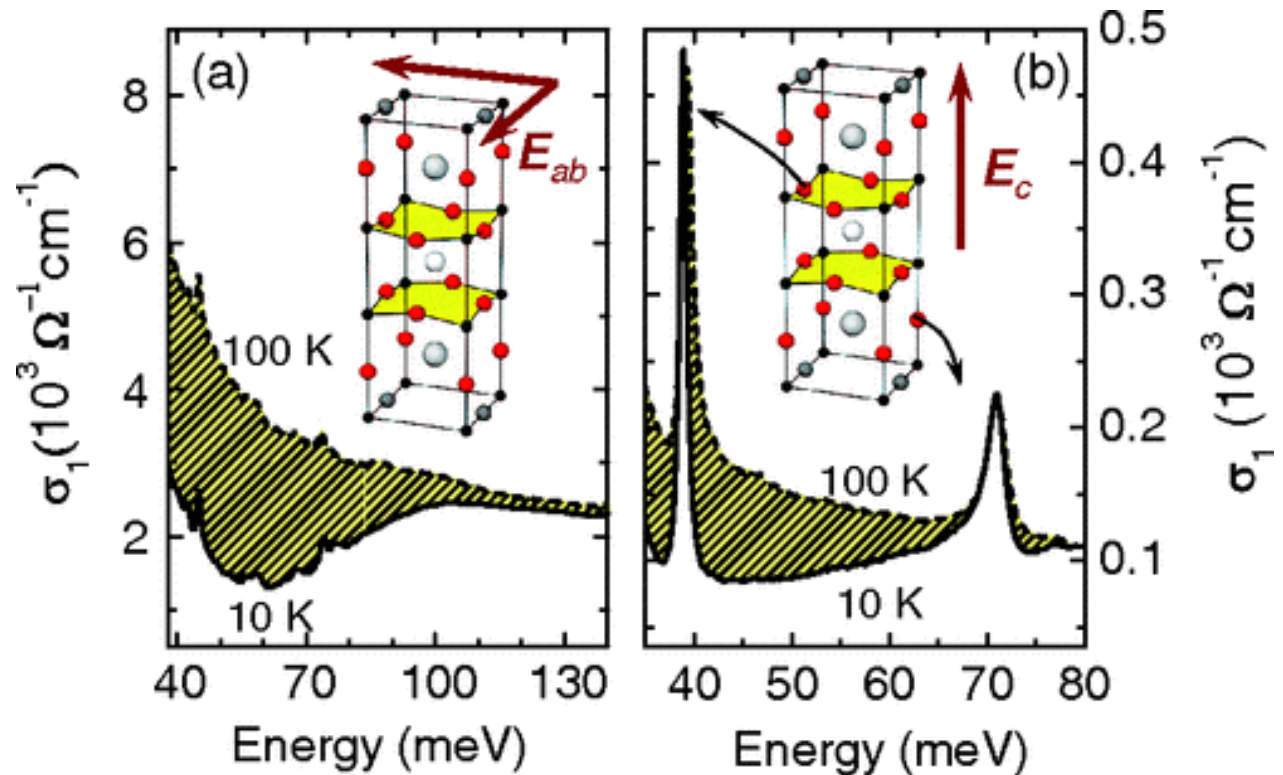
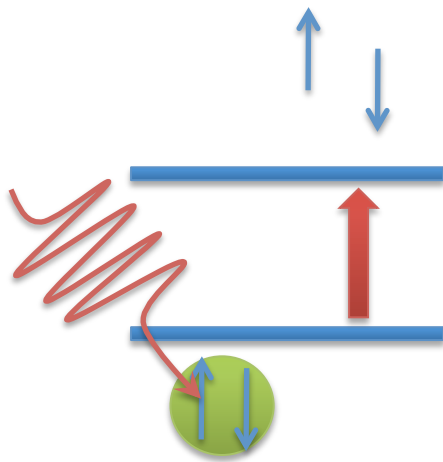
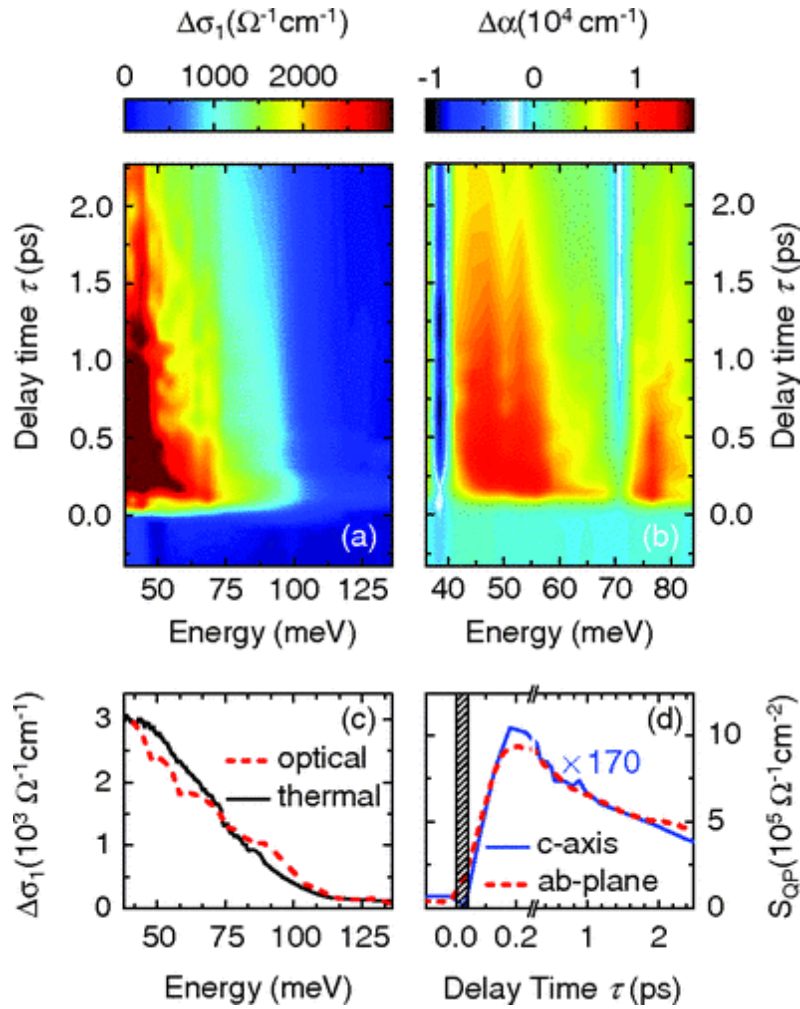


Photo-doping the SC gap

- Optical gap in YBCO

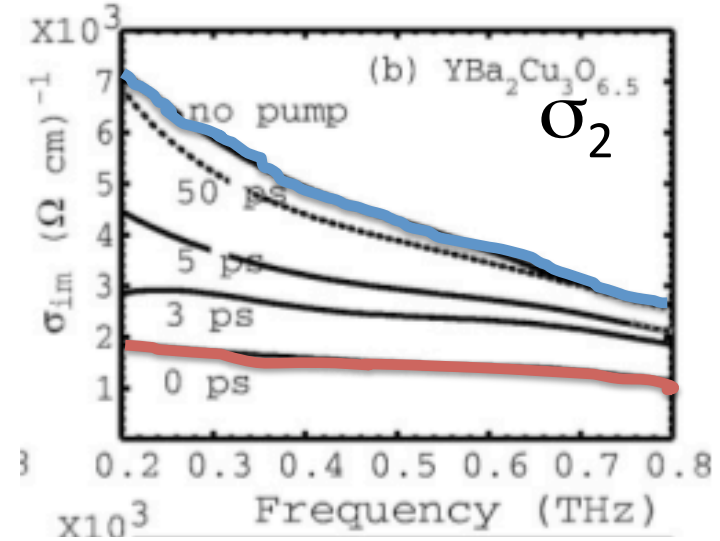


Spectroscopy of the gaps



A. Pashkin et al. PRL (2010)

- Full optical response

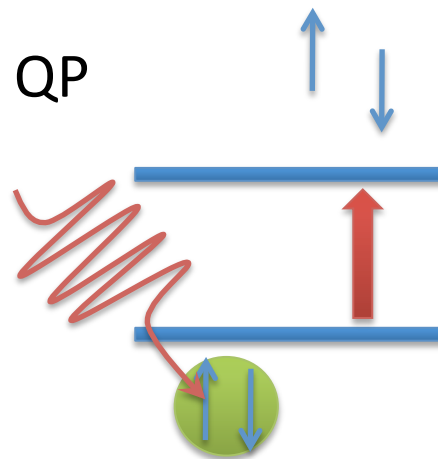


R. Averitt PRB (2001)

Time resolved pump-probe spectroscopy

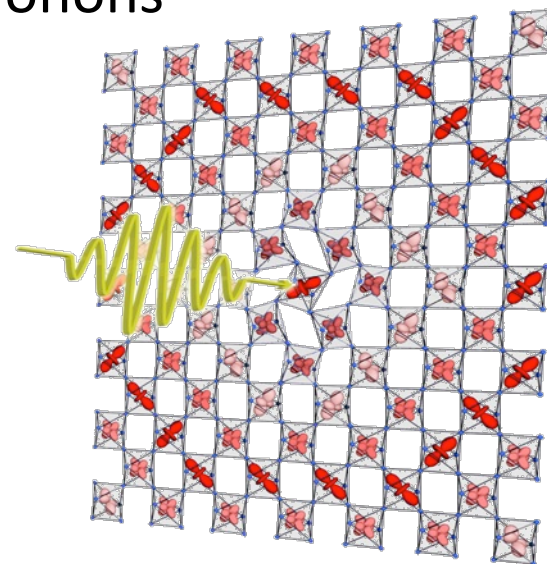
- Spectroscopy of non equilibrium QP

1.5 eV ~ 20,000 K



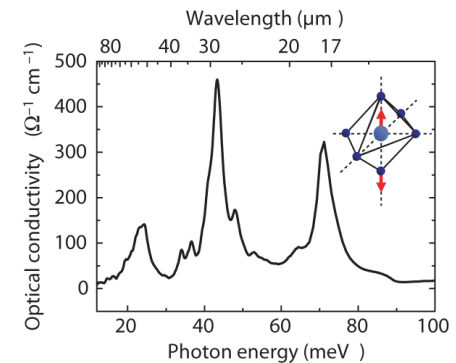
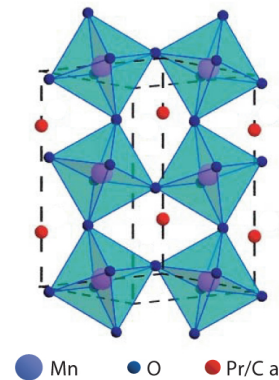
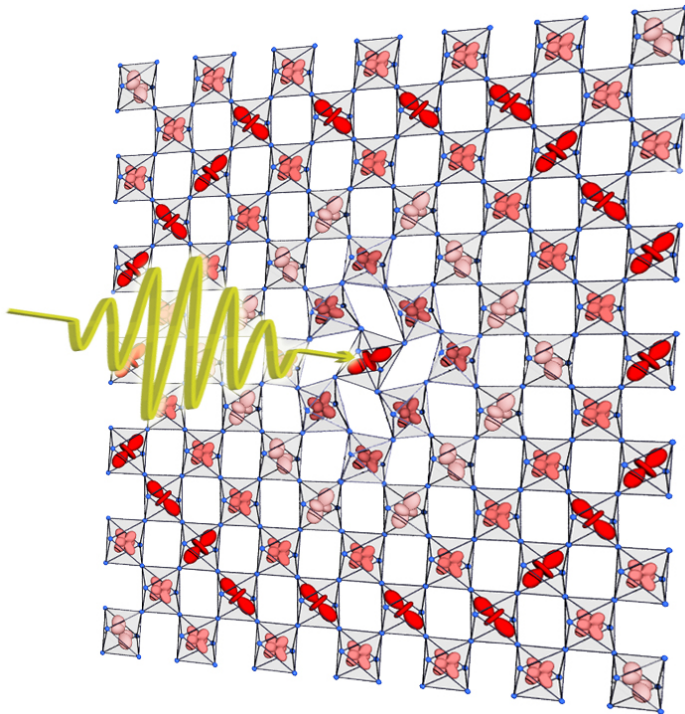
- THz control: Collective excitations, phonons

1 THz ~ 4 meV ~ 50 K



Selective lattice deformation

Resonant phonon excitation

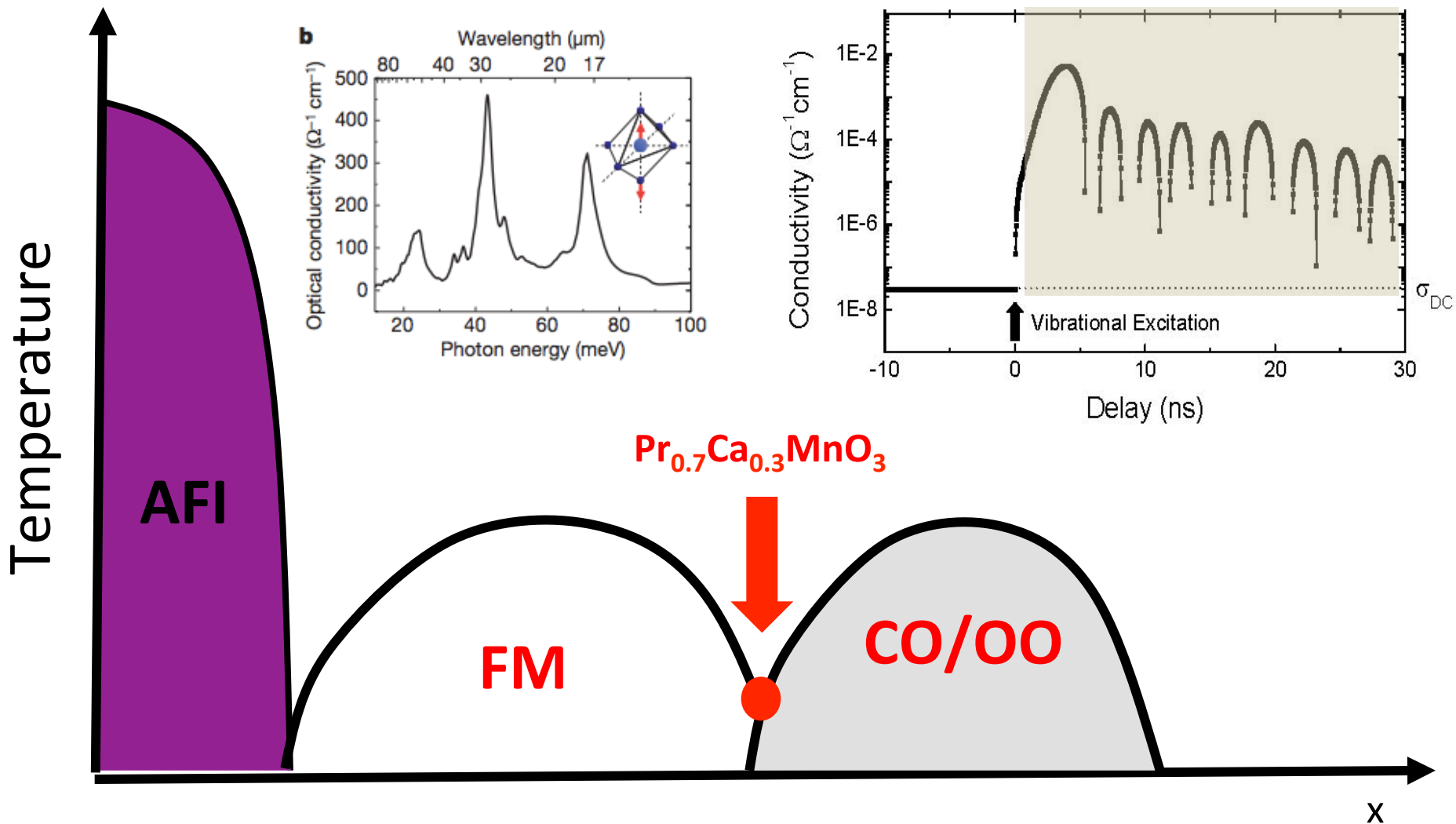


17 μm

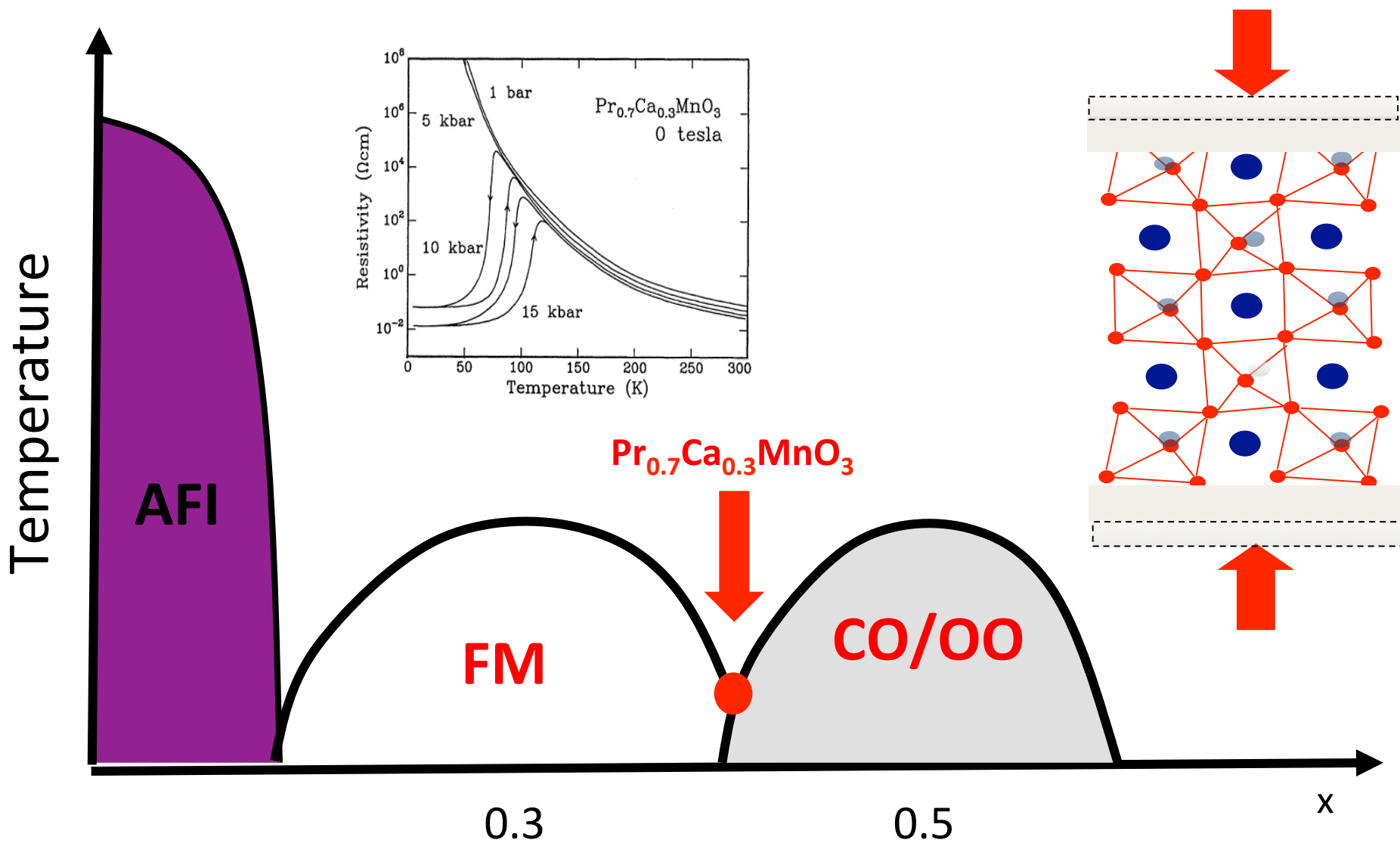
17 THz

70 meV

Non-thermal switching in $\text{Pr}_{0.7}\text{Ca}_{0.3}\text{MnO}_3$

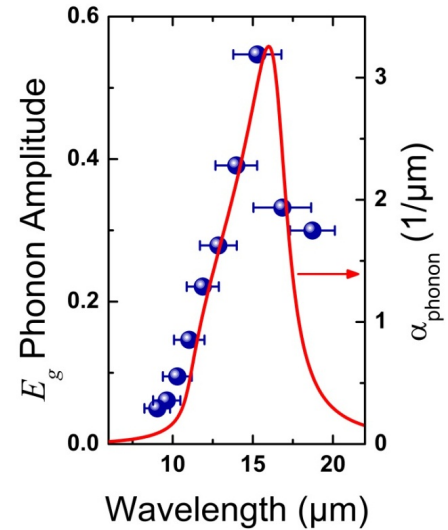
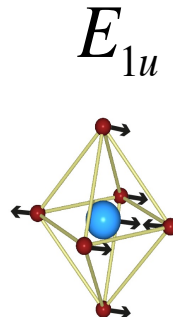
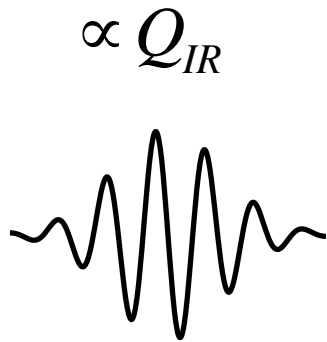


Pressure tuning of a metal insulator transition in PCMO

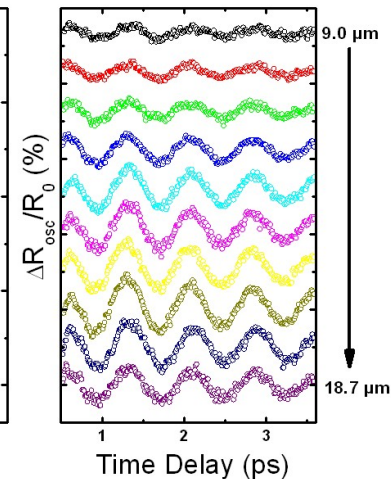
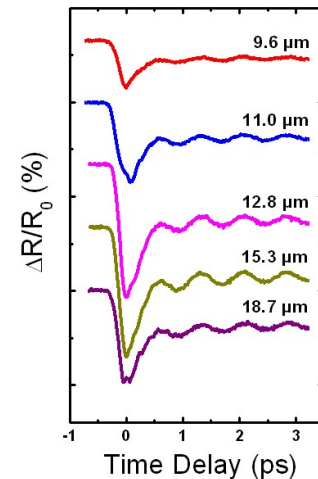
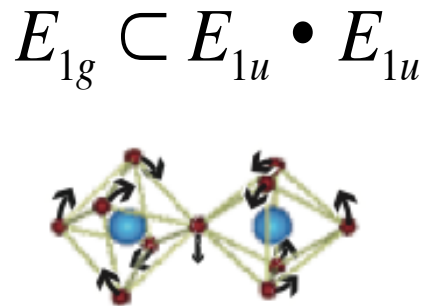
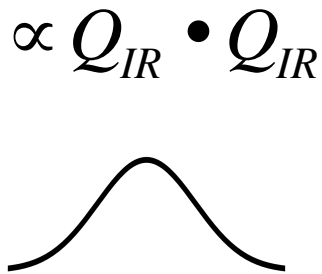


How does an optical pulse exert pressure ?

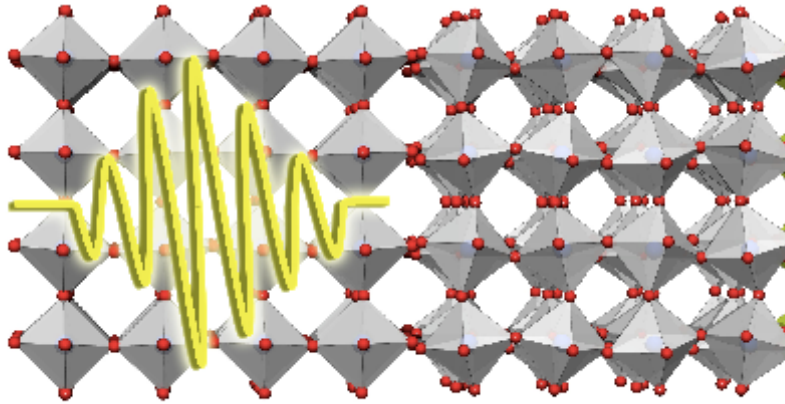
- resonant excitation at the 15 μm Mn-O mode



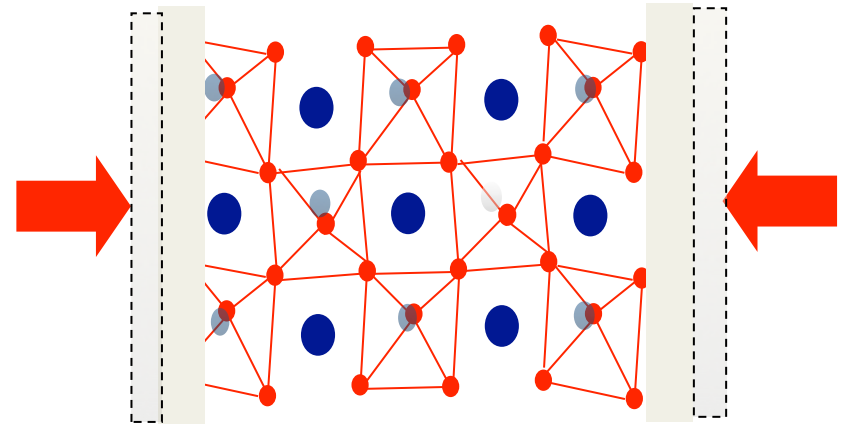
- Ionic Raman Scattering: Rectification into 1.2 THz mode



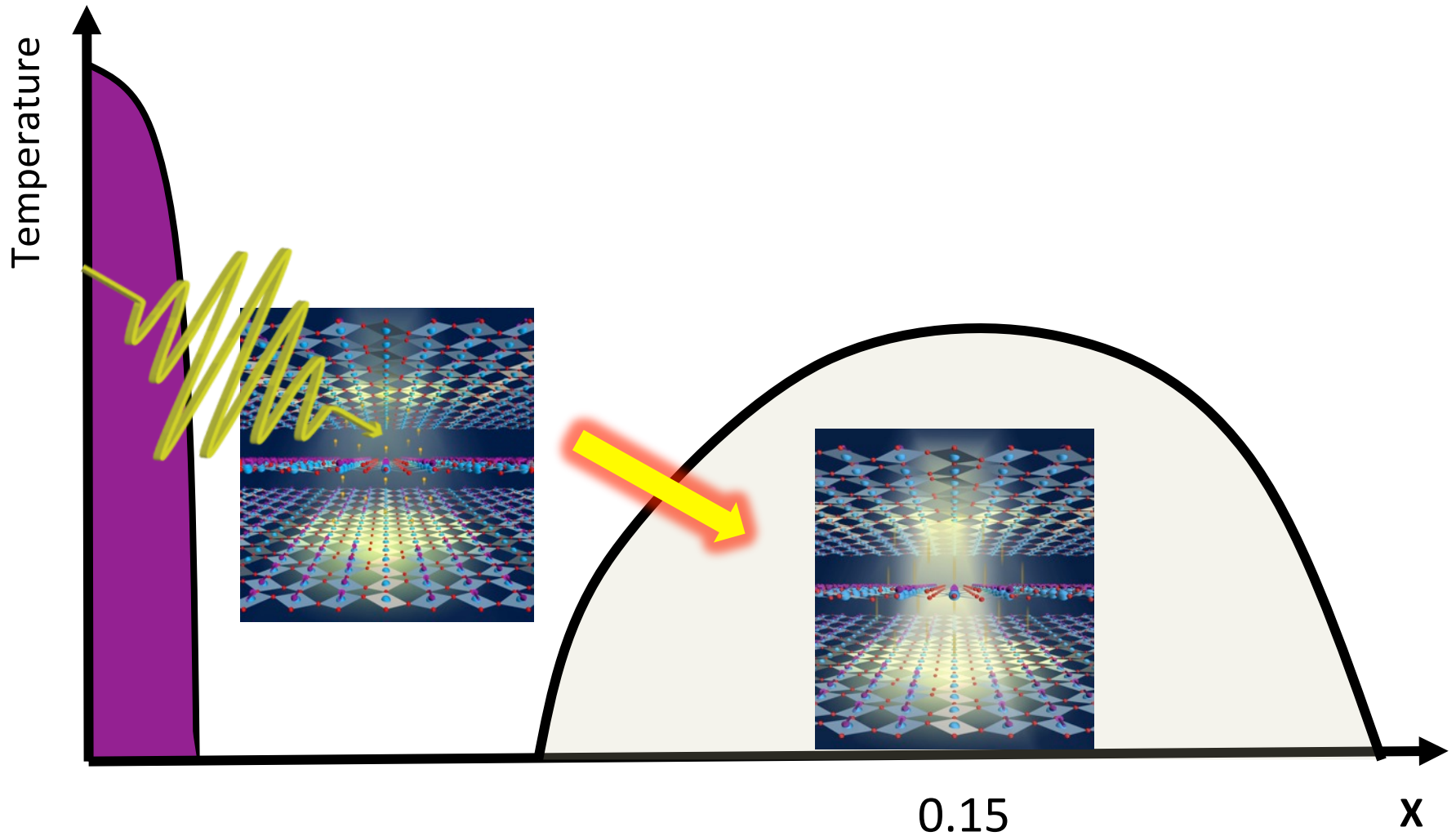
Mode selective pressure



- M. Rini et al., Nature (2007)
- R. Tobey et al, PRL (2008)
- M. Foerst et al., Nature Physics (2011)
- M. Foerst et al., PRB (2011)
- A. Caviglia et al, PRL (2011)



Can we control SC with light?

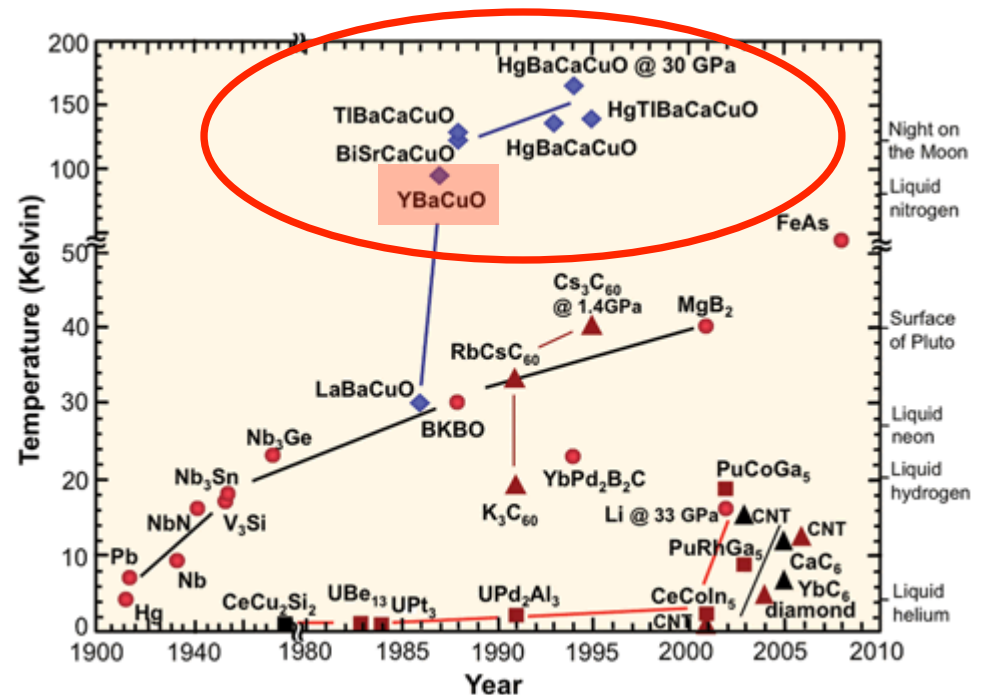


History of higher T_c

Increase of T_c :
Materials discovery

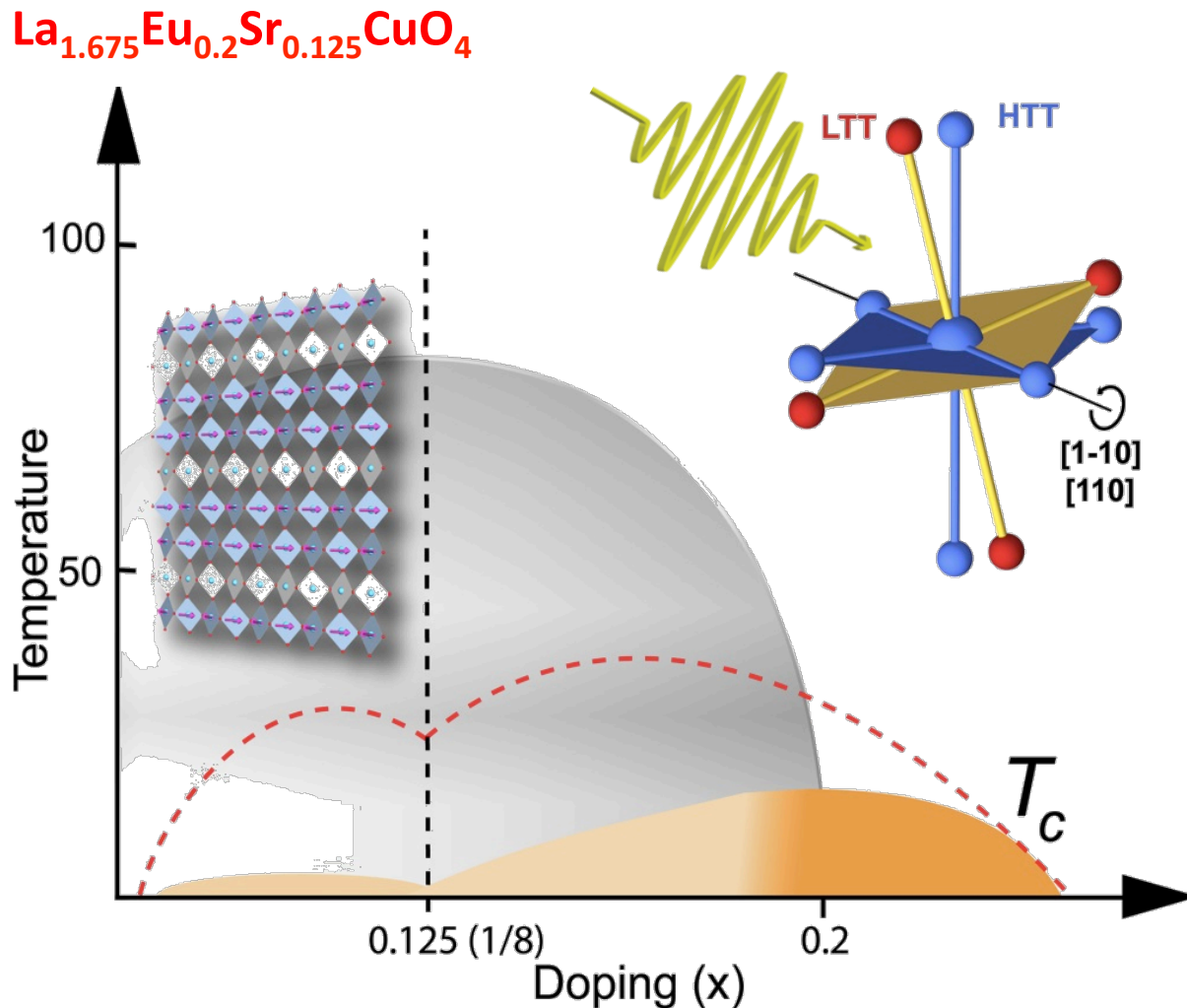
Layered Copper Oxides

$$T_{c,max} \ll T_{room}$$

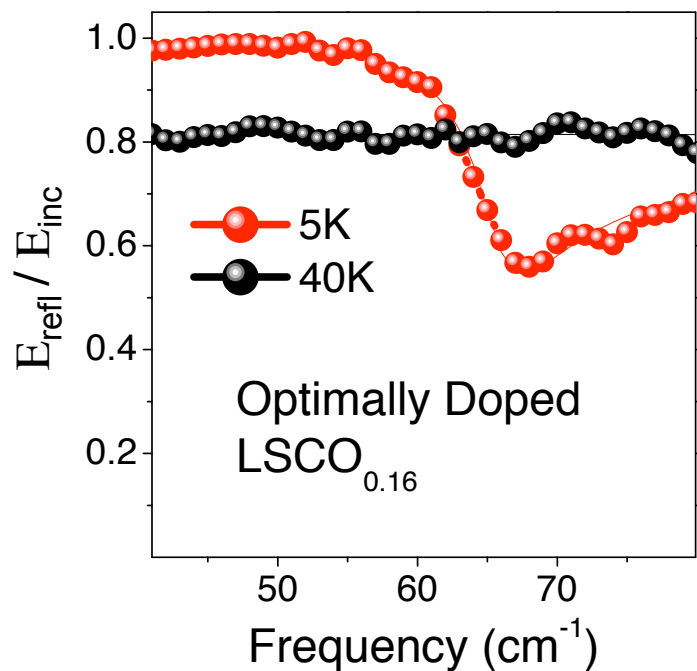


Can light induced phase transitions help us reaching room-temperature?

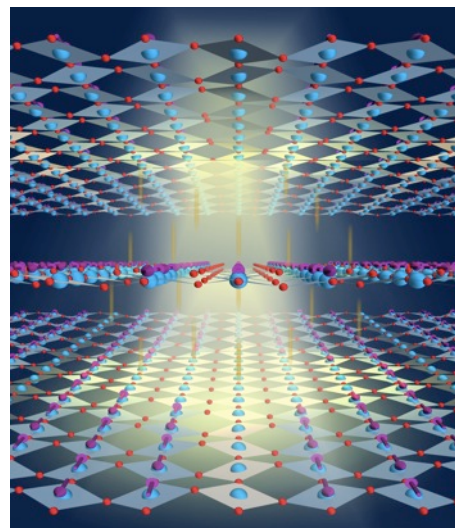
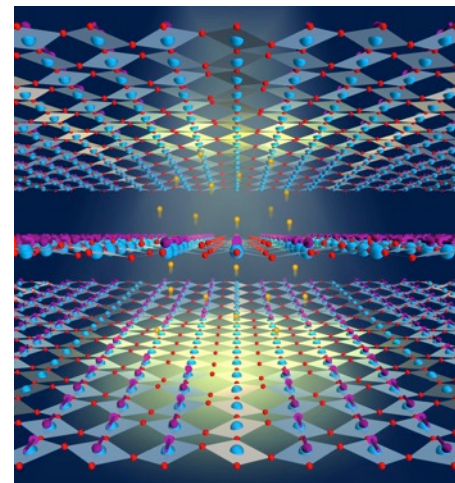
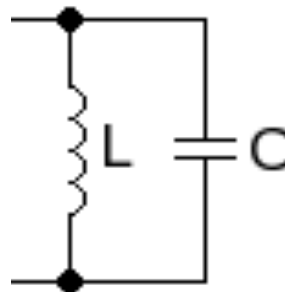
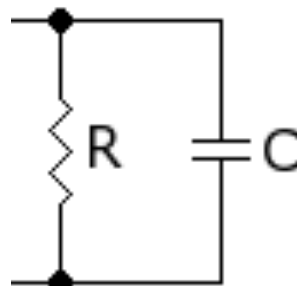
Stripe phase: Competing order to SC



How do we recognize a superconductor?

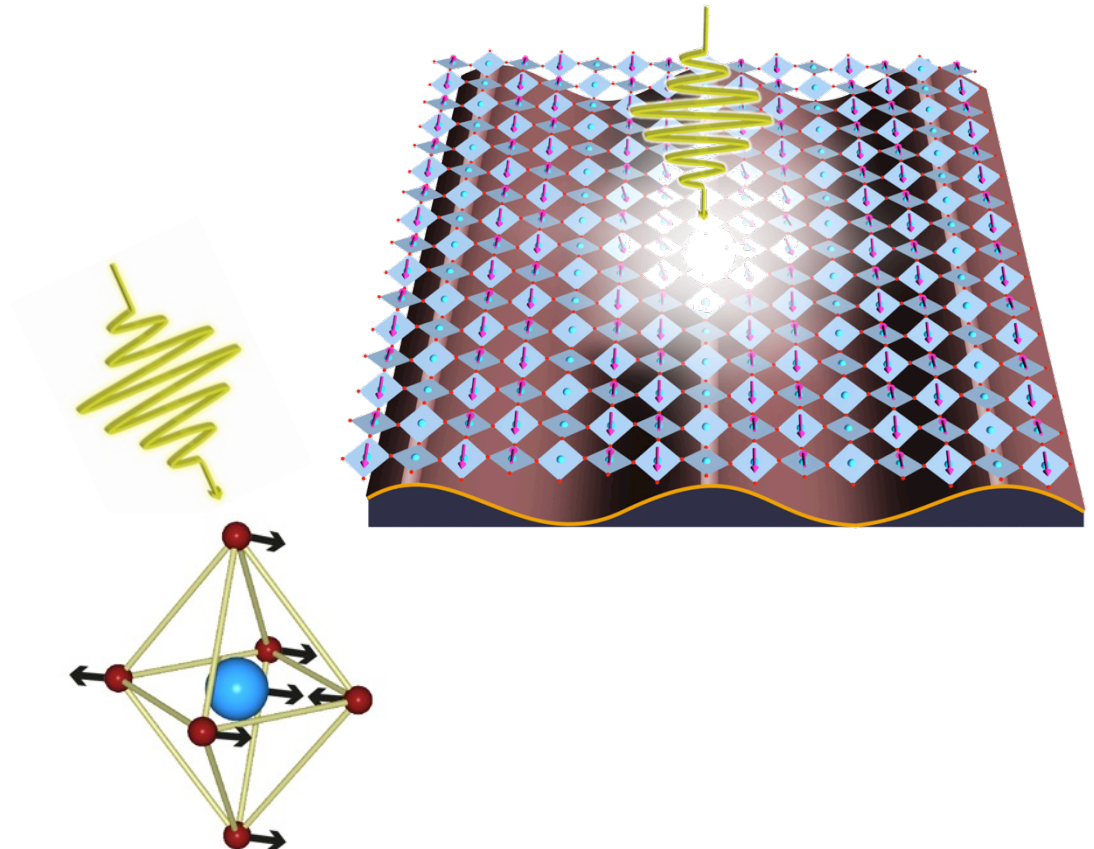
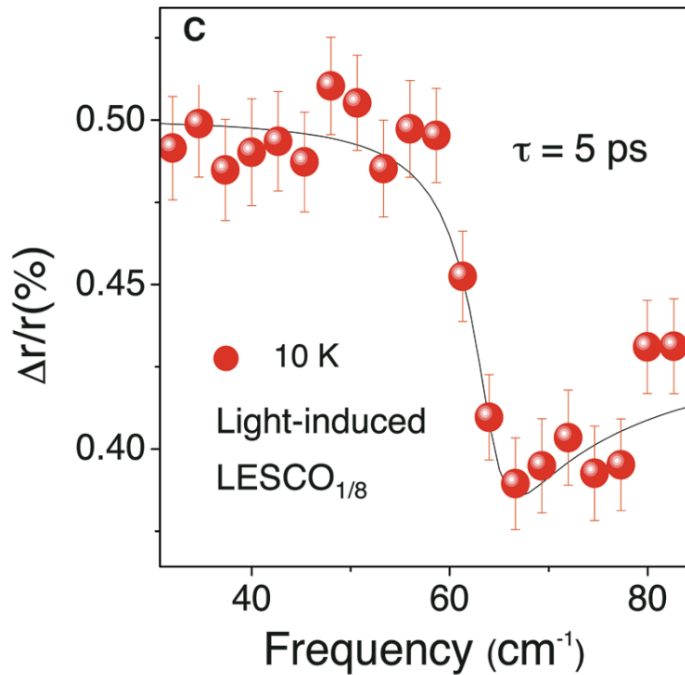


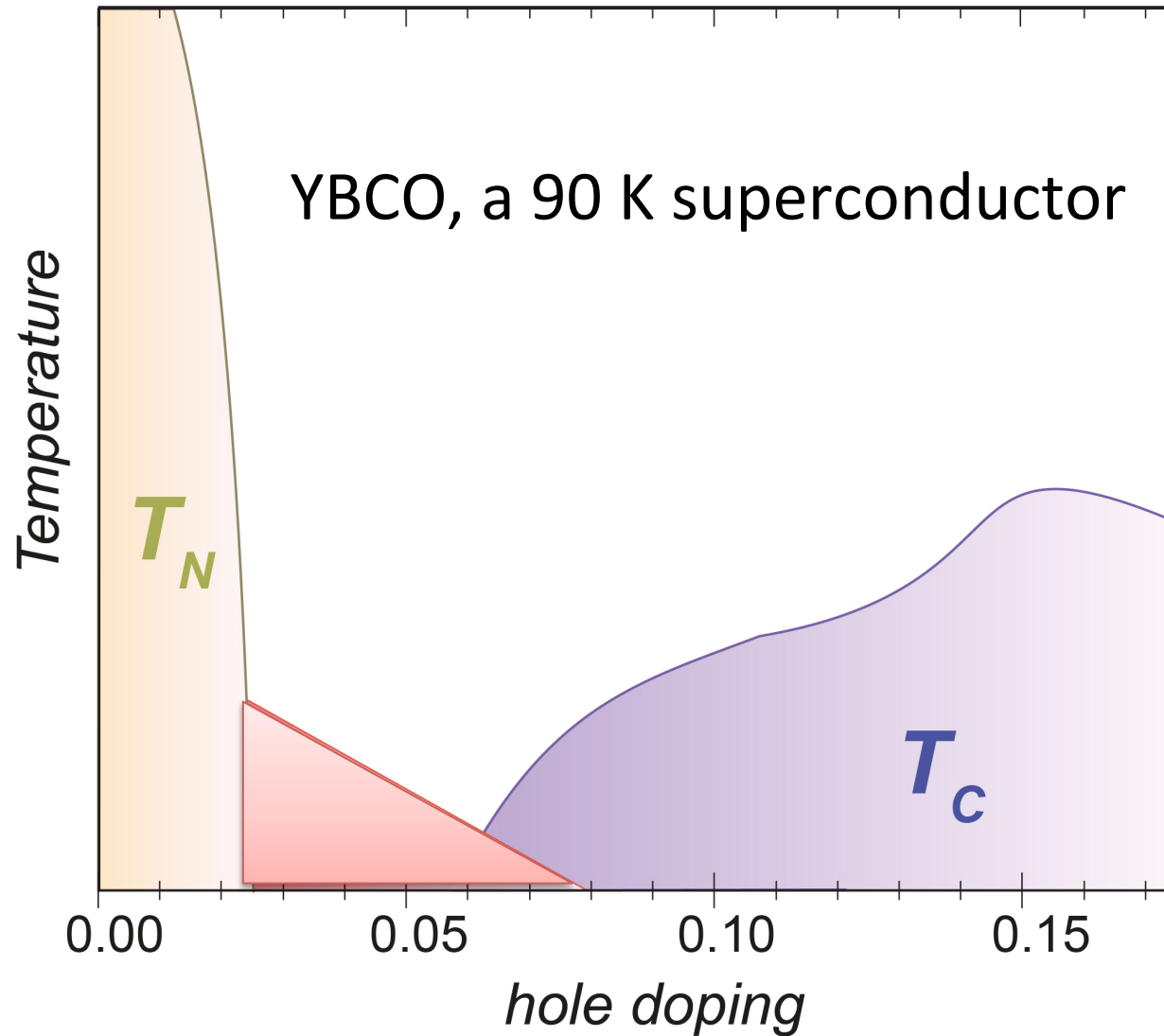
Josephson Plasma edge
of a superconductor



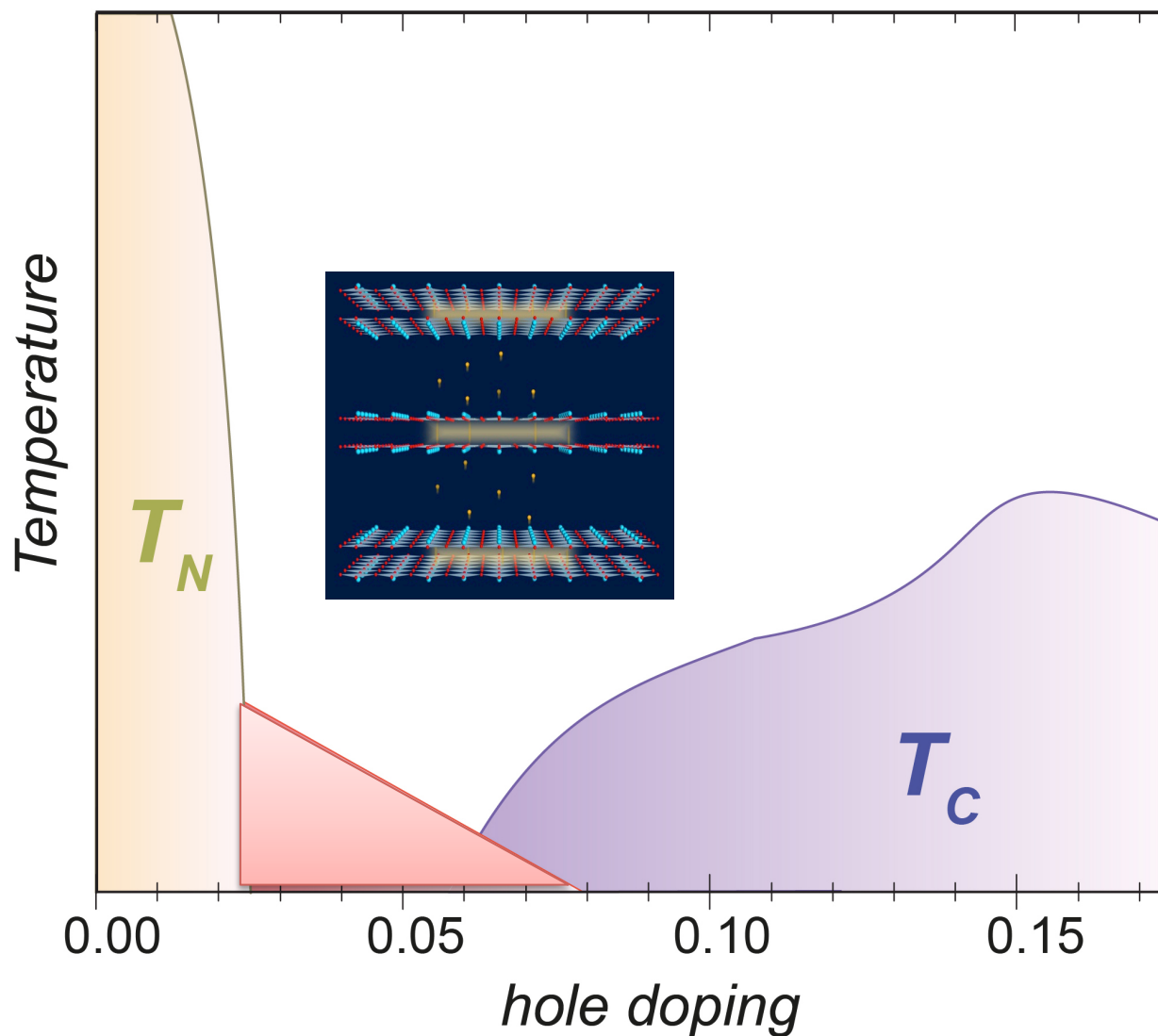
Light-induced superconductivity

Photo-induced JPR from the stripe phase





Is there superconducting coherence above T_c ?



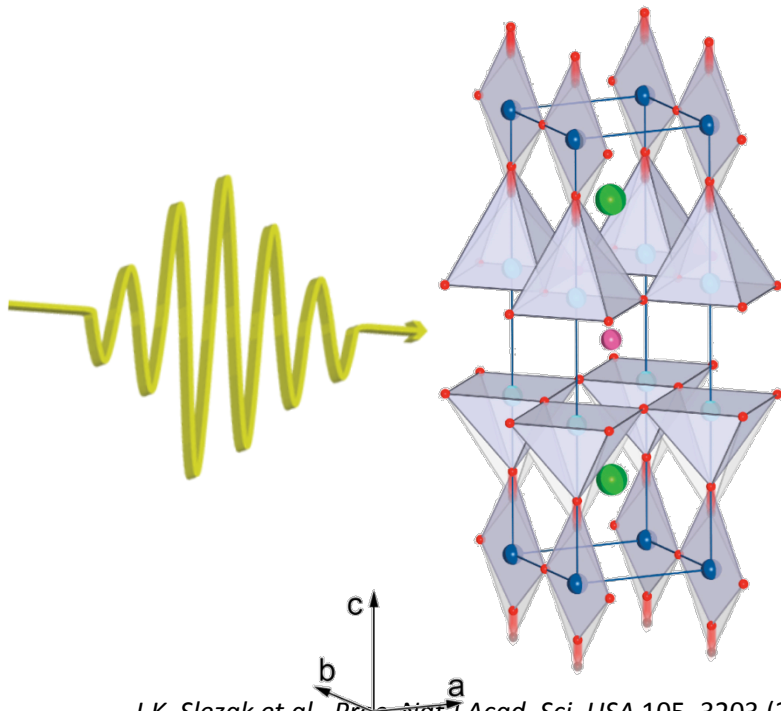
A. Dubroka et al., *Phys. Rev. Lett.* 107, 047006 (2011)

Can we modulate coherence in YBCO?

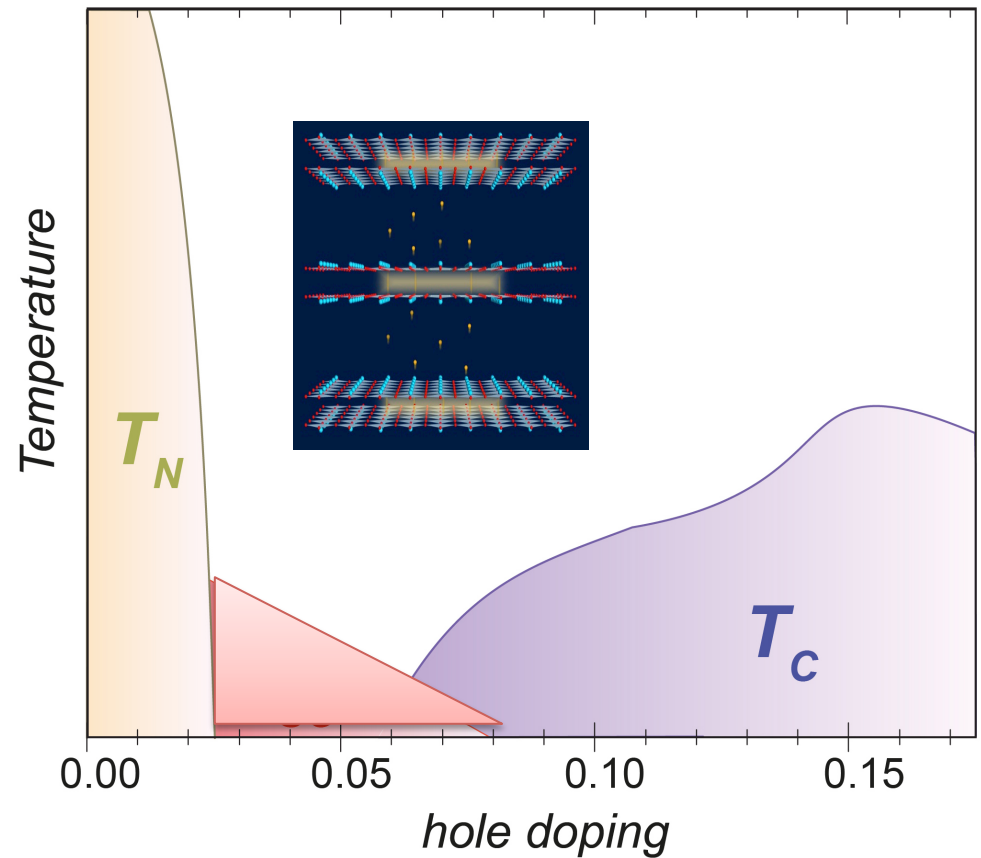
Apical oxygen in cuprates: stretching mode at $15.5 \mu\text{m}$

Direct control of

- In-plane hopping
- Exchange interaction



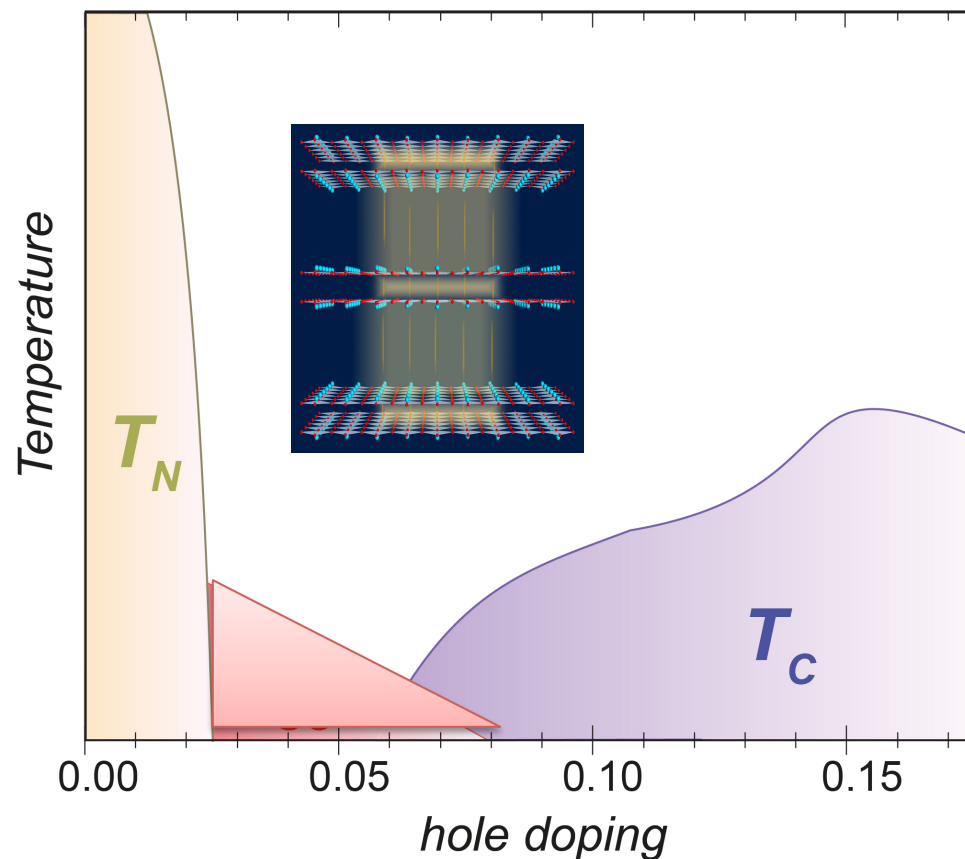
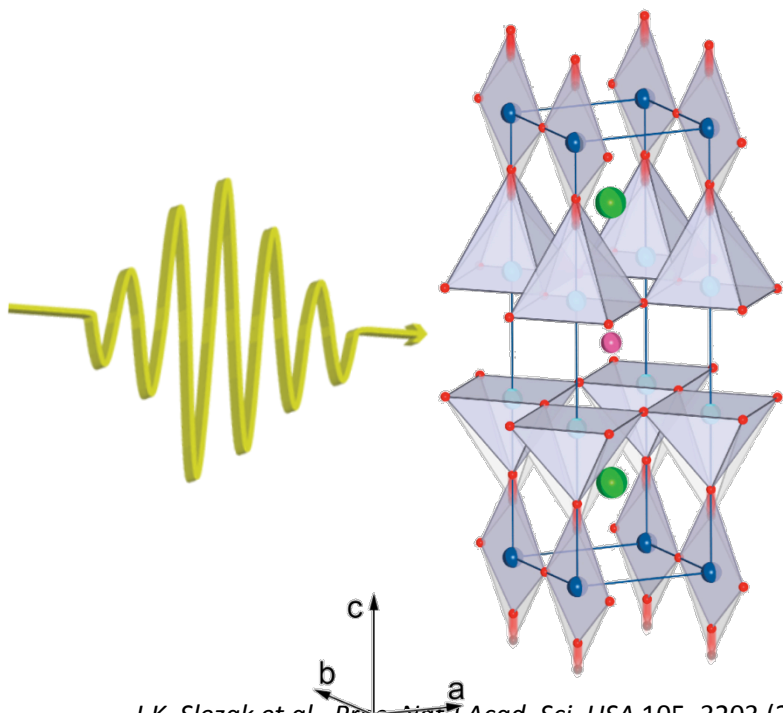
J.K. Slezak et al., Proc. Natl Acad. Sci. USA 105, 3203 (2008).



E. Pavarini et al. Phys. Rev. Lett. 87, 047003 (2001).

Dynamically modulated state in YBCO

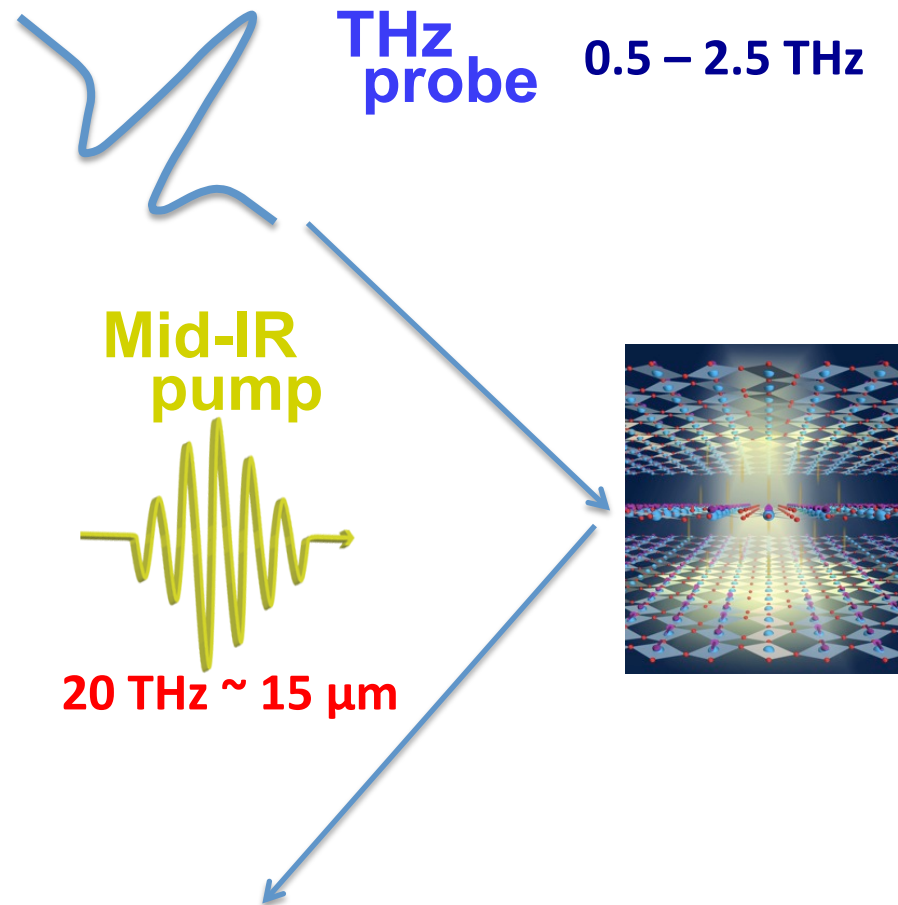
Is the modulated state more coherent ?



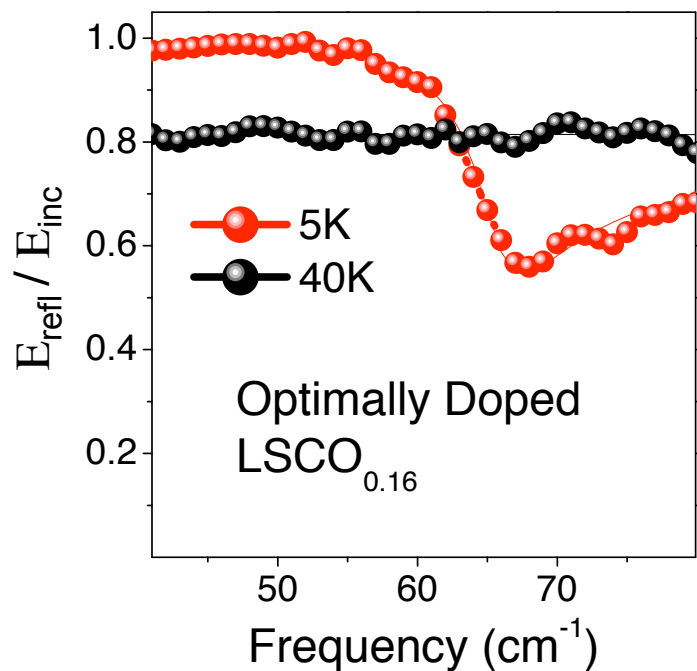
J.K. Slezak et al., *Proc. Natl Acad. Sci. USA* 105, 3203 (2008).

E. Pavarini et al. *Phys. Rev. Lett.* 87, 047003 (2001).

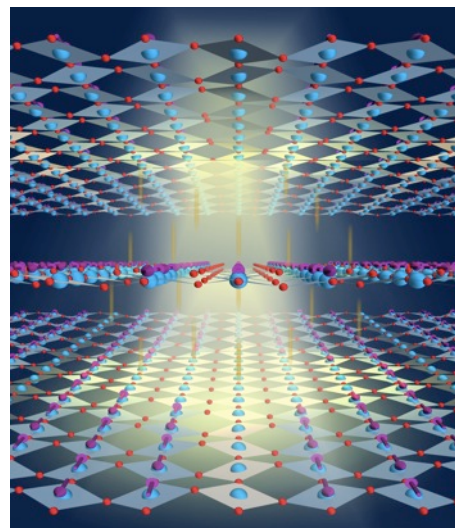
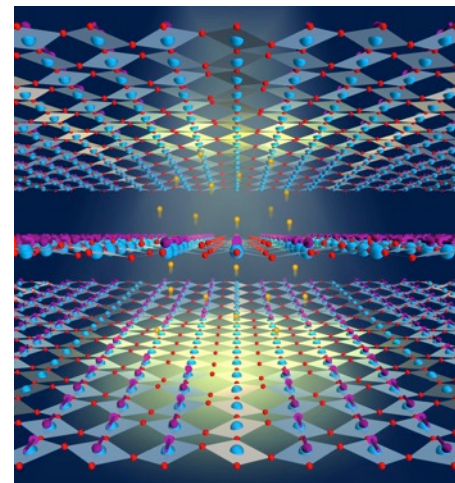
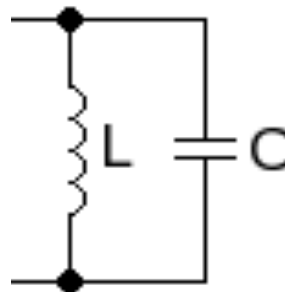
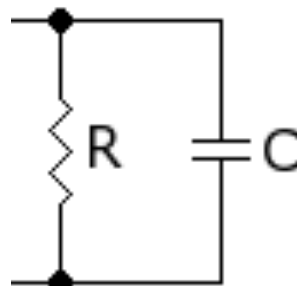
Vibrational modulation of a superconductor



How do we recognize a superconductor?



Josephson Plasma edge
of a superconductor

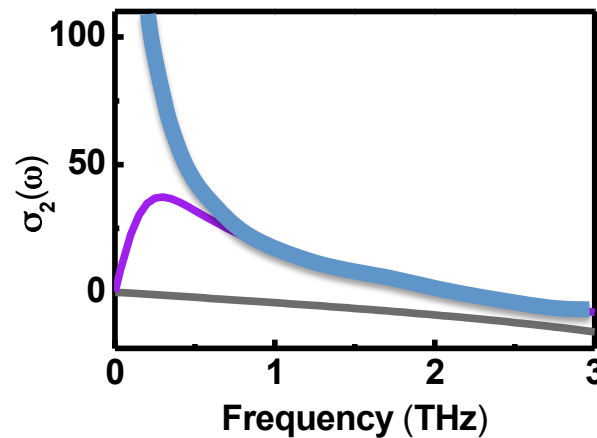
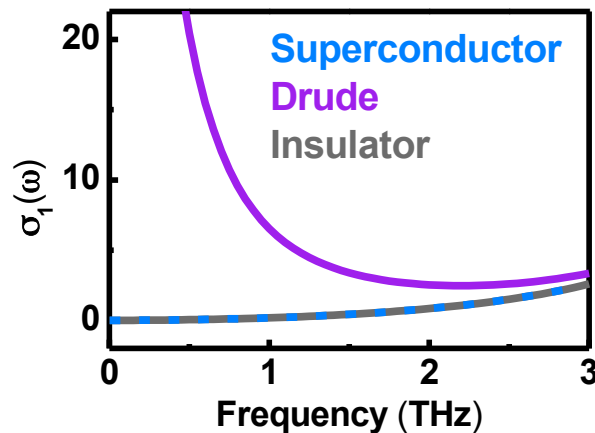


Signatures of superconductivity

- Zero resistivity $\rho_{\text{DC}}=0$
- Meissner effect

Signatures of superconductivity

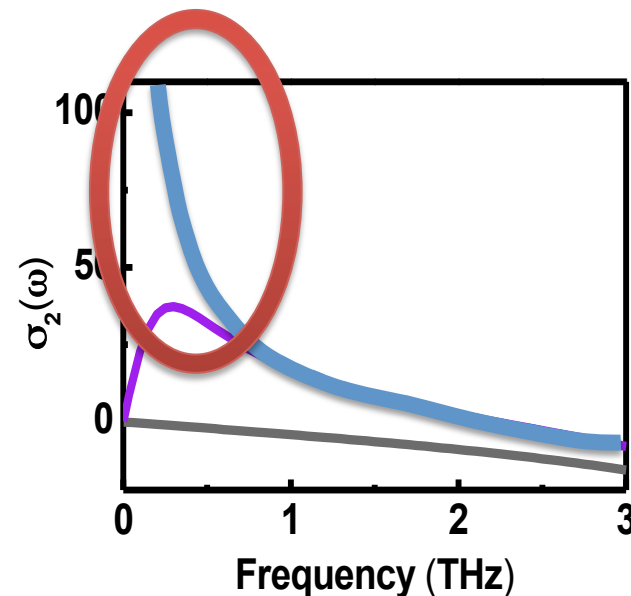
- Zero resistivity $\rho_{DC}=0$
- Meissner effect
- Fingerprints in the ac-optical conductivity
 - Delta peak at $\omega = 0$
 - London $1/\omega$ in σ_2



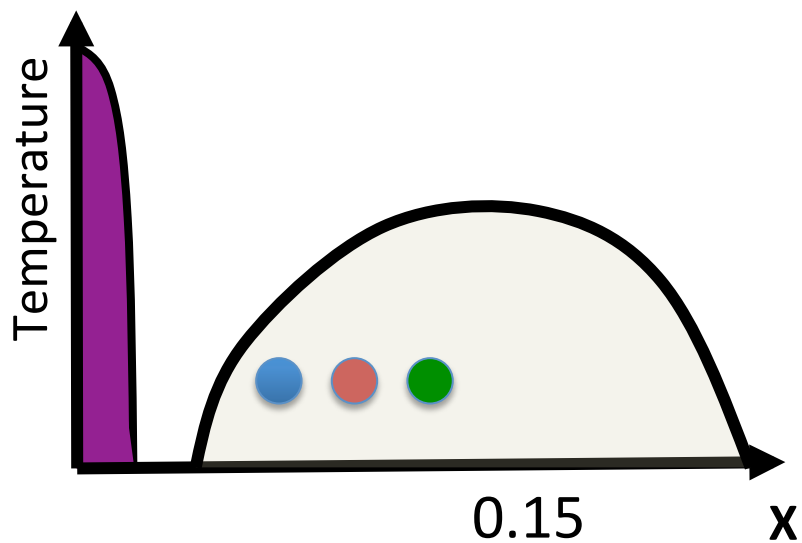
Signatures of superconductivity

- Zero resistivity $\rho_{DC}=0$
- Meissner effect
- Fingerprints in the ac-optical conductivity
 - Delta peak at $\omega = 0$
 - London $1/\omega$ in σ_2

→ Superfluid density: $\omega\sigma_2 \omega \rightarrow 0$



Vibrational modulation of a superconductor

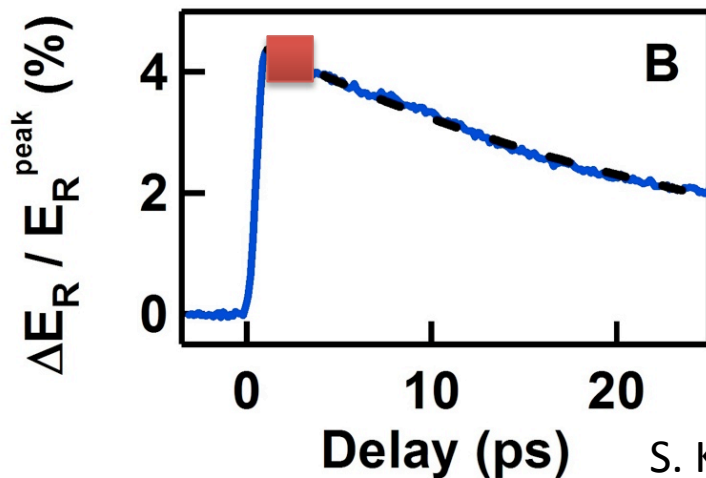
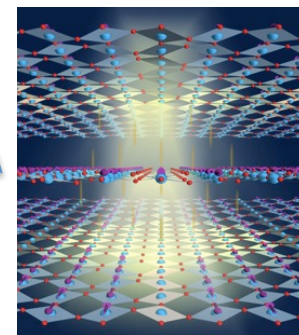


THz probe 0.5 – 2.5 THz

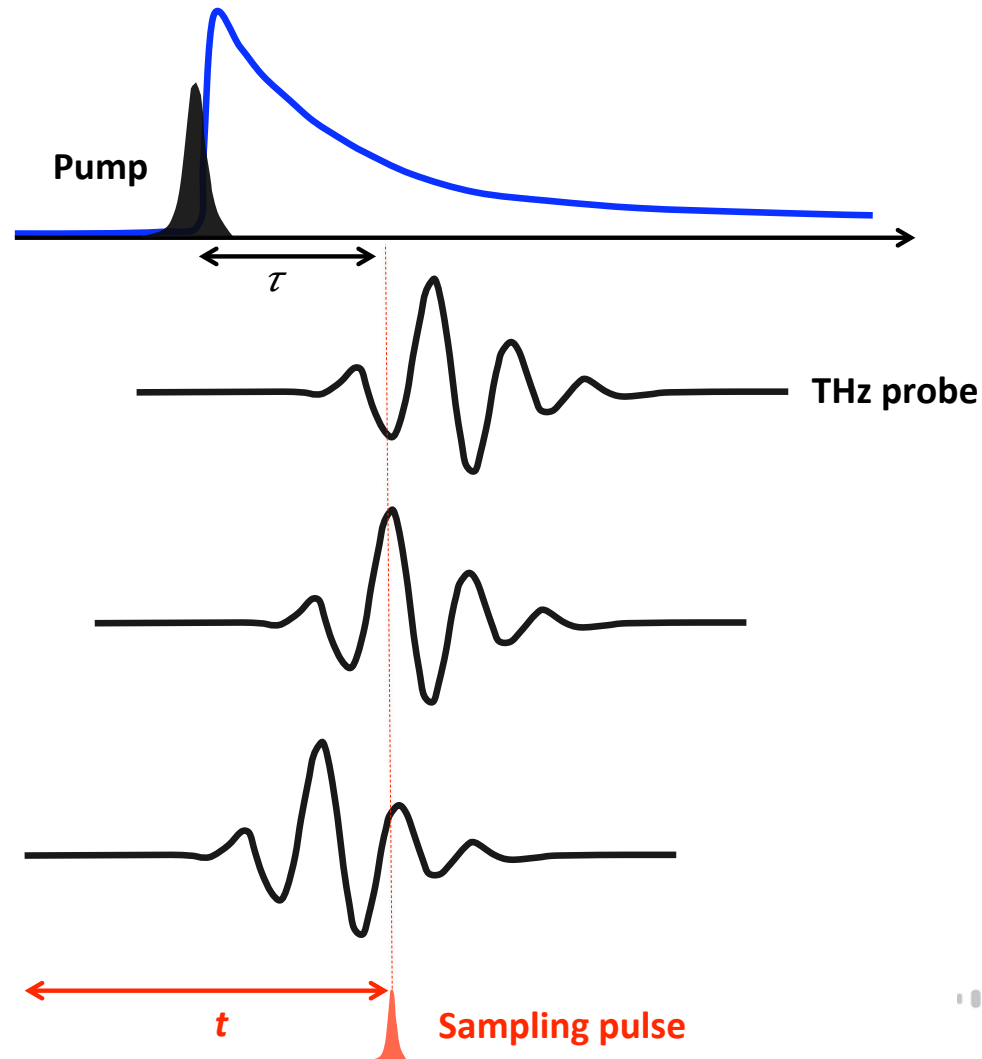
Mid-IR pump



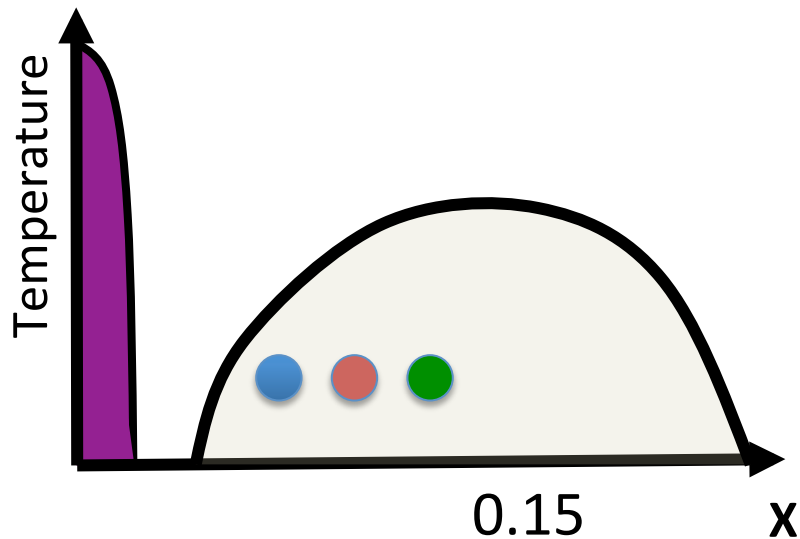
20 THz ~ 15 μm



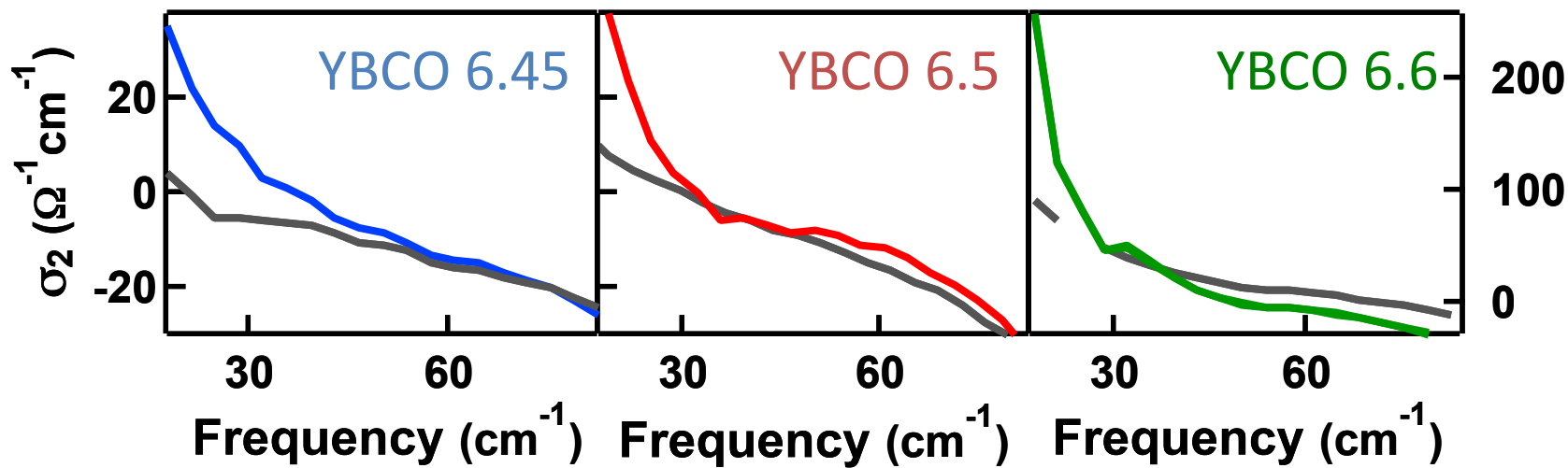
Electro optical sampling of the probe pulse



Enhanced coherence below T_c



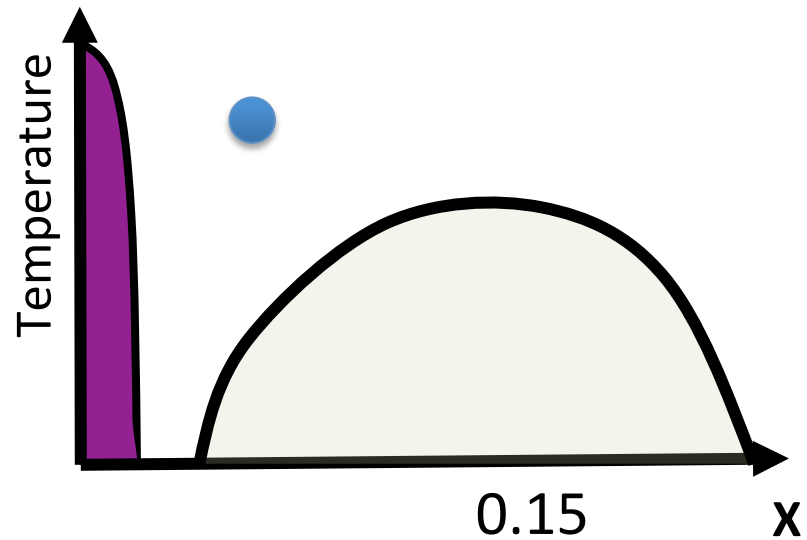
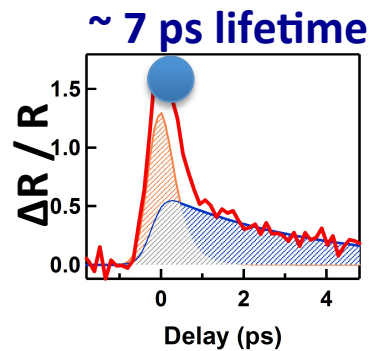
Increase of superfluid density:
enhanced low frequency JPR



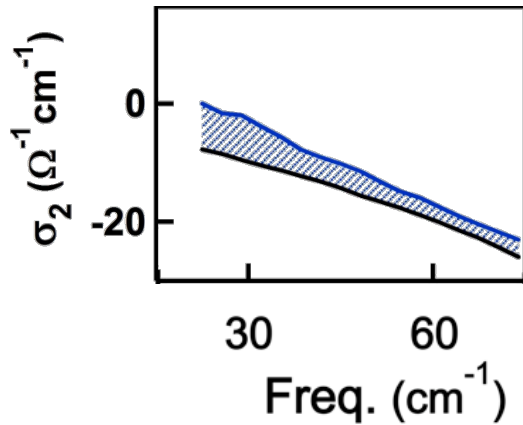
Dynamically modulated state above T_c

YBCO 6.45 ($T_c=35$ K)

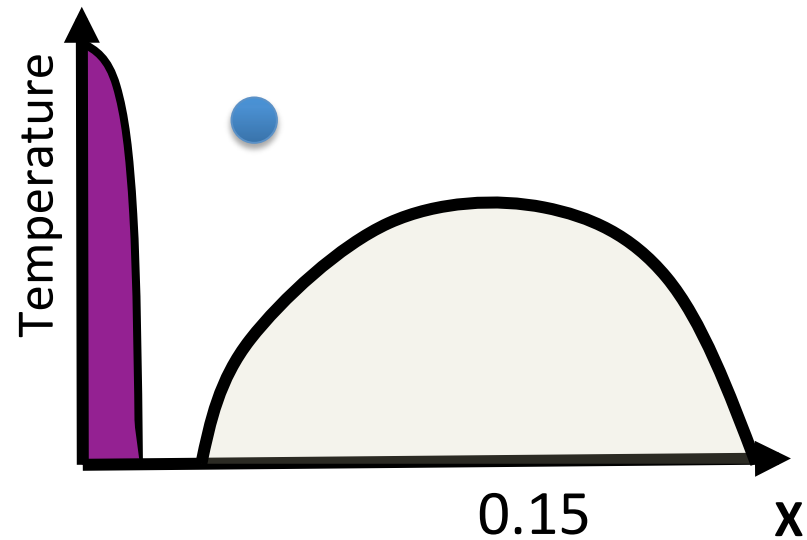
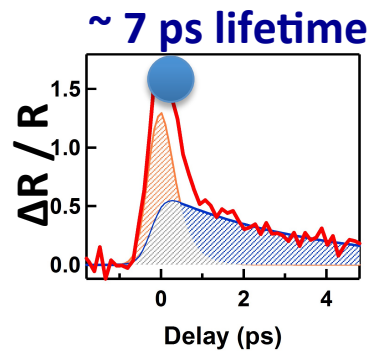
Base temperature 100 K



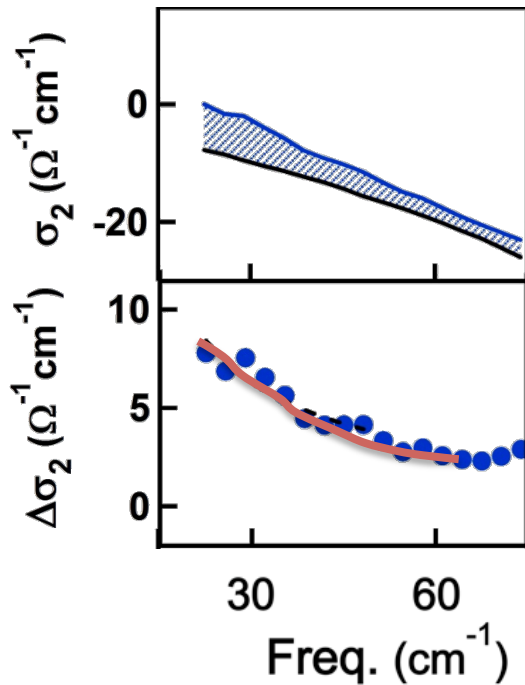
Optical response at 100 K YBCO 6.45 ($T_c=35$ K)



(1) Increased **inductive response**



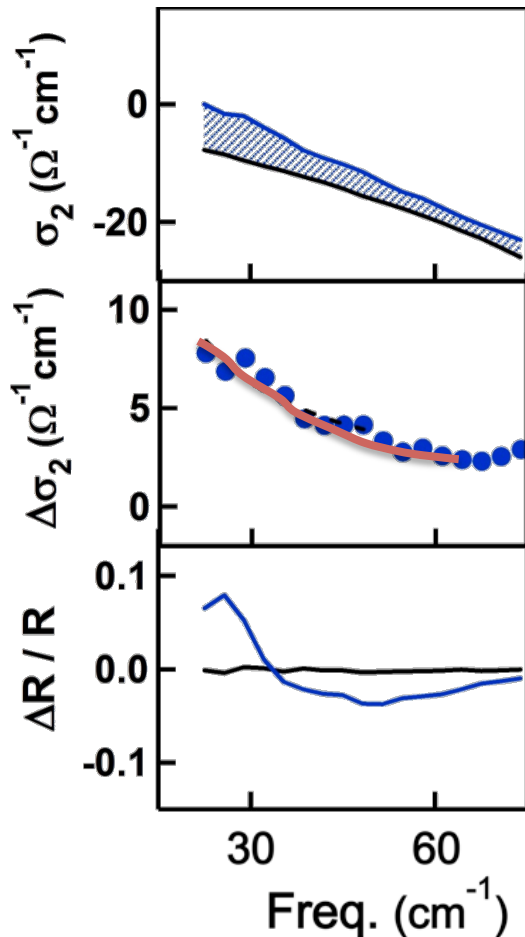
Optical response at 100 K YBCO 6.45 ($T_c=35$ K)



(1) Increased **inductive response**

(2) **1/ω** response in the σ_2 changes

Optical response at 100 K YBCO 6.45 ($T_c=35$ K)

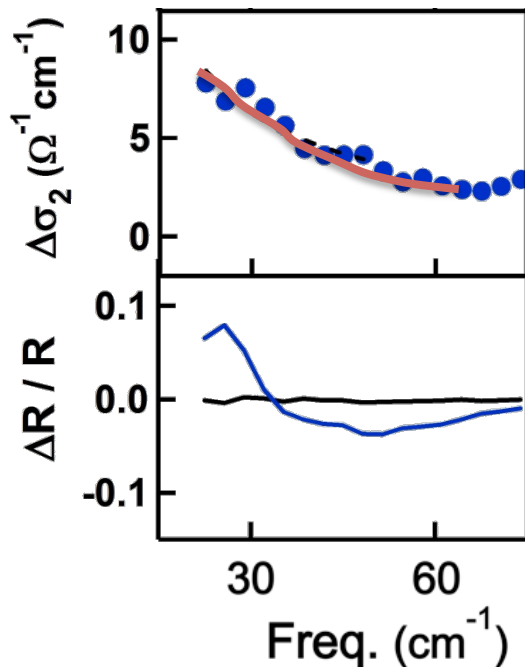


- (1) Increased **inductive response**
- (2) **1/ω** response in the σ_2 changes
- (3) Photoinduced **plasma edge**

Temperature dependence YBCO 6.45 ($T_c=35$ K)

Increase temperature:

100 K

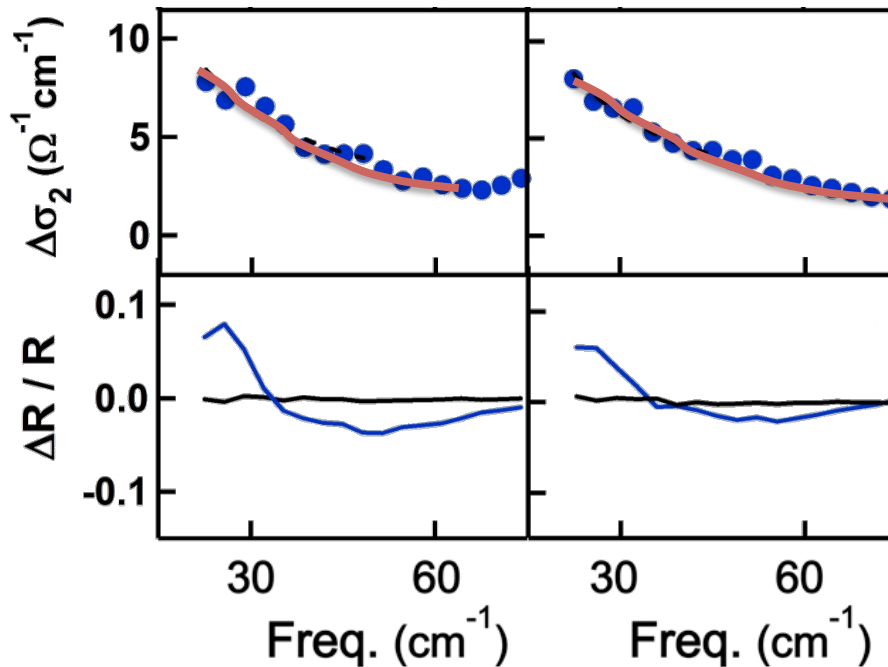


Temperature dependence YBCO 6.45 ($T_c=35$ K)

Increase temperature:

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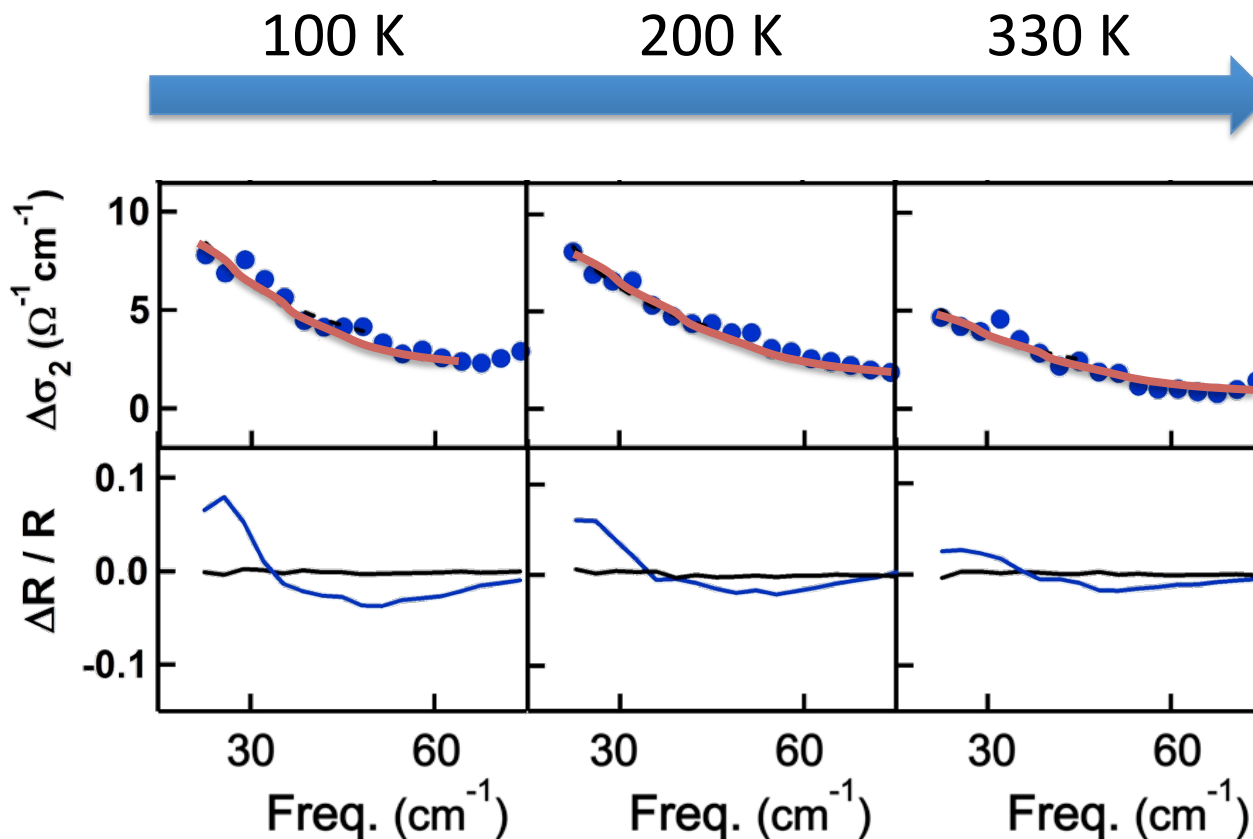
200 K



Temperature dependence YBCO 6.45 ($T_c=35$ K)

Increase temperature:

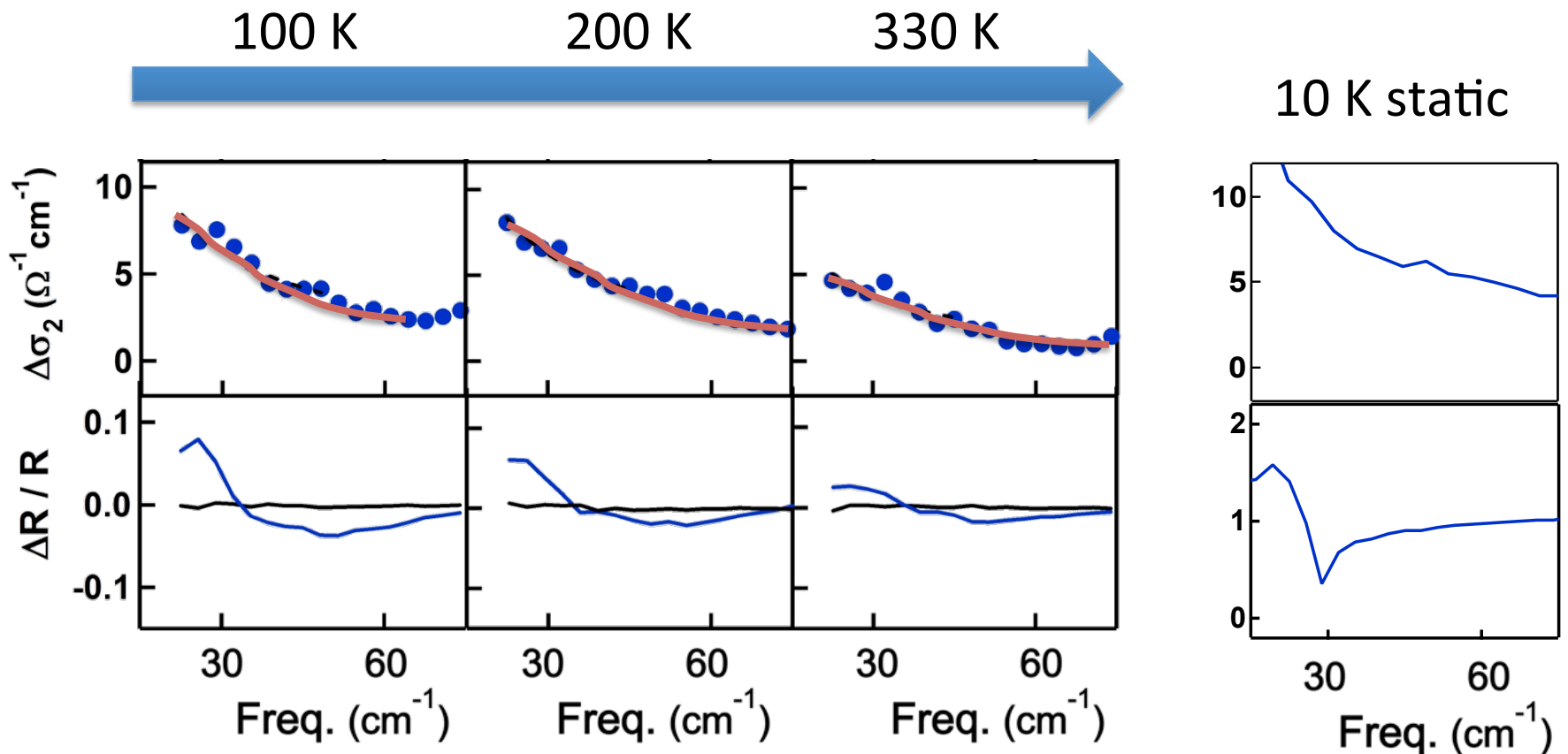
All effects persist!



Temperature dependence YBCO 6.45 ($T_c=35$ K)

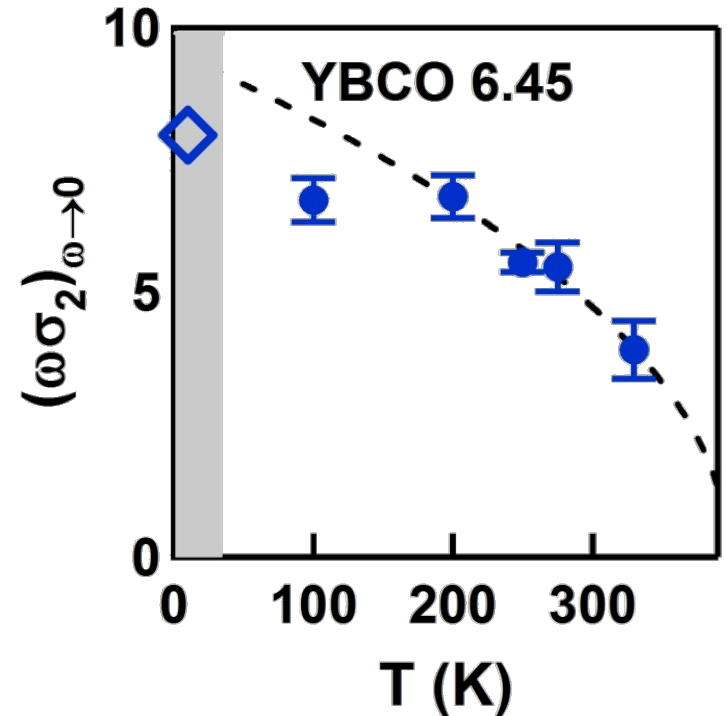
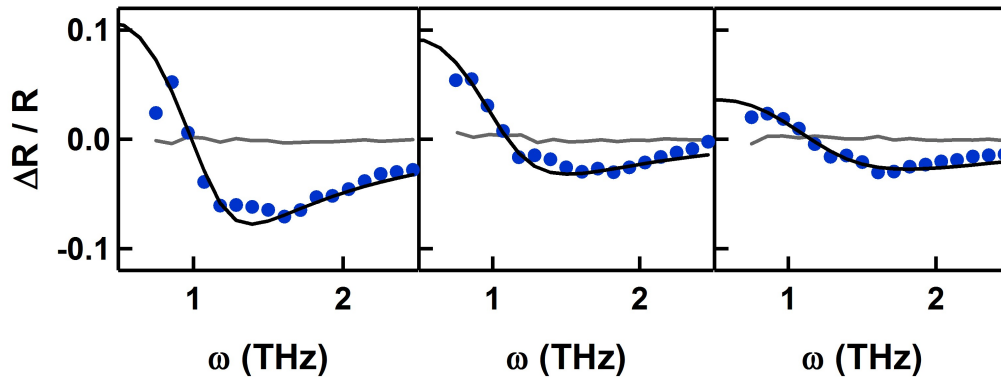
Increase temperature:

All effects persist!



Evaluating the “Superfluid” density

100 K 200 K 330 K



$$D^{SF} \propto (\omega\sigma_2)_{\omega \rightarrow 0}$$

$$D^{SF} \propto \sqrt{1 - \frac{T}{T'}}$$

Inhomogeneous nature of the SC state

Bruggeman Effective Medium

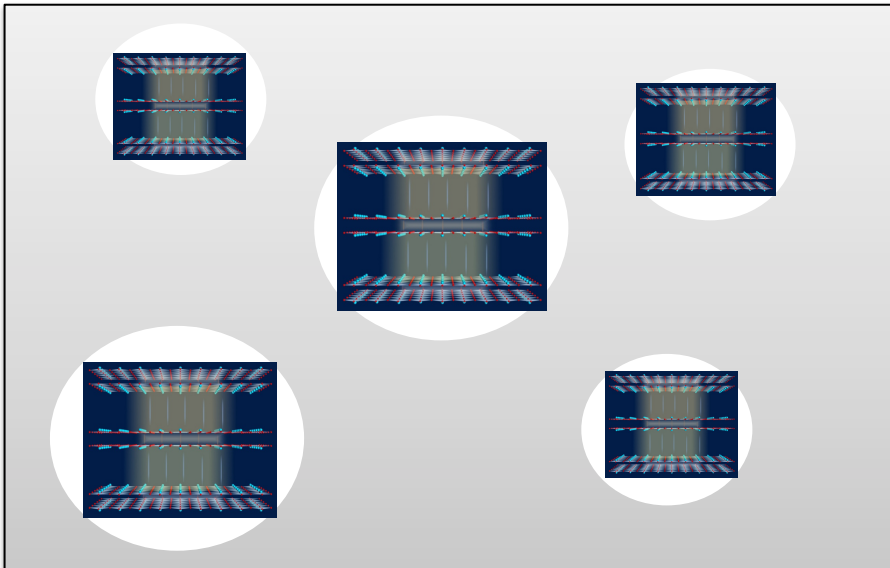
$$f \frac{\tilde{\epsilon}_S(\omega) - \tilde{\epsilon}_E(\omega)}{\tilde{\epsilon}_S(\omega) + 2\tilde{\epsilon}_E(\omega)} + (1 - f) \frac{\tilde{\epsilon}_{NS}(\omega) - \tilde{\epsilon}_E(\omega)}{\tilde{\epsilon}_{NS}(\omega) + 2\tilde{\epsilon}_E(\omega)} = 0,$$

ϵ_E : Effective medium dielectric function

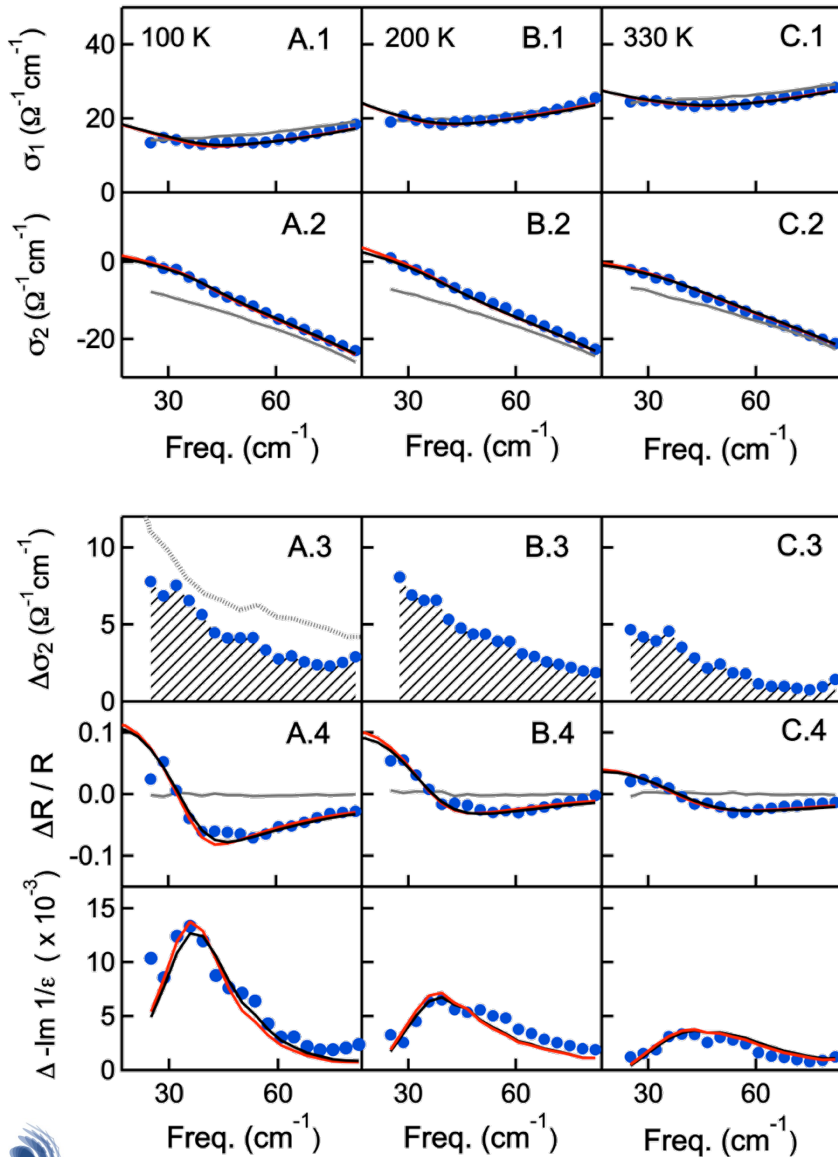
ϵ_S : Superconductor dielectric function
(Plasma frequency defines edge pos.)

ϵ_{NS} : Normal-state dielectric function

f : Superconducting volume fraction
(only free fit parameter)

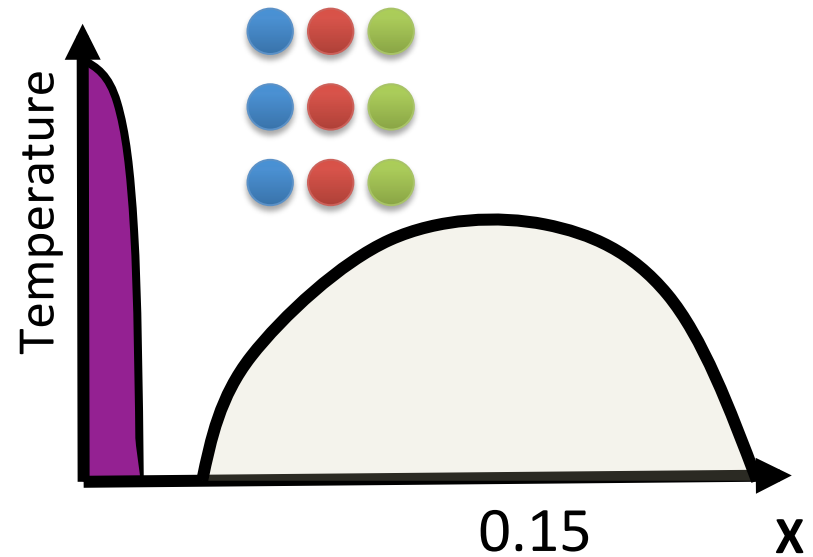


Effective medium fits to YBCO 6.45 ($T_c=35$ K)

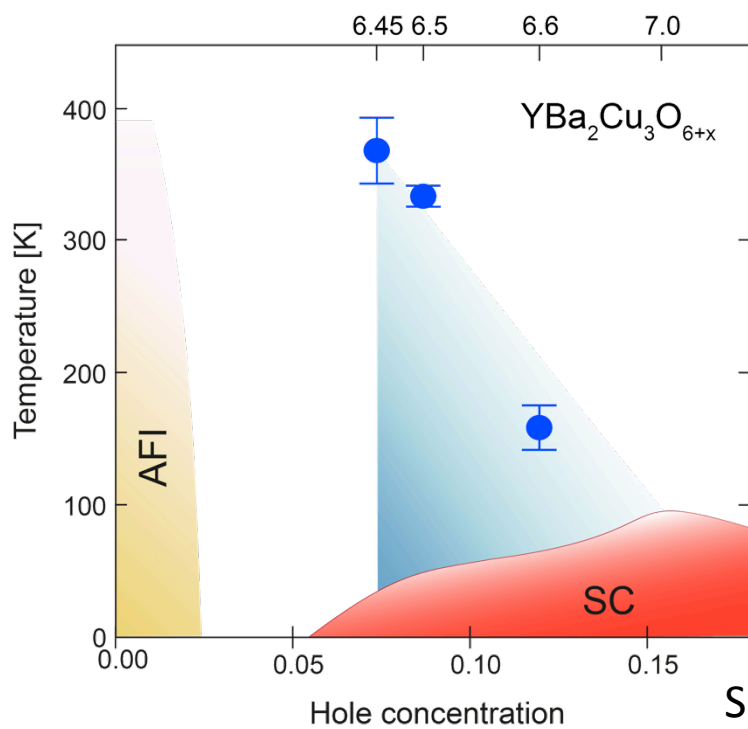
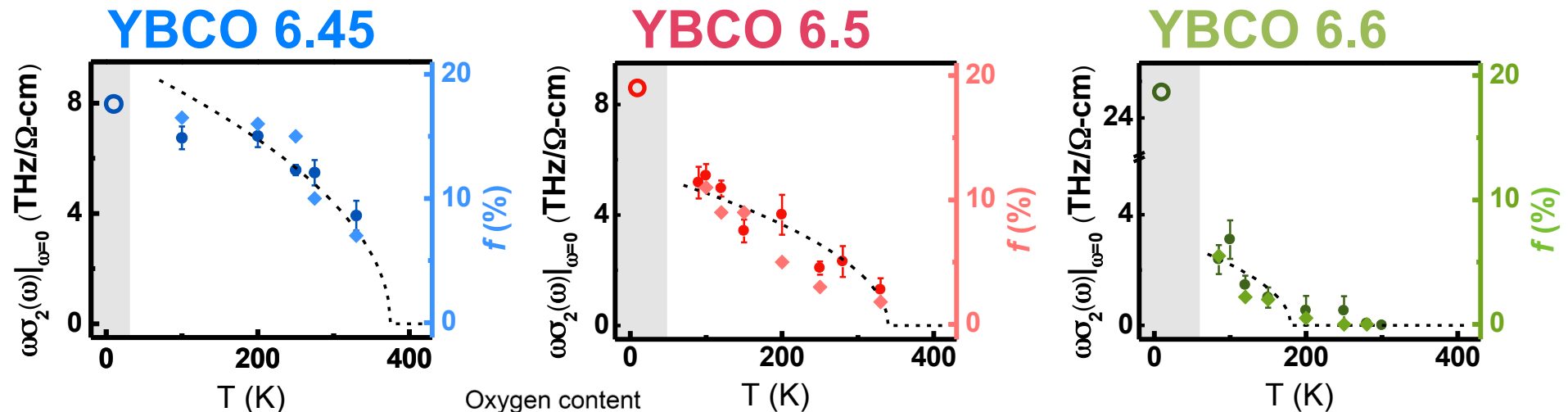


Excellent fit to all optical properties

All temperatures, all doping levels!



Temperature dependence of the superconducting properties



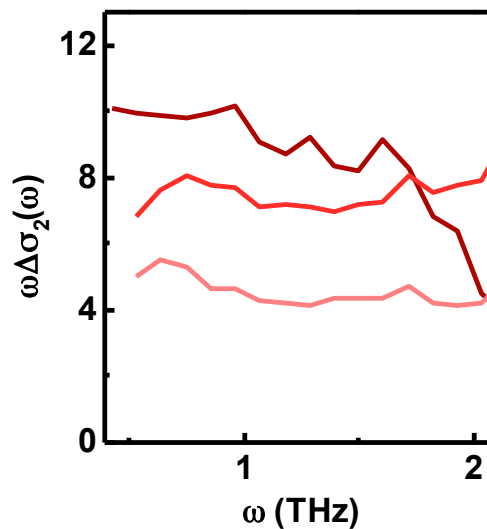
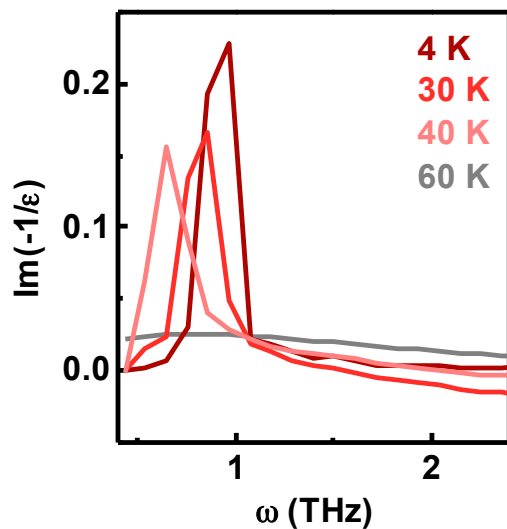
Filling fraction:

$$f(T) \propto \sqrt{1 - \frac{T}{T'}}$$

Phase diagram for photo-induced transient superconductivity

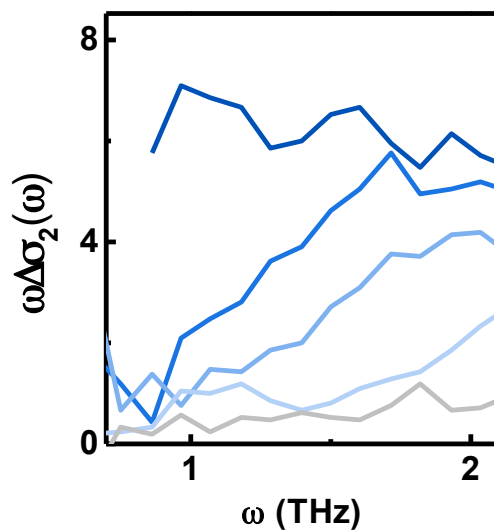
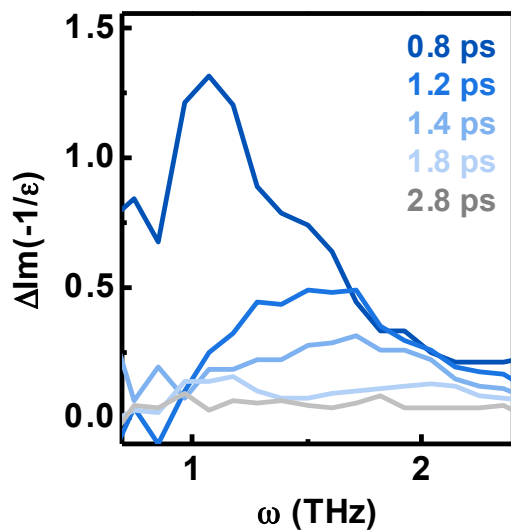
Relaxation of the transient state

Equilibrium:

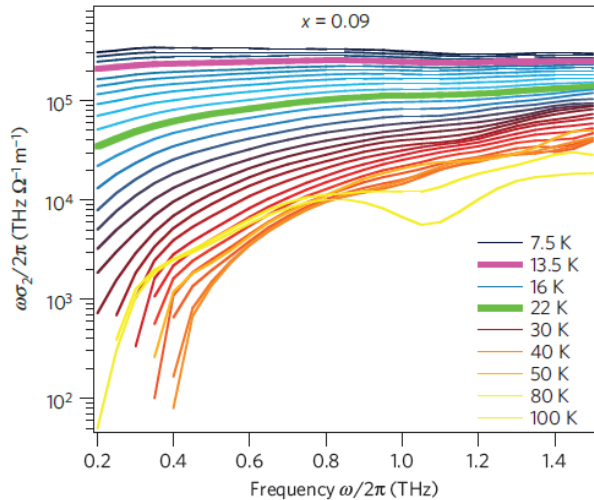


100 K

Transient:



Relaxation of the transient state



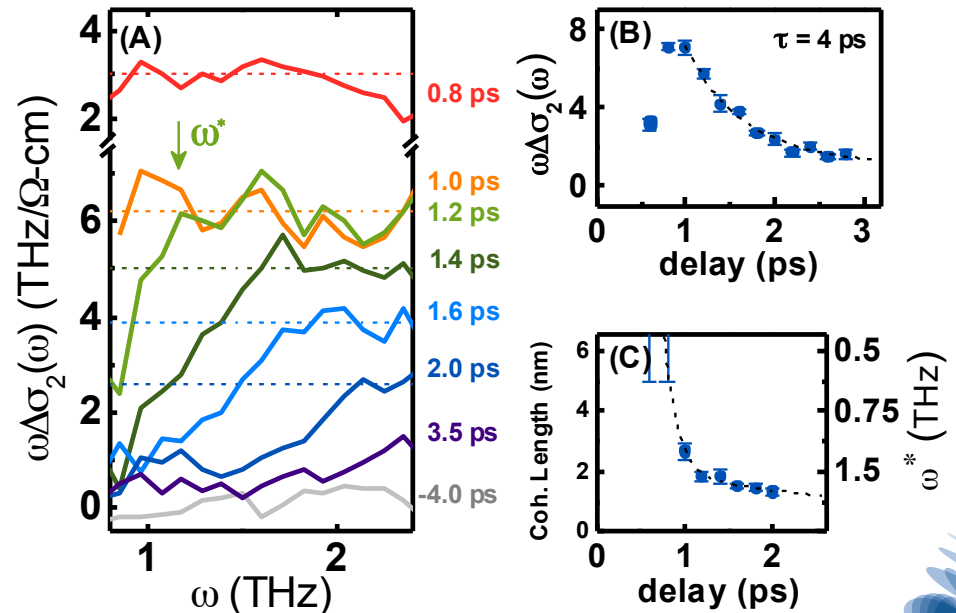
Divergence of σ_2 a measure of coherence

For instance, LSCO—fluctuations above the KTB temperature. [L. S. Bilbro, *et al. Nature Physics*, **7** 298 (2011).]

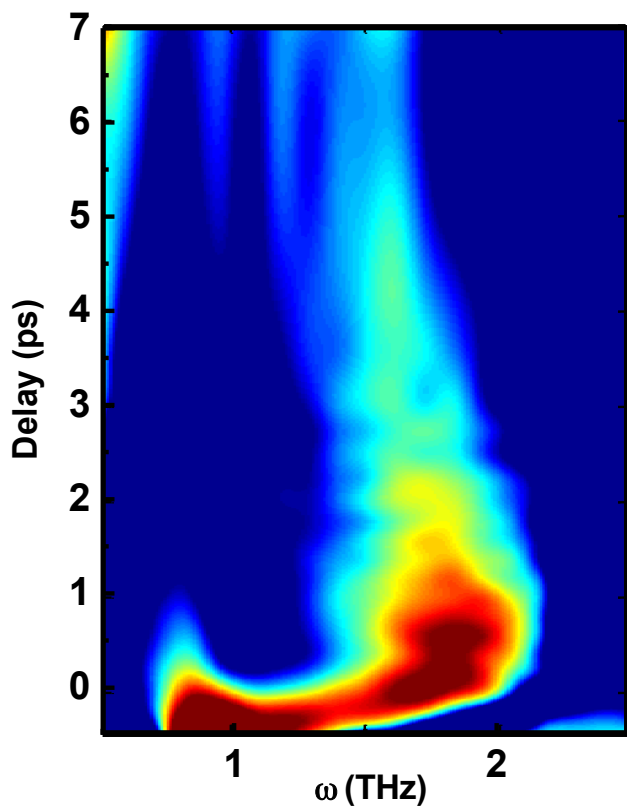
Fluctuations on long time/length scales \rightarrow divergence in $1/\omega$ behavior seen at low frequencies.

Can use “roll off” frequency as a measure of the coherence length scale.

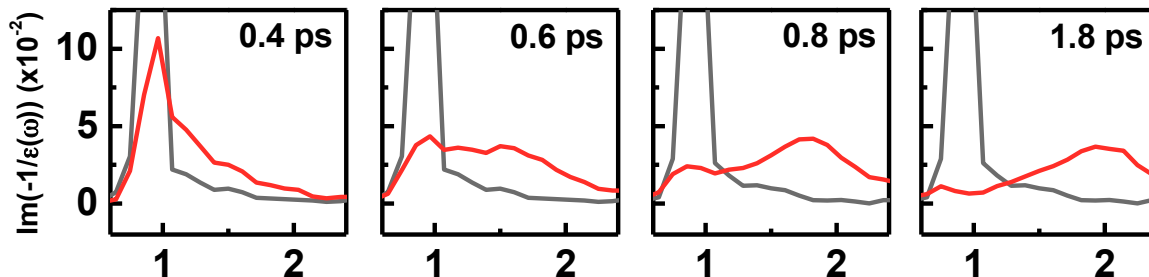
$$d\downarrow coh \sim 2d\downarrow jctn \omega\downarrow J / \omega\uparrow^*$$



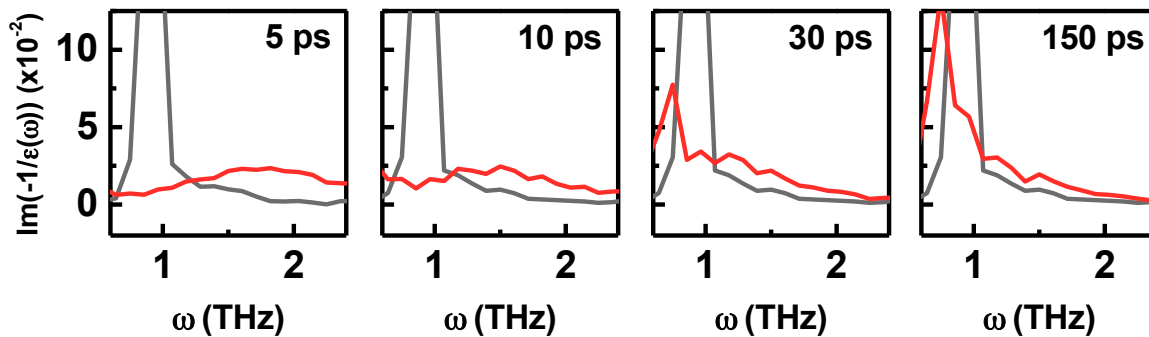
Relaxation of the transient plasmon, $T < T_c$



Blue Shift of the Josephson Plasmon

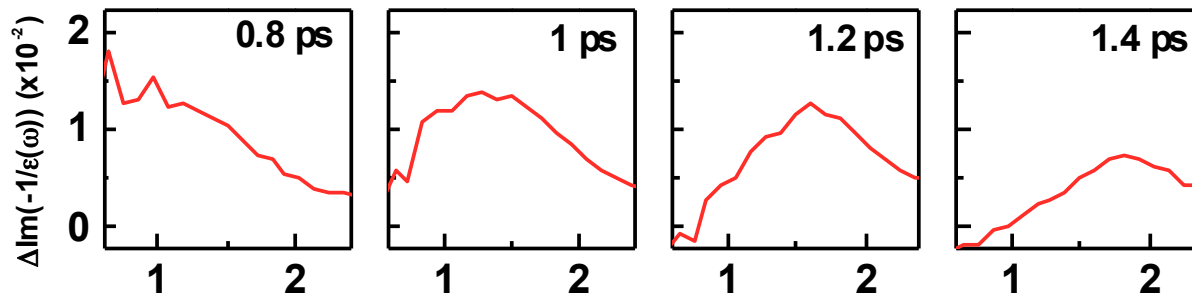


Relaxation of the Josephson Plasmon

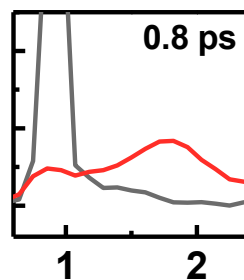
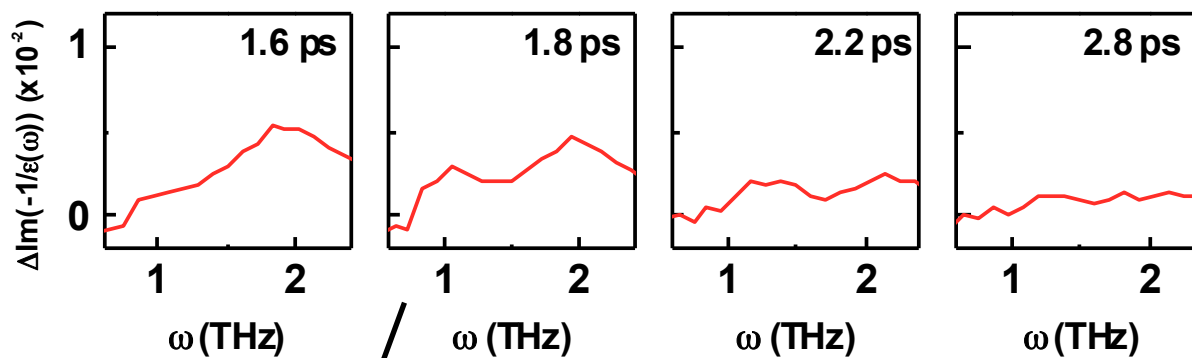


Relaxation of the transient plasmon, $T > T_c$

Appearance of the Plasmon

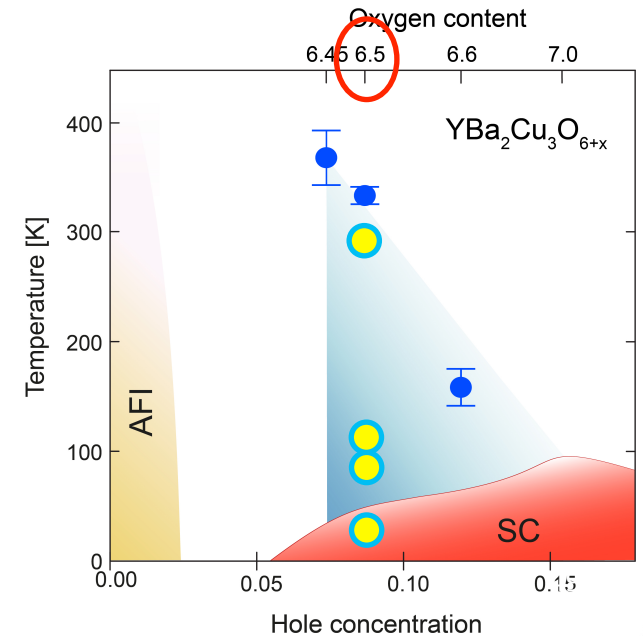


Relaxation of the Plasmon

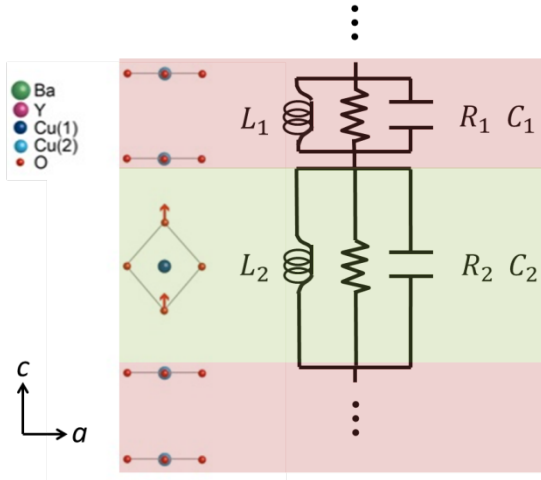


Light-induced superconductivity can be found up to room temperature

➤ Where are the transient superconducting carriers coming from?



Where are the carriers coming from?

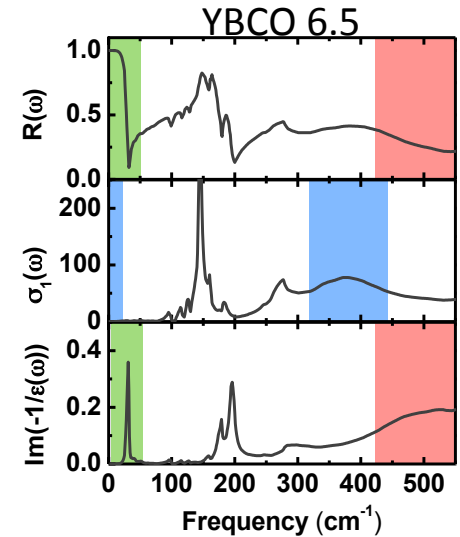
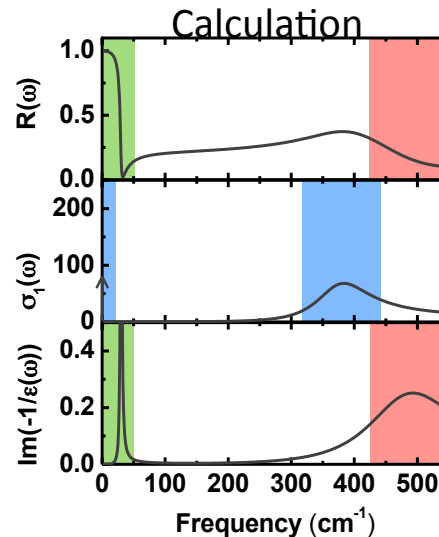
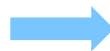


Consider below T_c again.

A second equilibrium Josephson plasma resonance (JPR) at high frequency.

Two modes share spectral weight

- DC delta function
- Finite frequency transverse mode



[calculation based on: D. van der Marel, and A. Tsvetkova. Czech. Journal of Physics 46, 3165]

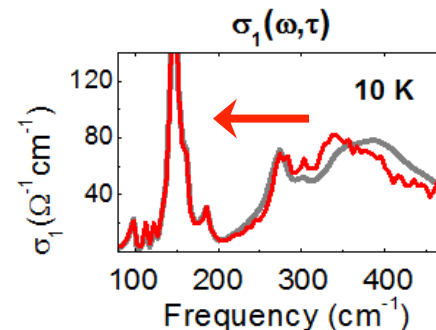
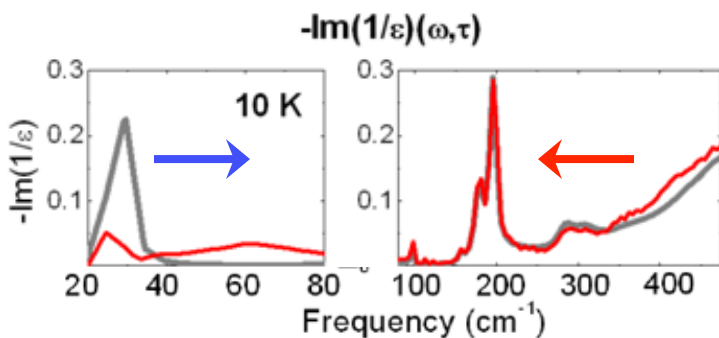
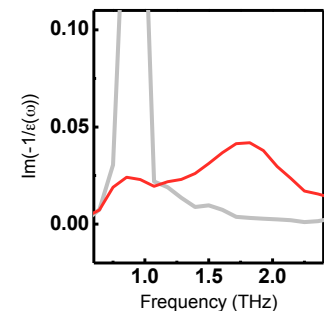
Broadband Response at $T < T_c$: High frequency plasmon weakens

ω_{JPR1} blue shifts

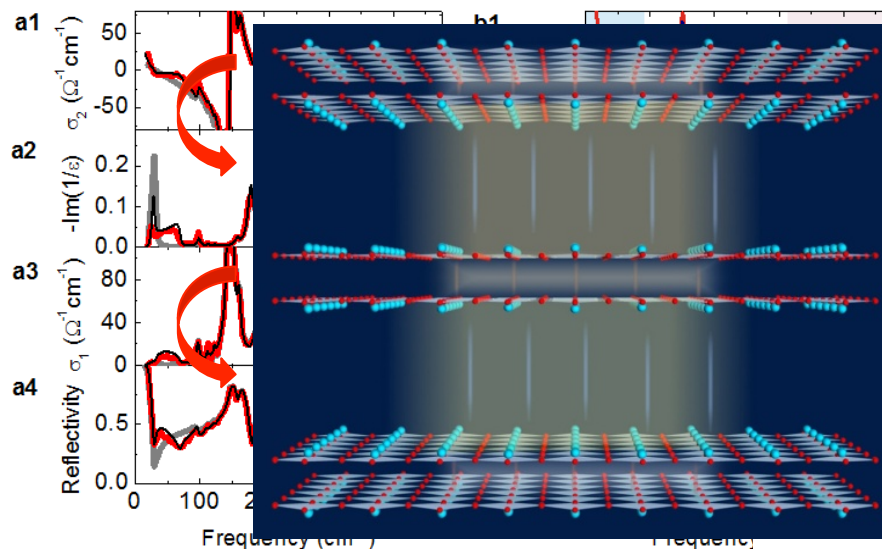
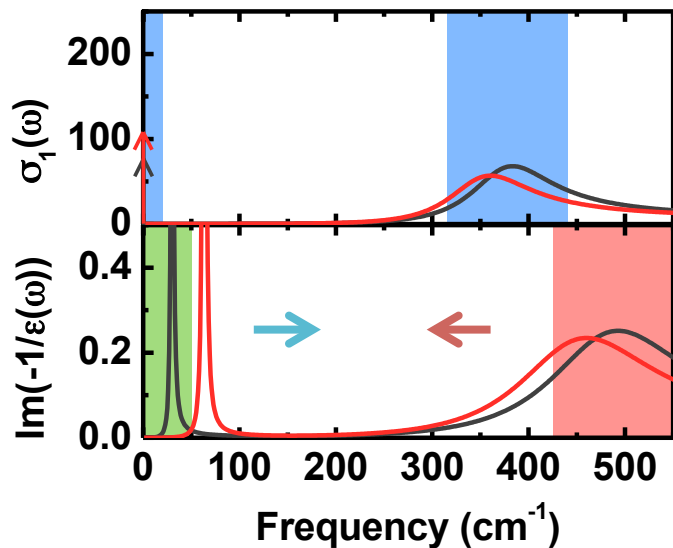
ω_{JPR2} red shifts

Transverse mode

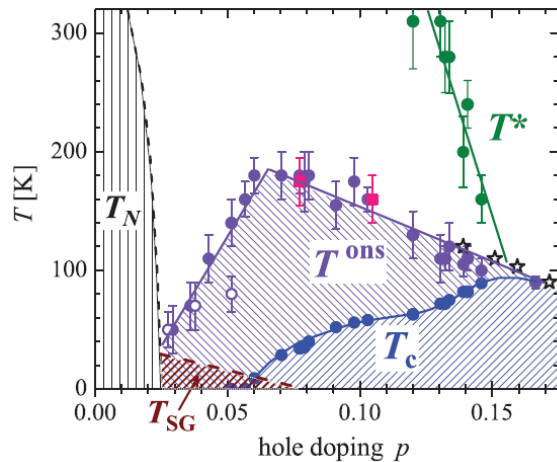
ω_T red shifts



Calculation



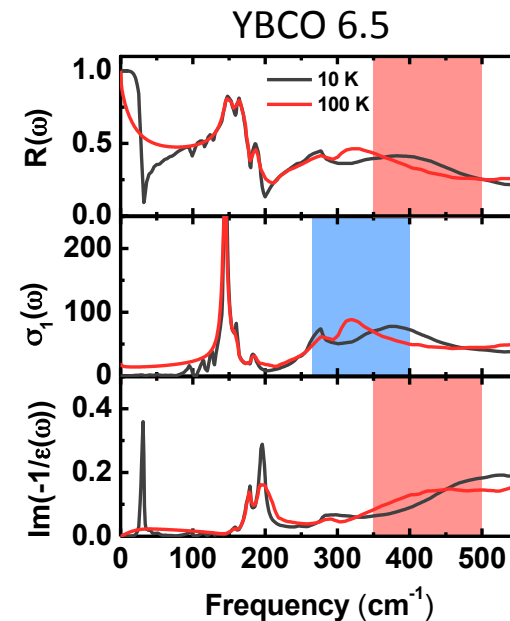
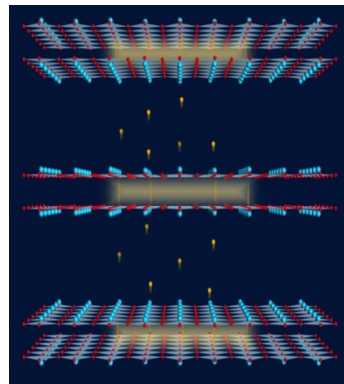
What about above T_c ?



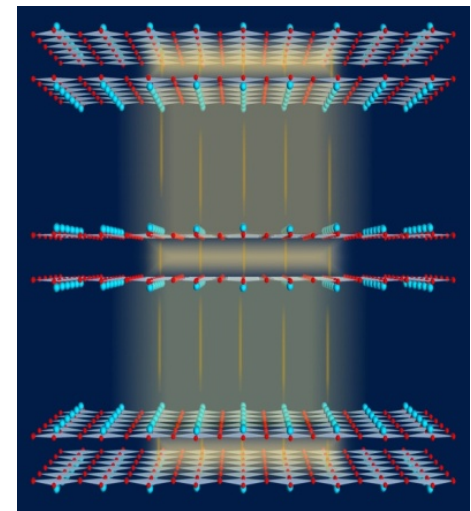
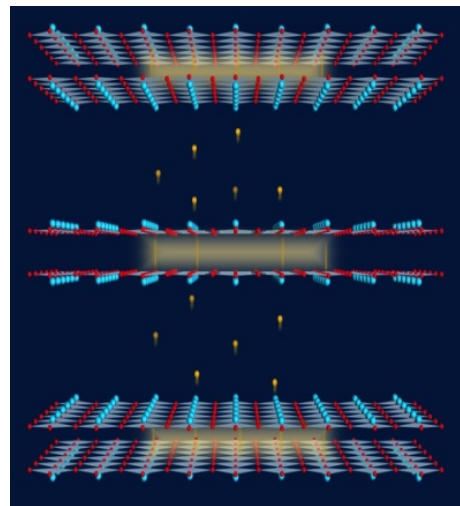
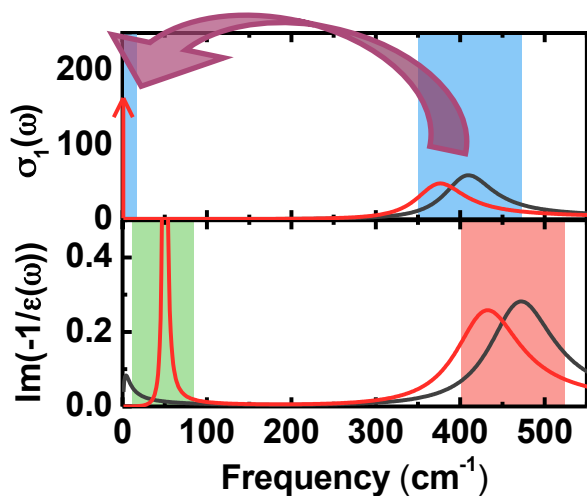
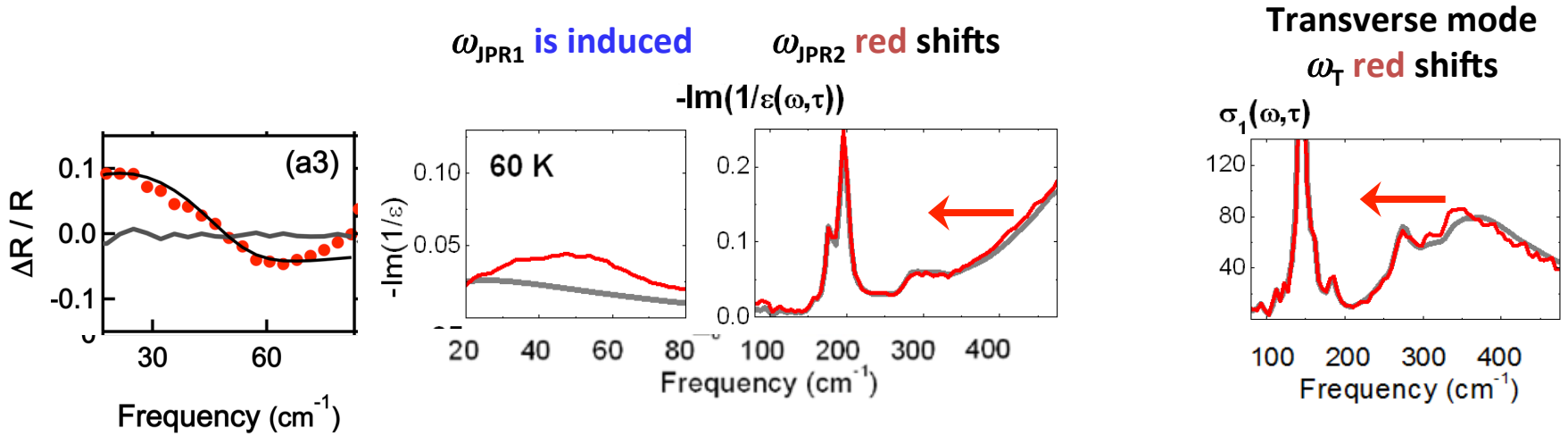
Precursor superconductivity far above T_c
YBCO 6.6: up to 180 K!

[A. Dubroka, *et al.*, PRL **106**, 047006 (2011).]

Effect attributed to conductivity across **bilayers** only

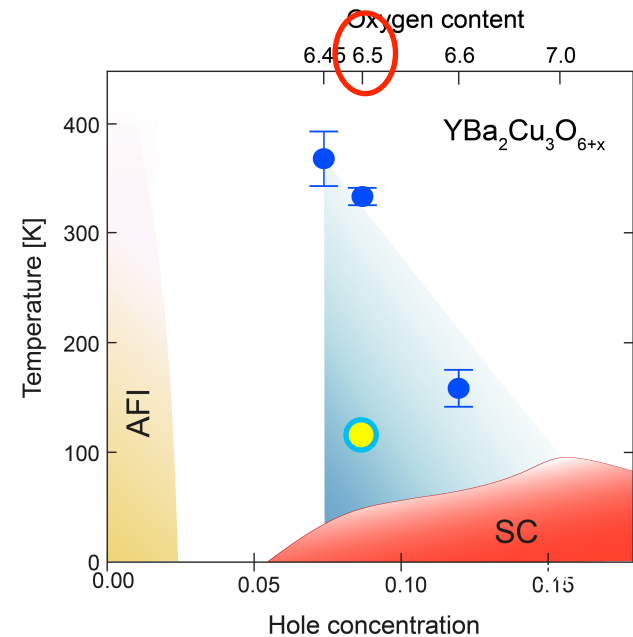
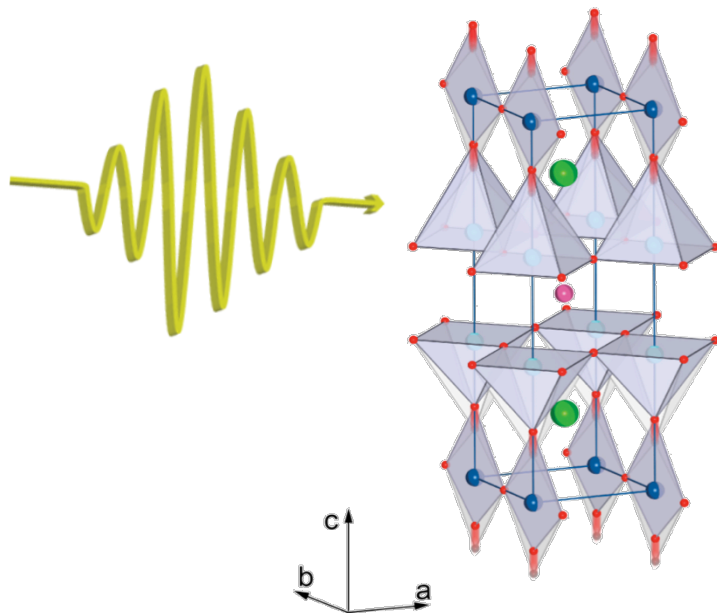


Broadband Response at $T > T_c$: Redistribution of coherence



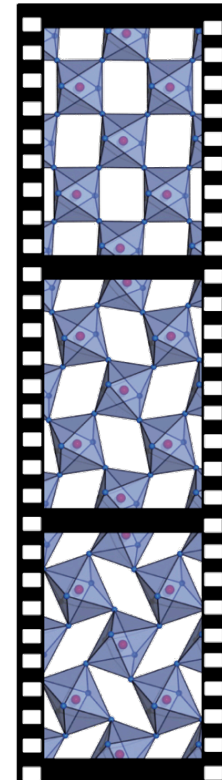
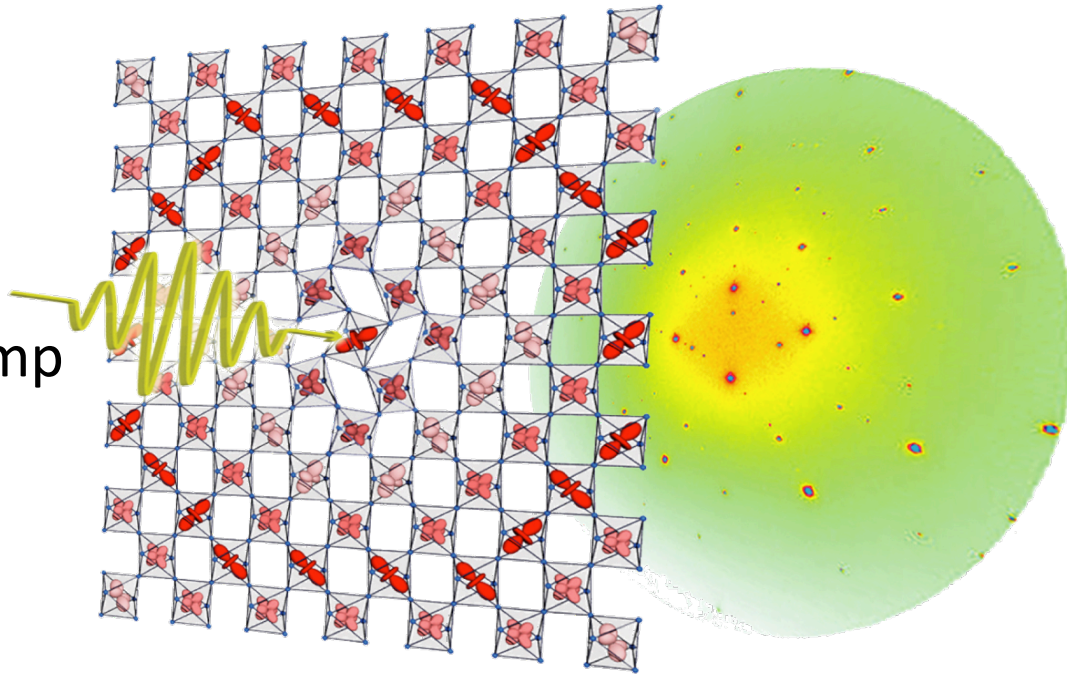
Light-induced superconductivity can be found
Cooper pairs transfer between the bilayers

What are the atoms doing?



fs-crystallography

15 μm
phonon pump



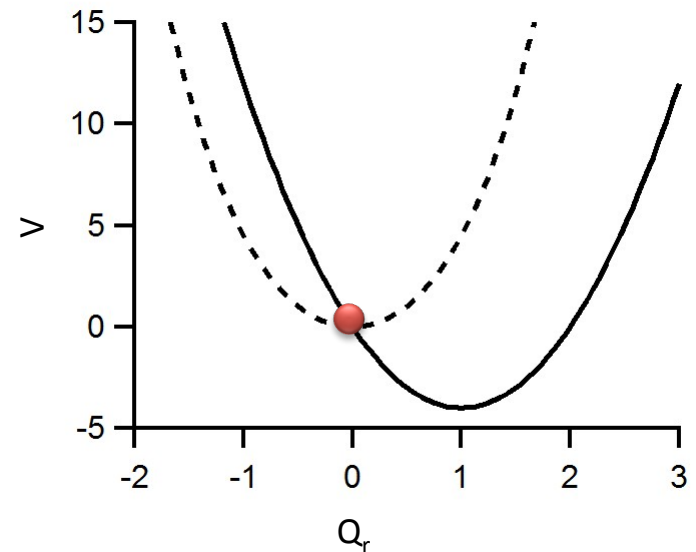
• Simple Classical Model of Nonlinear Phononics

- Potential energy of the Raman mode as function of IR mode Amplitude

$$V_r = \frac{1}{2} \mu_r \omega_r^2 Q_r^2 - \boxed{A Q_{ir}^2 Q_r}$$



Displacement of Raman mode



- Equation of Motion $\mu_r \ddot{Q}_r + 2\gamma_r \mu_r \dot{Q}_r + \mu_r \omega_r^2 Q_r = A Q_{ir}^2$

•Equation of Motion $\mu_r \ddot{Q}_r + 2\gamma_r \mu_r \dot{Q}_r + \mu_r \omega_r^2 Q_r = A Q_{ir}^2$

damped harmonic oscillator

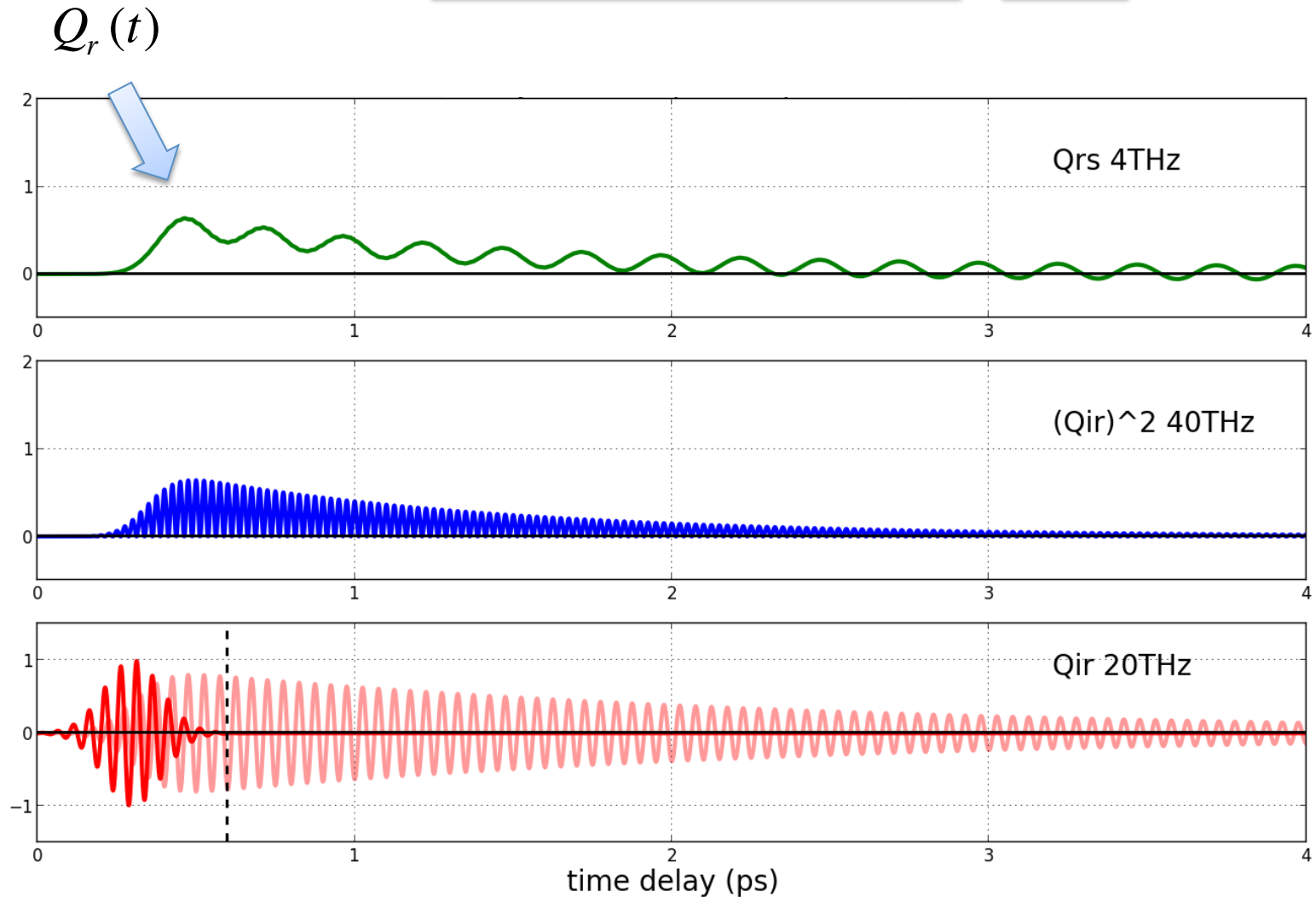
•Equation of Motion

$$\mu_r \ddot{Q}_r + 2\gamma_r \mu_r \dot{Q}_r + \mu_r \omega_r^2 Q_r = A Q_{ir}^2$$

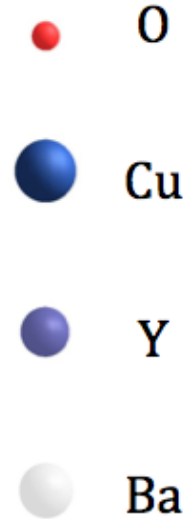
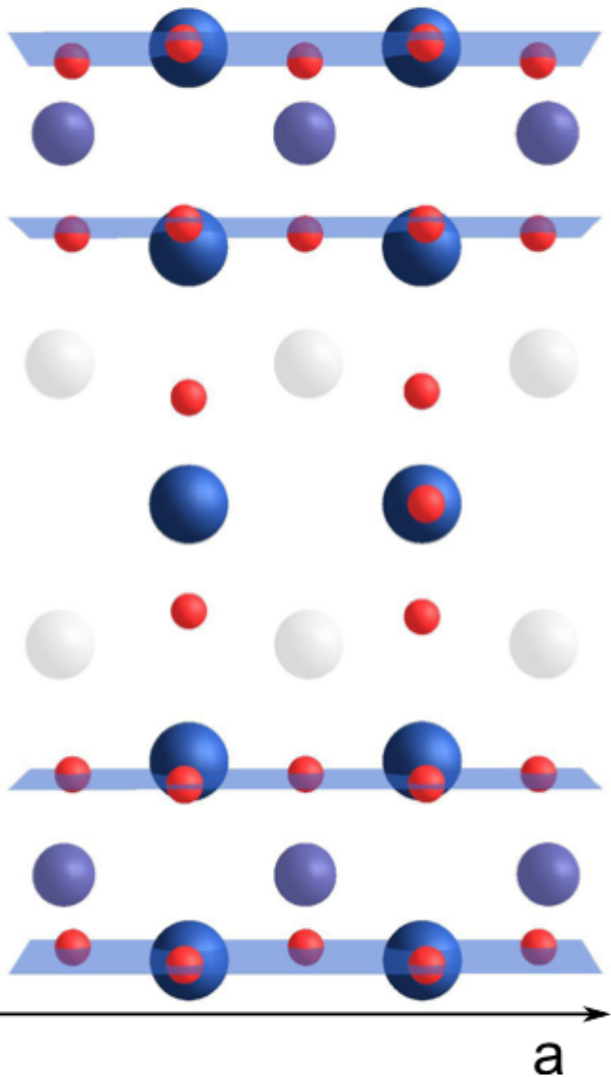
damped harmonic oscillator

driving force

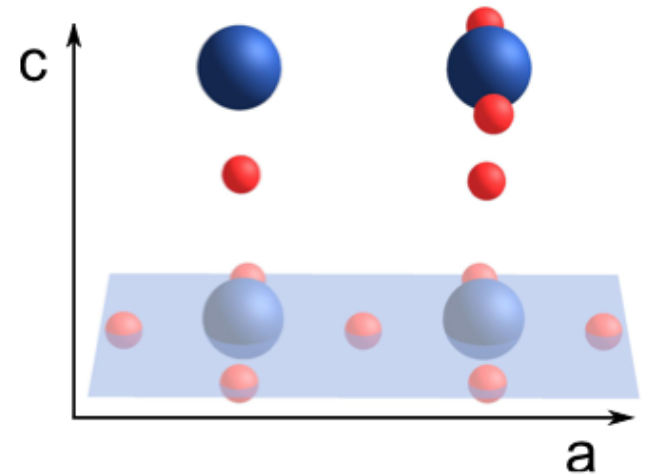
• Equation of Motion $\mu_r \ddot{Q}_r + 2\gamma_r \mu_r \dot{Q}_r + \mu_r \omega_r^2 Q_r = A Q_{ir}^2$



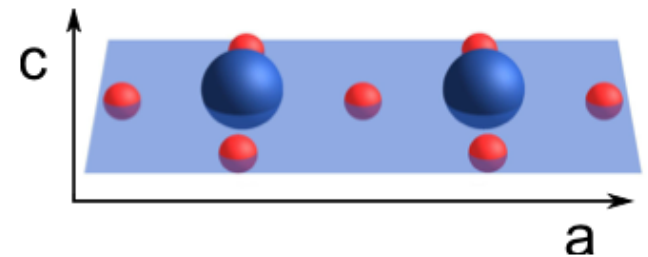
- YBCO6.5 OrthoII



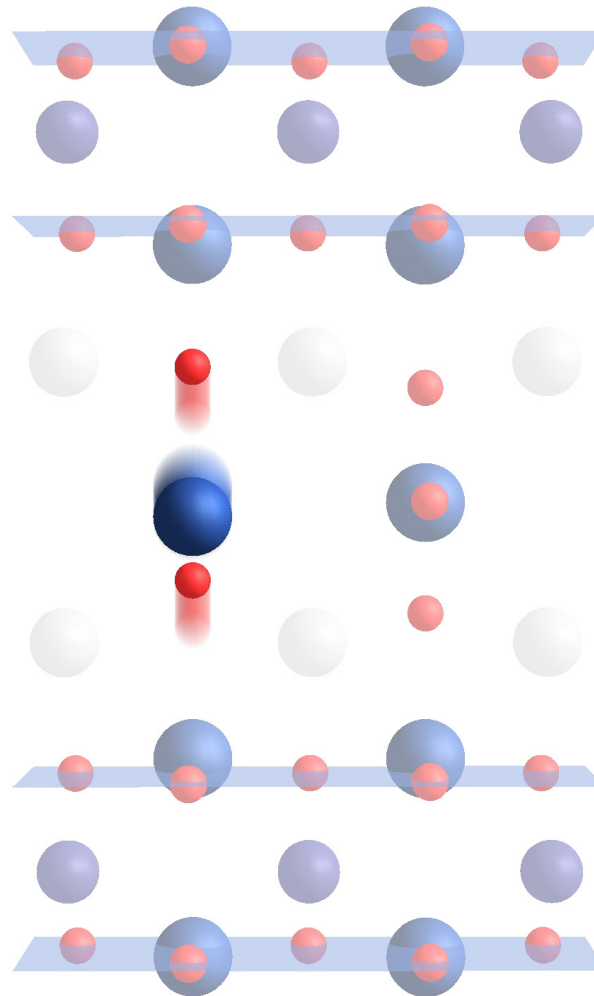
Chains



CuO₂ Planes



B_{1u} Mode



•Symmetry Analysis of YBCO

\otimes	A_g	B_{1g}	B_{2g}	B_{3g}	A_u	B_{1u}	B_{2u}	B_{3u}
A_g	A_g	B_{1g}	B_{2g}	B_{3g}	A_u	B_{1u}	B_{2u}	B_{3u}
B_{1g}	B_{1g}	A_g	B_{3g}	B_{2g}	B_{1u}	A_u	B_{3u}	B_{2u}
B_{2g}	B_{2g}	B_{3g}	A_g	B_{1g}	B_{2u}	B_{3u}	A_u	B_{1u}
B_{3g}	B_{3g}	B_{2g}	B_{1g}	A_g	B_{3u}	B_{2u}	B_{1u}	A_u
A_u	A_u	B_{1u}	B_{2u}	B_{3u}	A_g	B_{1g}	B_{2g}	B_{3g}
B_{1u}	B_{1u}	A_u	B_{3u}	B_{2u}	B_{1g}	A_g	B_{3g}	B_{2g}
B_{2u}	B_{2u}	B_{3u}	A_u	B_{1u}	B_{2g}	B_{3g}	A_g	B_{1g}
B_{3u}	B_{3u}	B_{2u}	B_{1u}	A_u	B_{3g}	B_{2g}	B_{1g}	A_g

A_g modes

$$B_{1u} \times B_{1u} = A_g$$

•Symmetry Analysis of YBCO

\otimes	A_g	B_{1g}	B_{2g}	B_{3g}	A_u	B_{1u}	B_{2u}	B_{3u}
A_g	A_g	B_{1g}	B_{2g}	B_{3g}	A_u	B_{1u}	B_{2u}	B_{3u}
B_{1g}	B_{1g}	A_g	B_{3g}	B_{2g}	B_{1u}	A_u	B_{3u}	B_{2u}
B_{2g}	B_{2g}	B_{3g}	A_g	B_{1g}	B_{2u}	B_{3u}	A_u	B_{1u}
B_{3g}	B_{3g}	B_{2g}	B_{1g}	A_g	B_{3u}	B_{2u}	B_{1u}	A_u
A_u	A_u	B_{1u}	B_{2u}	B_{3u}	A_g	B_{1g}	B_{2g}	B_{3g}
B_{1u}	B_{1u}	A_u	B_{3u}	B_{2u}	B_{1g}	A_g	B_{3g}	B_{2g}
B_{2u}	B_{2u}	B_{3u}	A_u	B_{1u}	B_{2g}	B_{3g}	A_g	B_{1g}
B_{3u}	B_{3u}	B_{2u}	B_{1u}	A_u	B_{3g}	B_{2g}	B_{1g}	A_g

YBCO OrthoII – 11 A_g modes

A_g modes

B_g modes

$$B_{1u} \times B_{1u} = A_g$$



M
P
S
D

Determination of Coupling Strength



DFT Calculations

(Alaska Subedi, Antoine Georges)

DFT Calculations

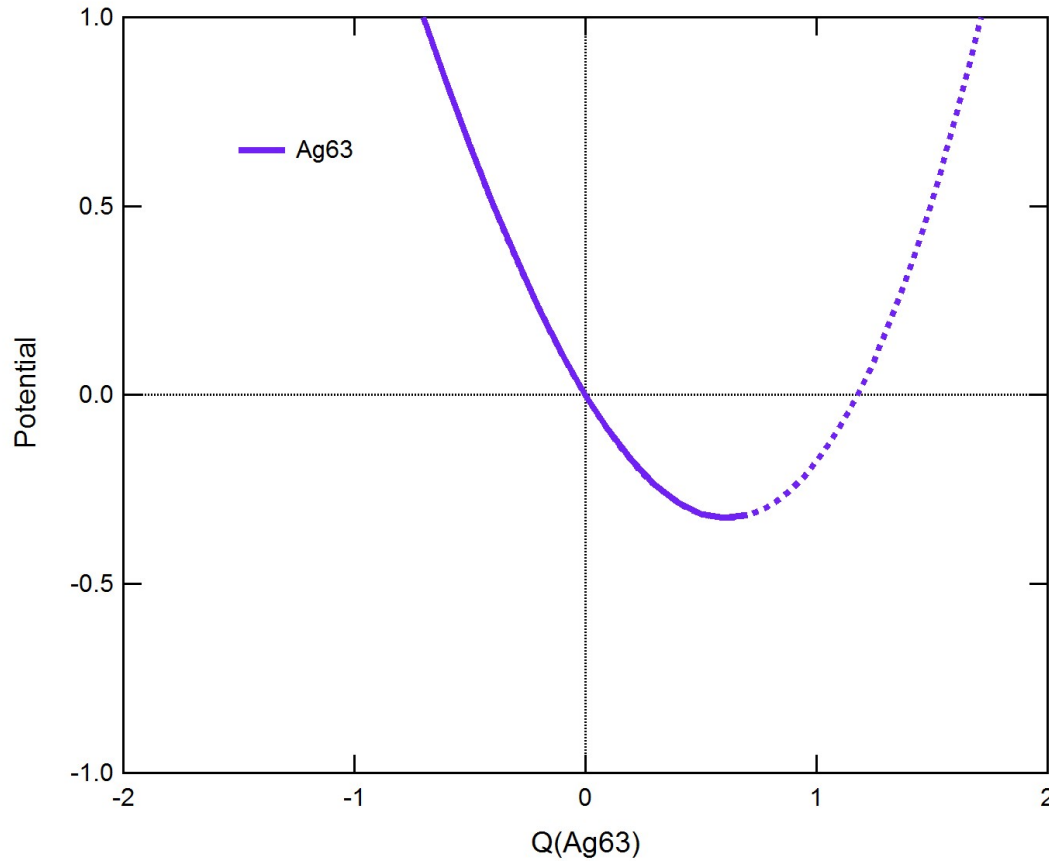
(Alaska Subedi, Antoine Georges)

Input B1u amplitude

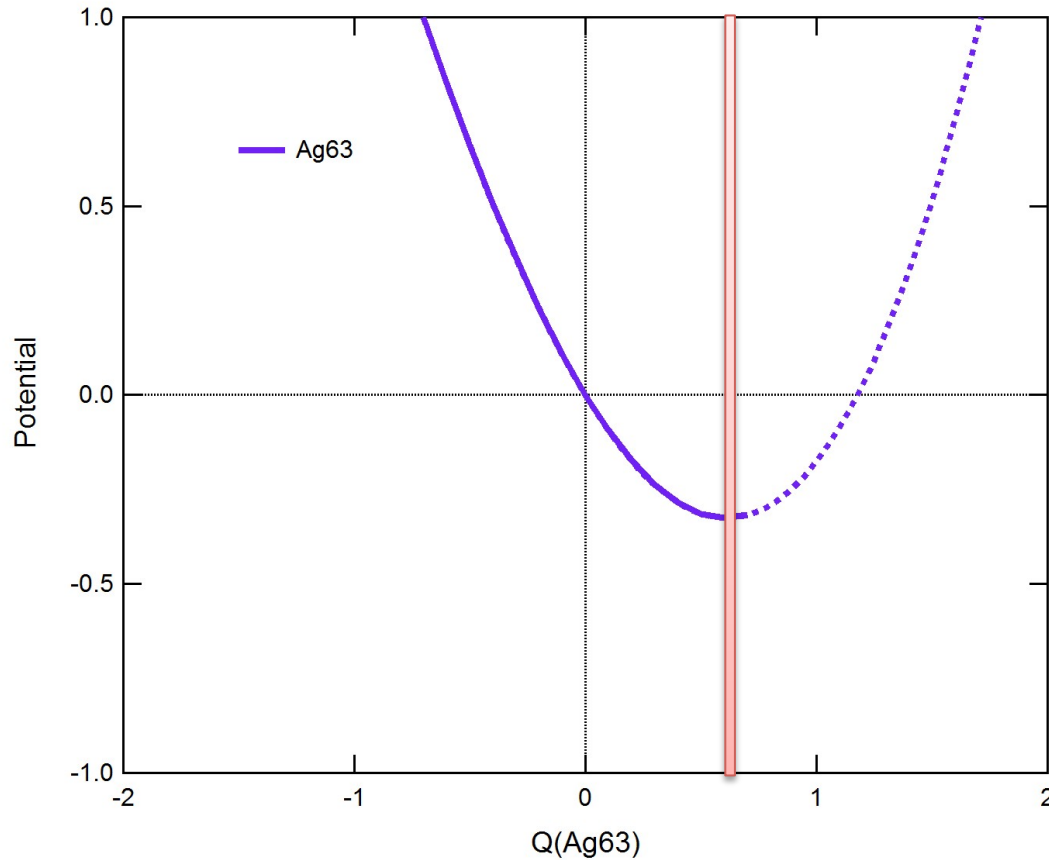


Output Raman mode coupling strengths

- Results of DFT calculations

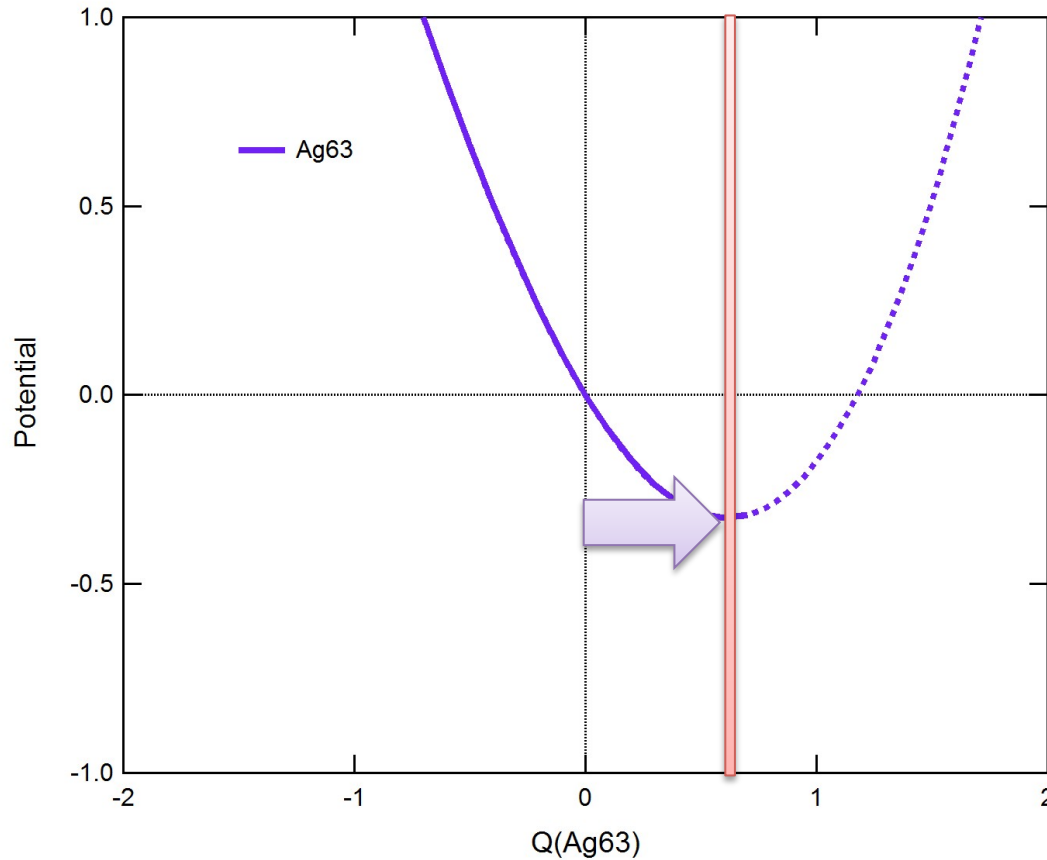


- Example for determination of coupling strength



Raman mode Energy potential shifts

• Example for determination of coupling strength



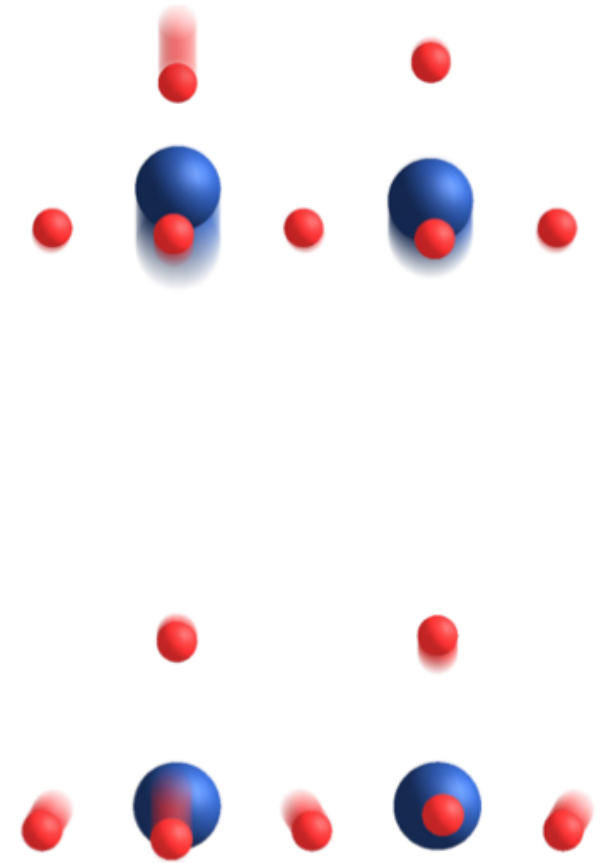
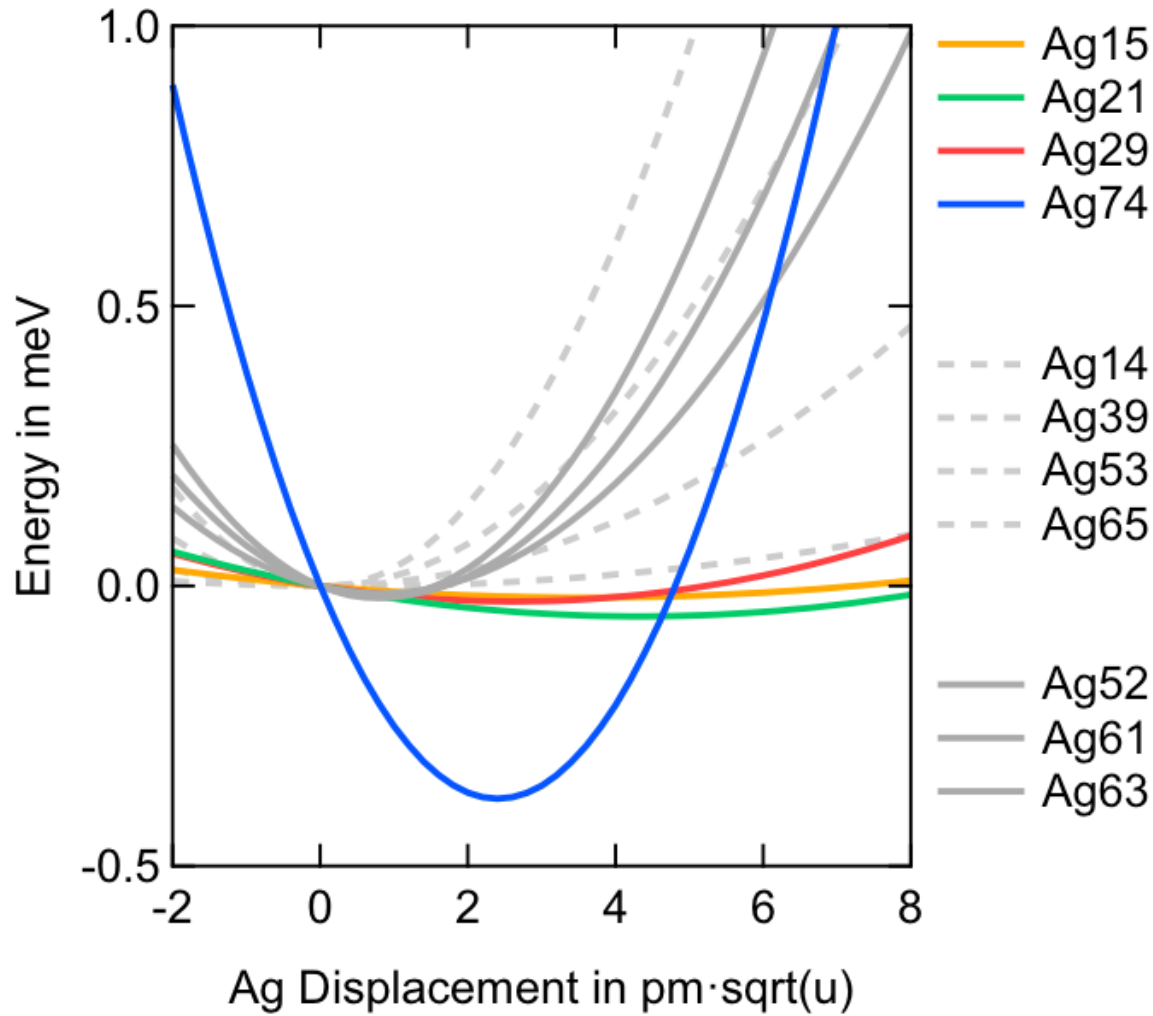
Raman mode Energy potential shifts

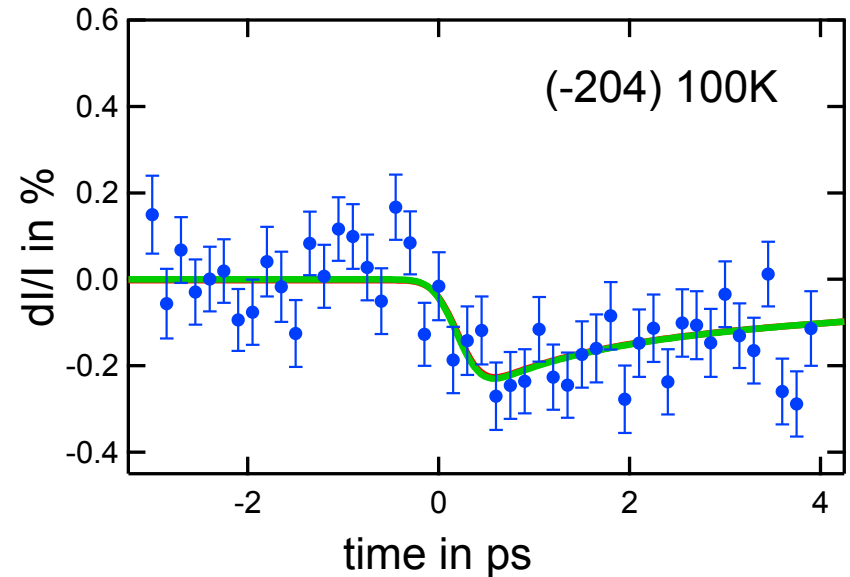
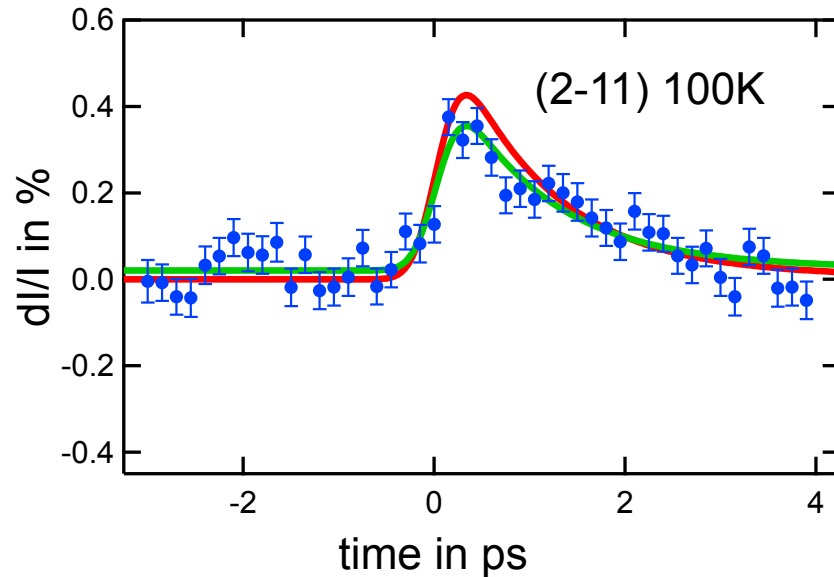
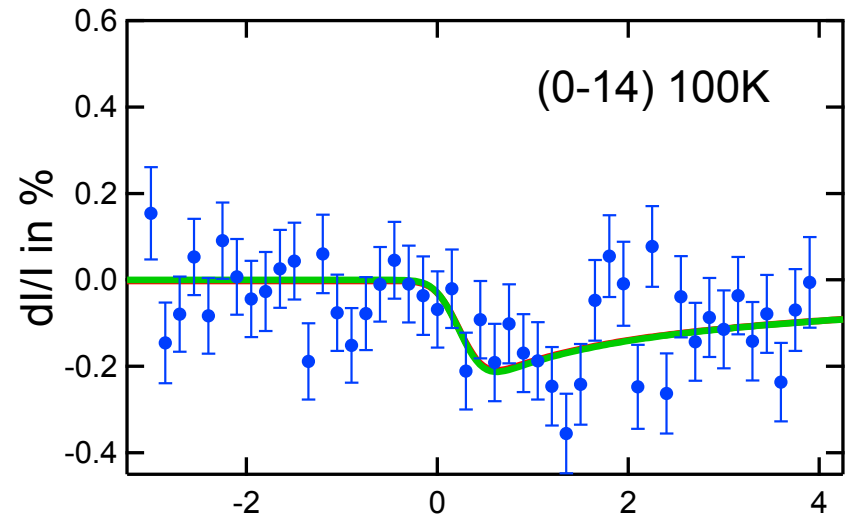
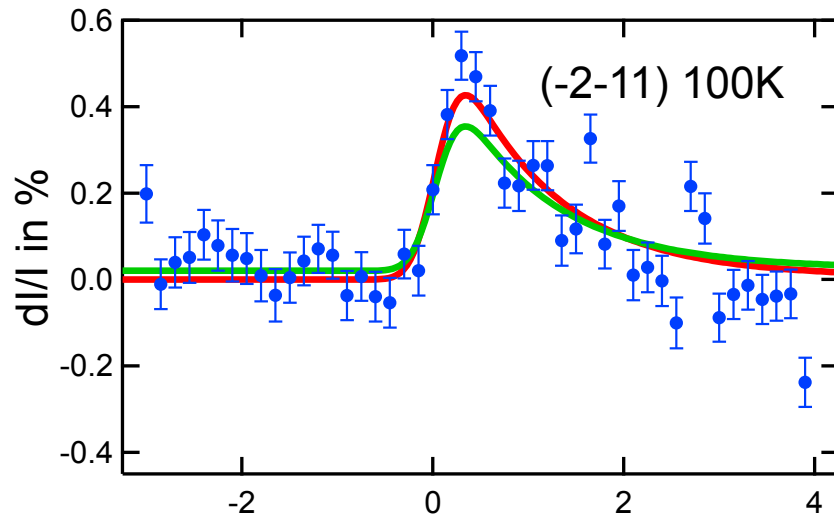
Size of the shift gives coupling strength



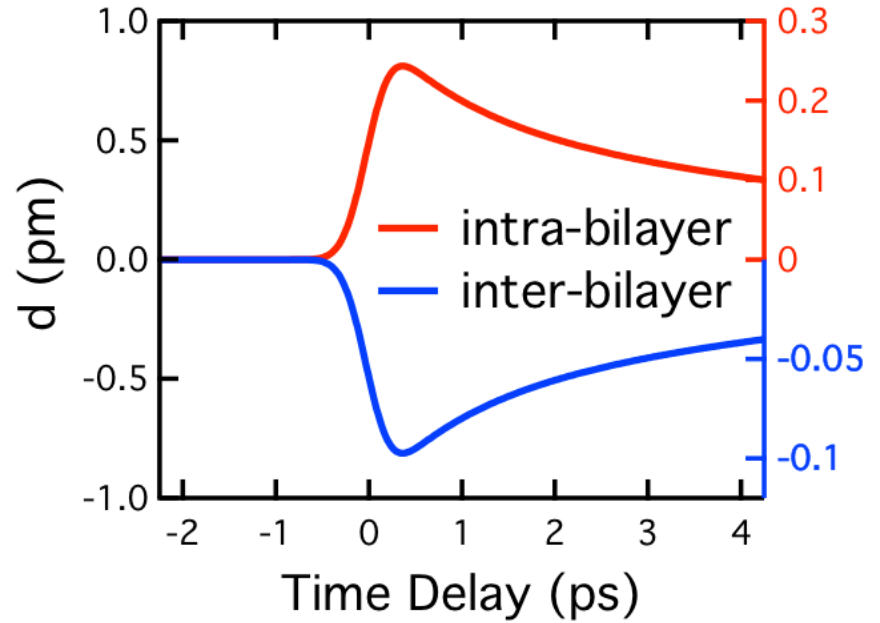
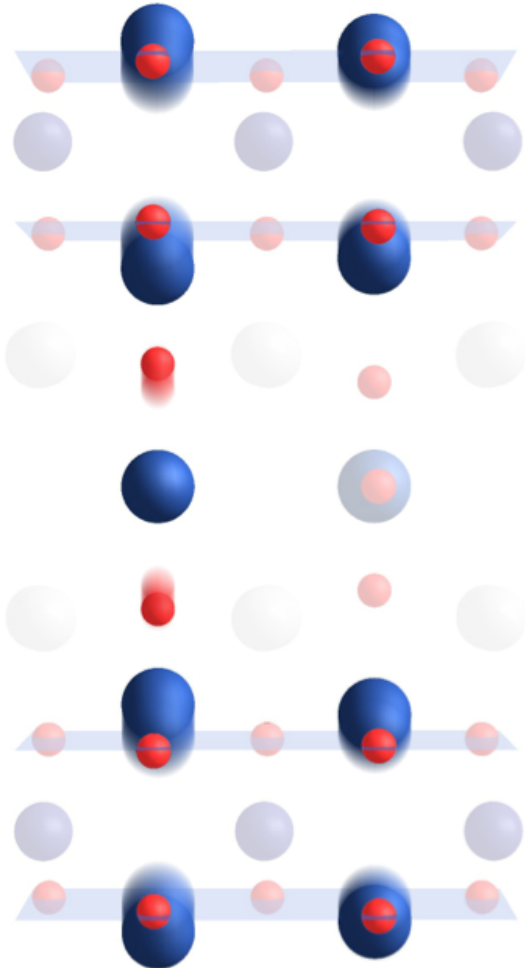
Fit only one Amplitude to experimental Data

• Relative coupling strengths



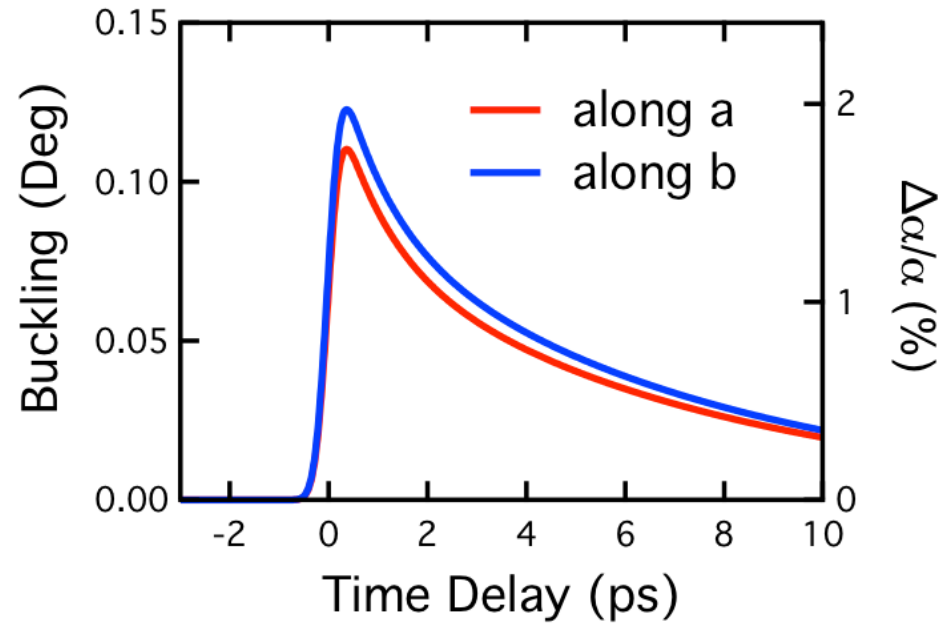
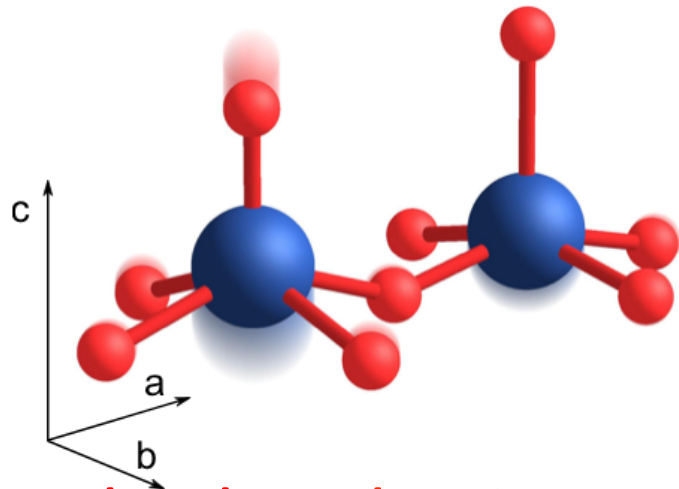
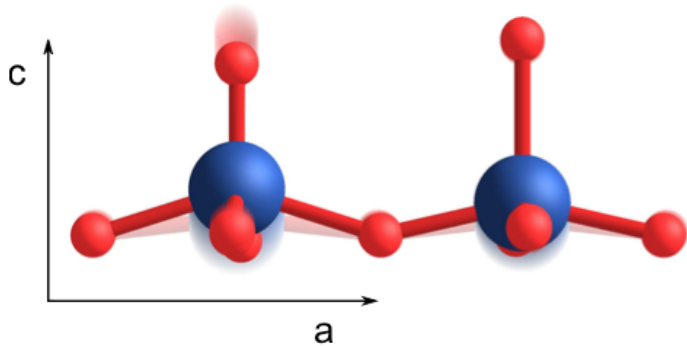


- Lattice Rearrangement

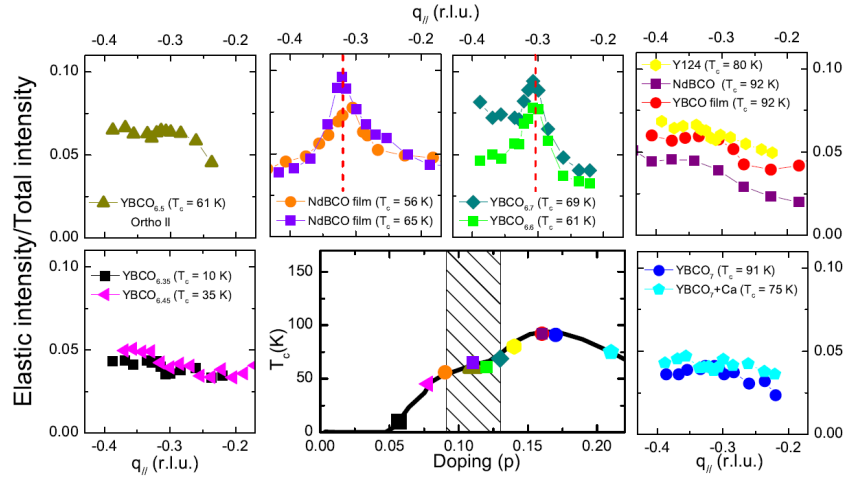


Is this the (unstable) atomic structure of a room temperature superconductor?

- Buckling



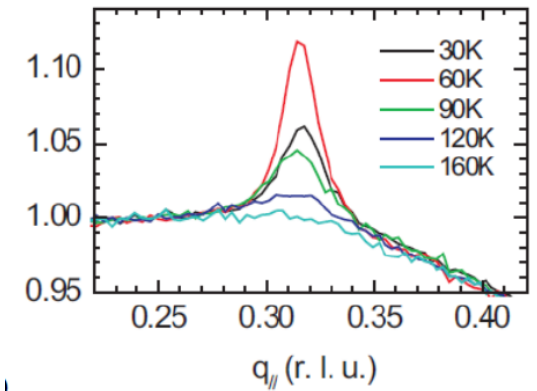
Microscopic Physics: Hypothesis nr. 2



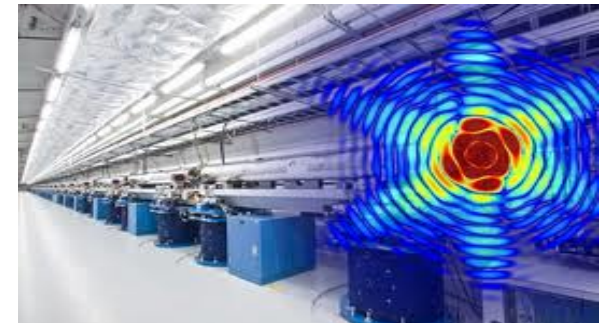
Is this atomic motion melting competing orders ?

Ghiringhelli et al. Science 337, 821 (2012)

Resonant Soft X-ray Diffraction (Cu L-edge) reveals bi-axial in-plane charge density wave, peaking at T_c

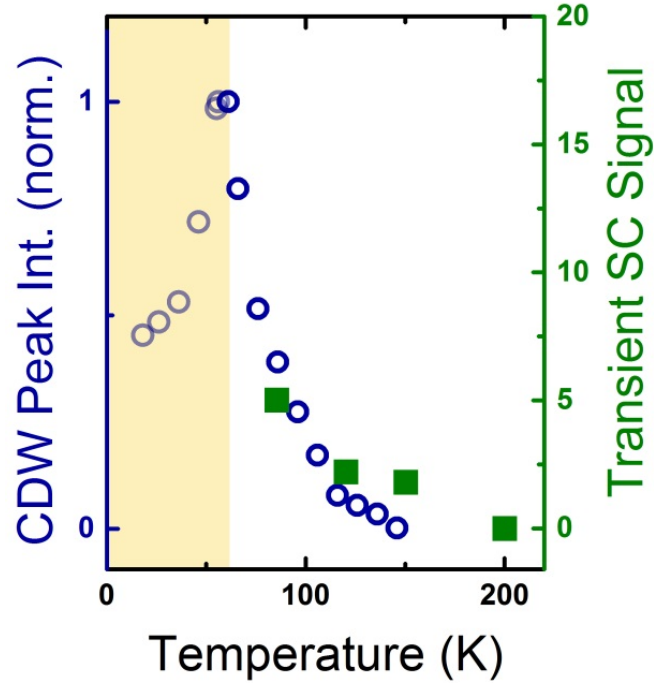
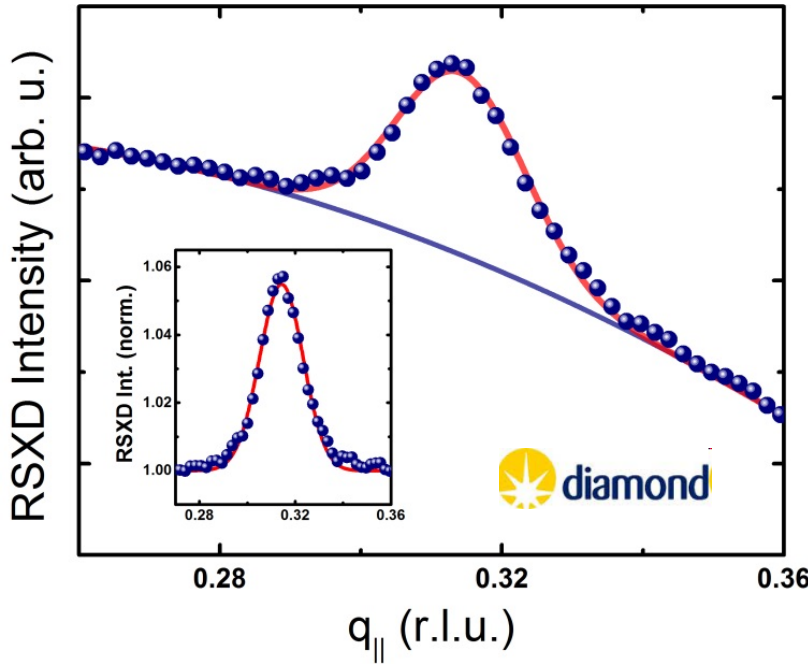


LCLS soft x-ray experiments



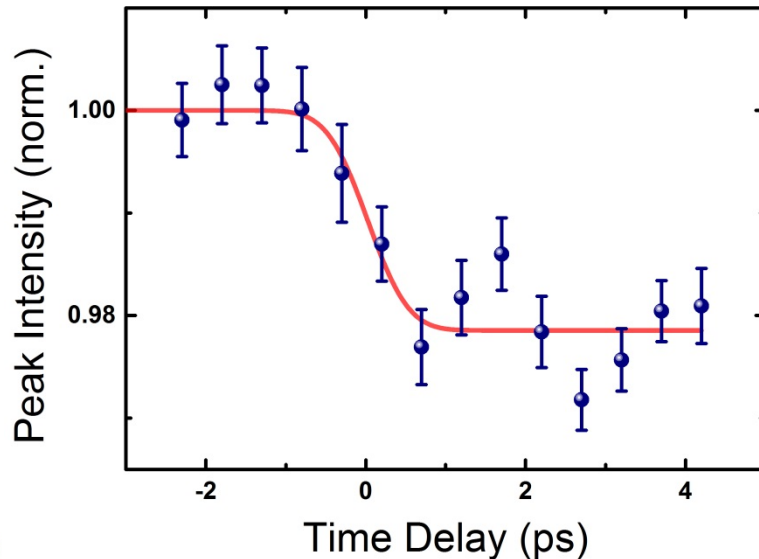
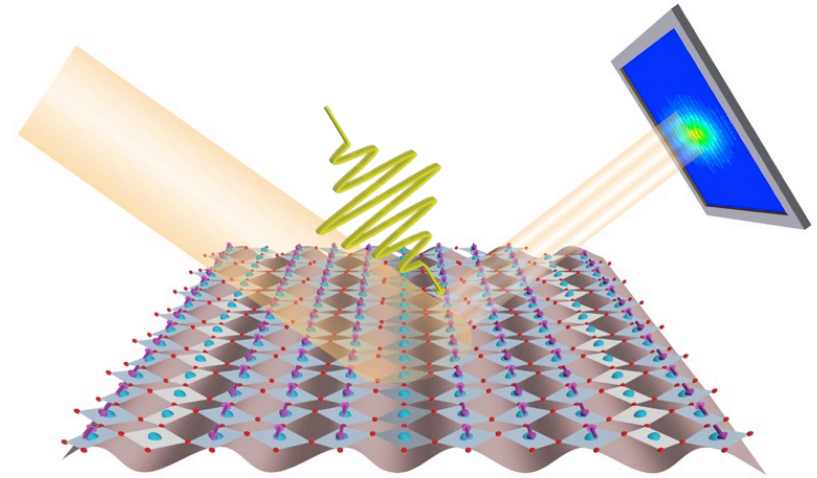
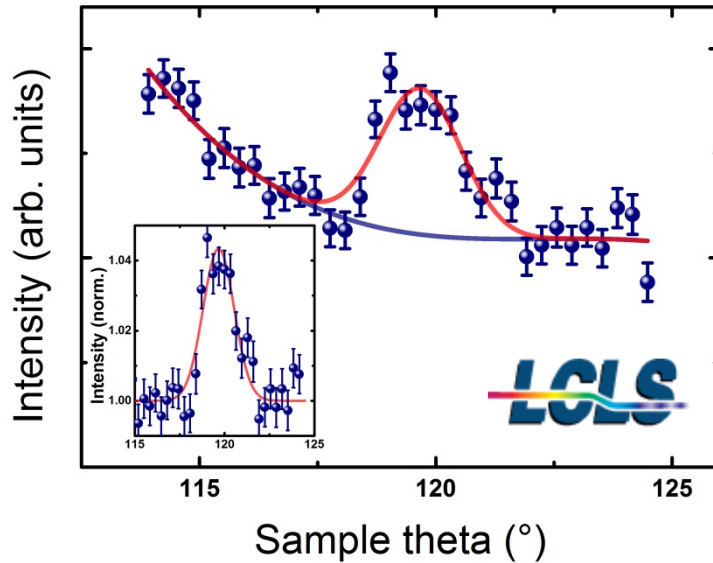
Temperature dependence of the coherent and CDW state

Comparison to the strength of charge correlations above T_c



YBCO 6.6, $T_c=60$ K

Transient reduction of the CDW order

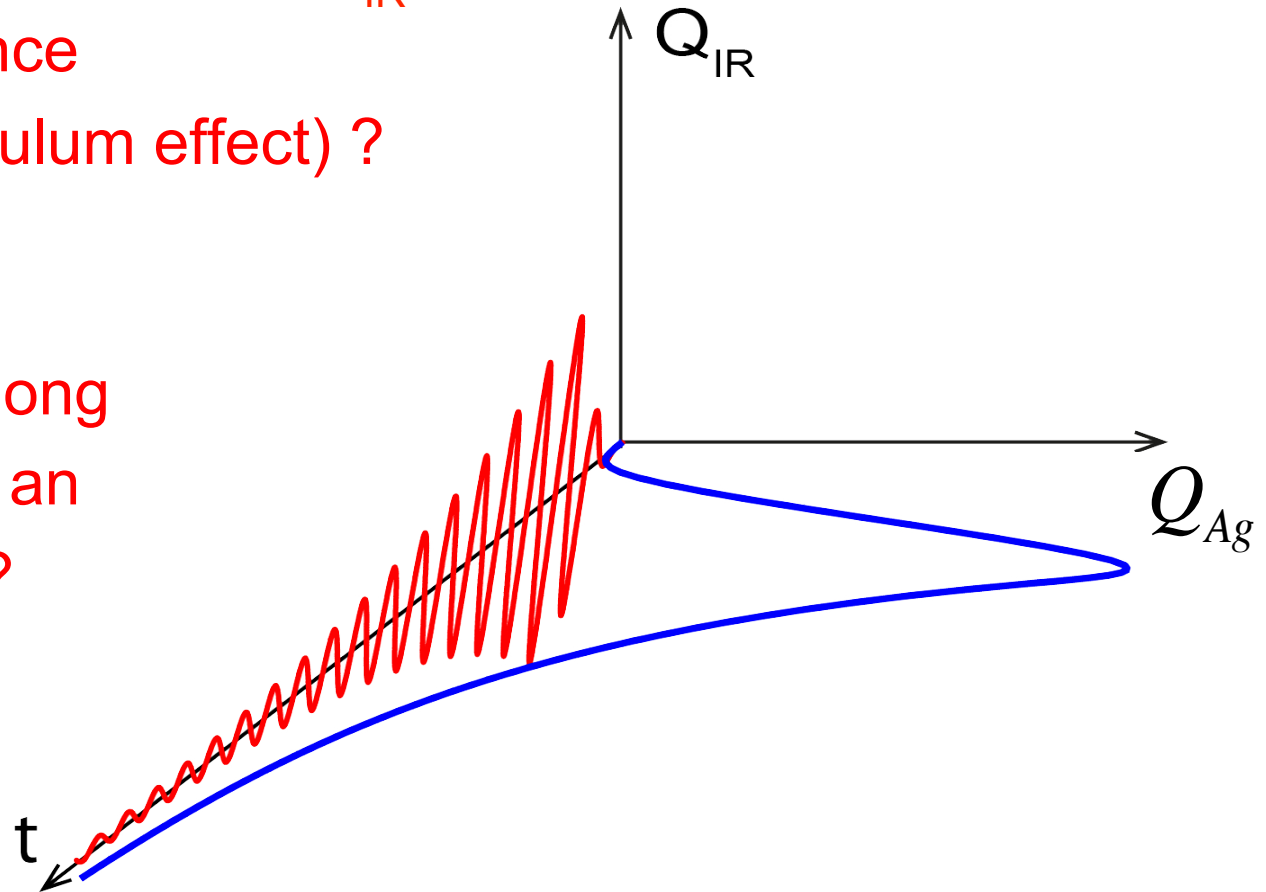


- We observe 50% melting of the CDW peak
- Time scale comparable to the appearance of transient SC
- Lifetime exceeds the lifetime of transient SC

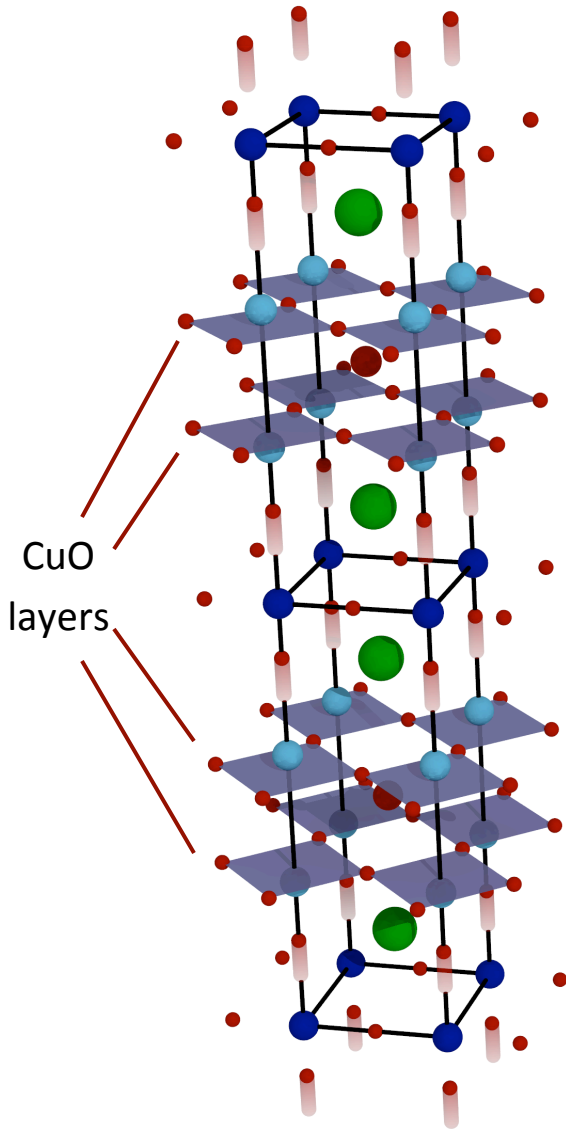
Microscopic Physics: Hypothesis nr. 3

Is the dynamic modulation of Q_{IR}
stabilizing coherence
(see Kaptiza pendulum effect) ?

Is the distortion along
the A_g coordinate an
epiphenomenon ?



Reducing phase fluctuations via driving

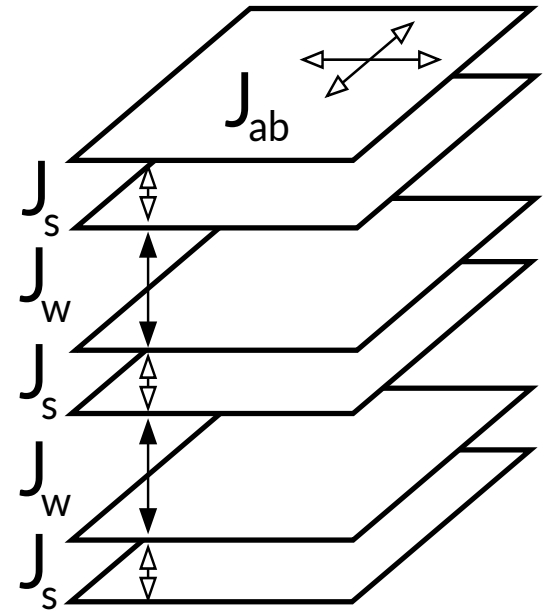


Modeled via a layered system of coupled Josephson junctions

$$H_{\theta} = - \sum_{\langle ij \rangle} J_{ij} \cos(\theta_i - \theta_j)$$

Anisotropy:

$$J_{ab} : J_s : J_w \sim 1000 : 100 : 1$$

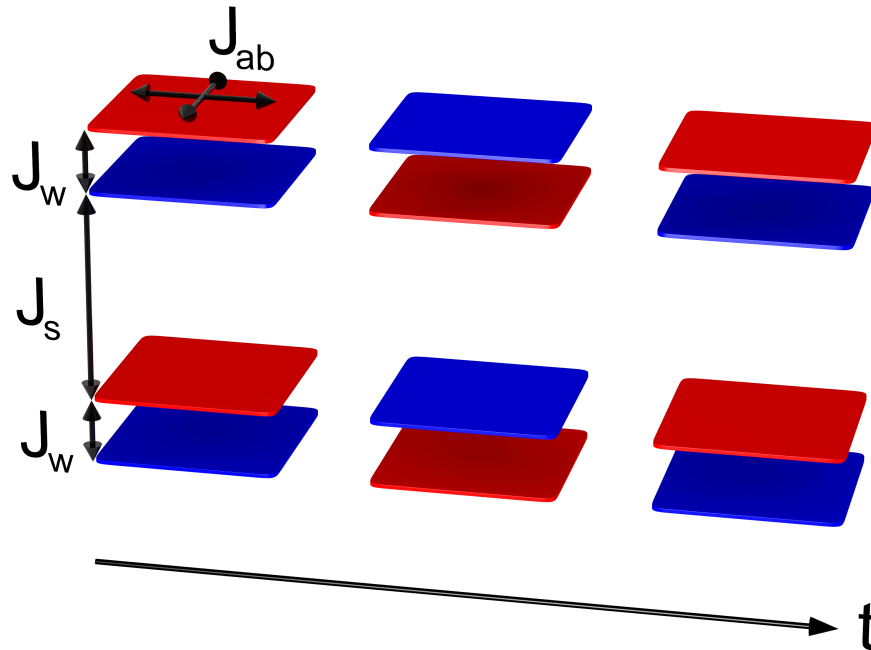


Arxiv:1406.3609

Robert Höppner, Beilei Zhu, Tobias Rexin, Ludwig Mathey

Driving

Modelled by a staggered potential, oscillating with frequency ω_m



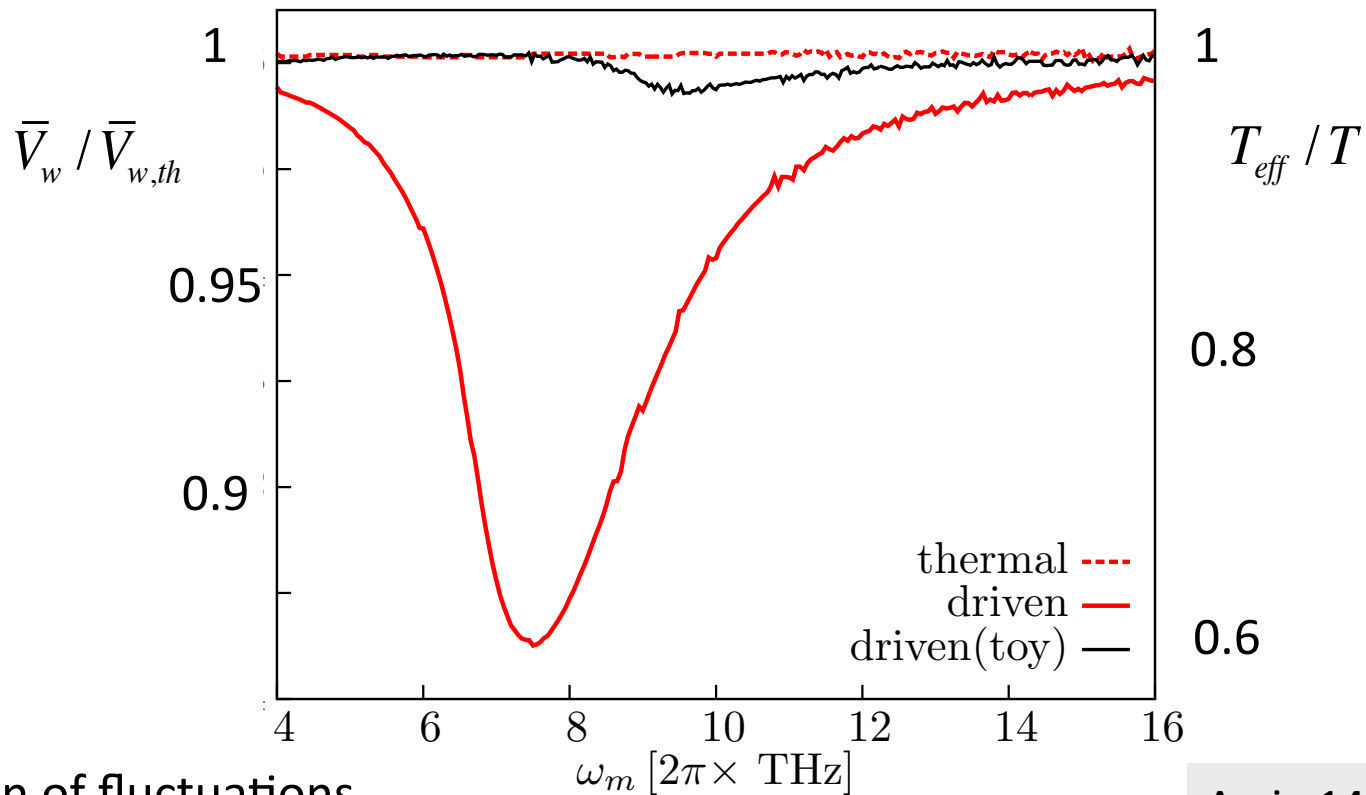
Current across the weak junctions: $j_w = 2J_w \sin \theta$

Arxiv:1406.3609

Robert Höppner, Beilei Zhu, Tobias Rexin, Ludwig Mathey

Frequency scan

Current fluctuations: $V_w(t) = \overline{(j_w(t))^2} - \left(\overline{j_w(t)}\right)^2$ $\xrightarrow{\text{time-averaged}}$ \bar{V}_w

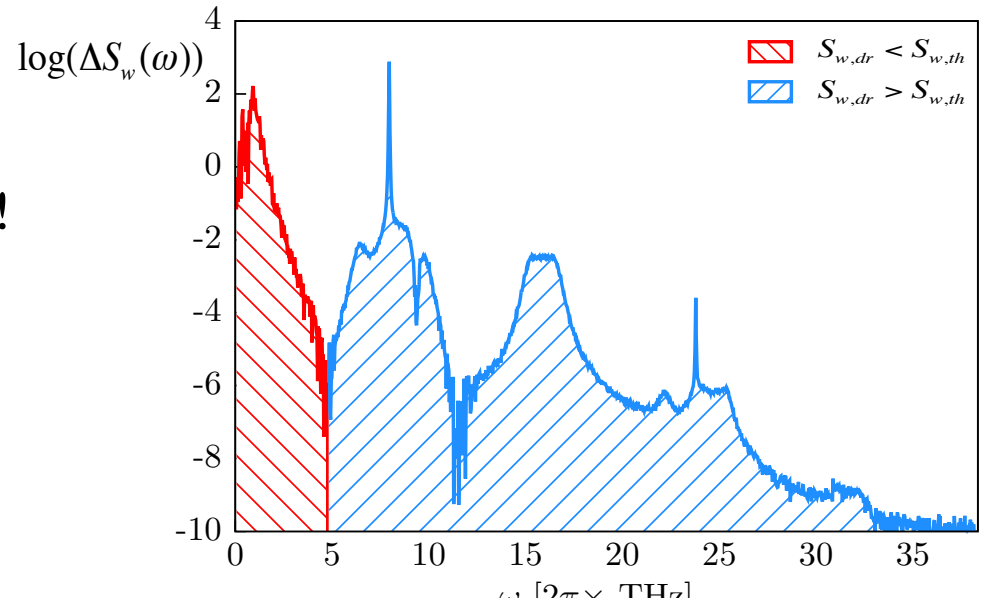
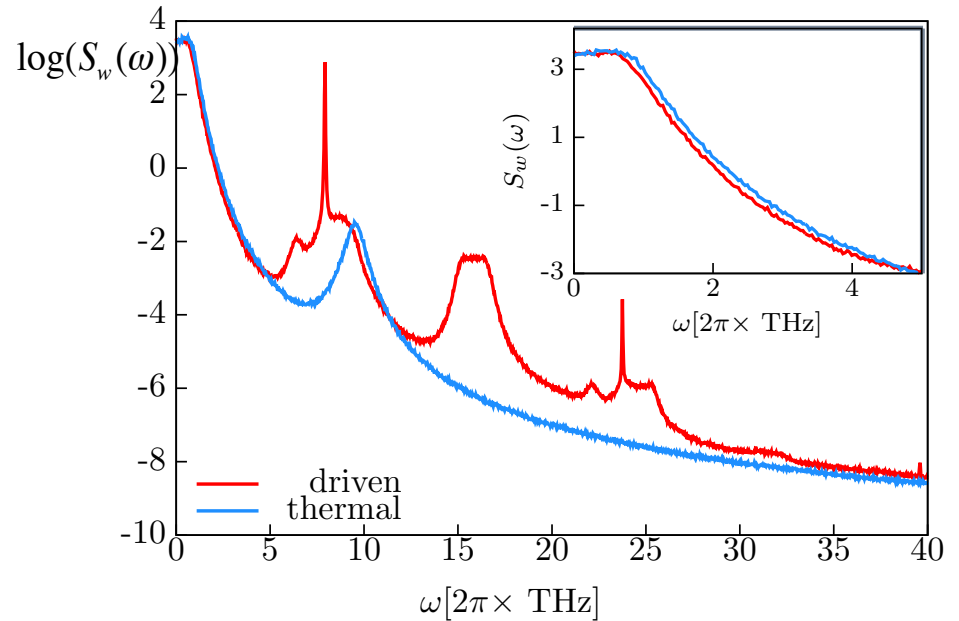


- ✓ Reduction of fluctuations
- ✓ Plasmon mode of the strong layers serve as an amplifier of the driving

Power spectrum

$$S_w(\omega) = \langle j_w(-\omega)j_w(\omega) \rangle$$

- ✓ Low-frequency modes suppressed
- ✓ High-frequency modes enhanced



Up-conversion of spectral weight!

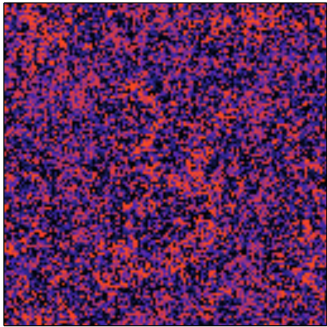
In-plane dynamics

Spatially resolved
fluctuations:

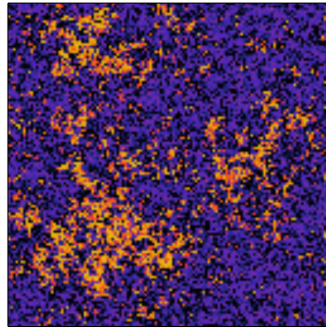
$$\left(j_w(\vec{r}, t) - \bar{j}_w(t) \right)^2$$

Spatial average $\bar{j}_w(t)$

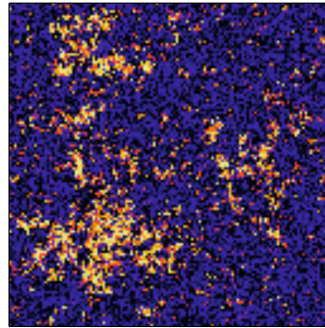
thermal



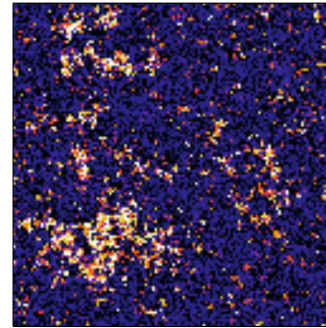
1/8 cycle



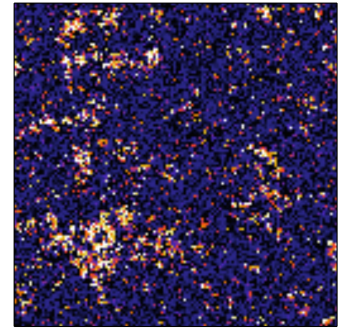
1/4 cycle



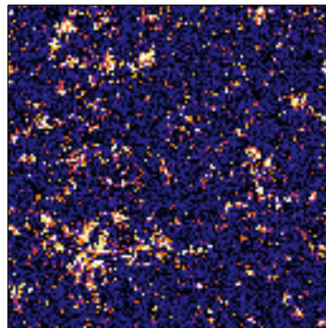
3/8 cycle



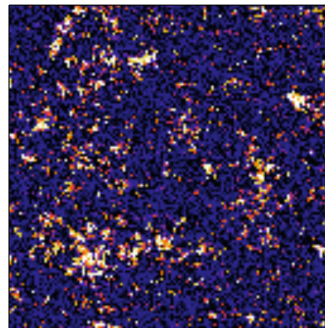
1/2 cycle



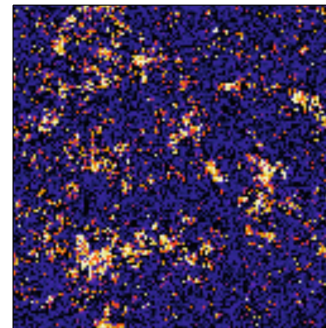
5/8 cycle



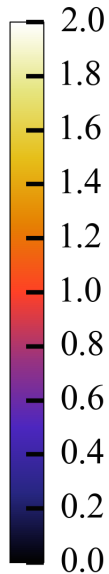
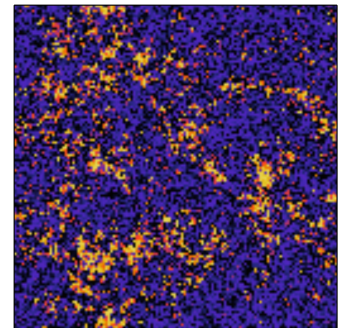
3/4 cycle



7/8 cycle



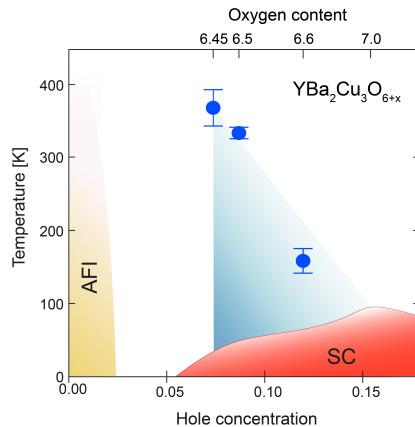
1 cycle



- ✓ Fluctuations reduced overall
- ✓ Local hotspots: short range fluctuations increased, long-range reduced

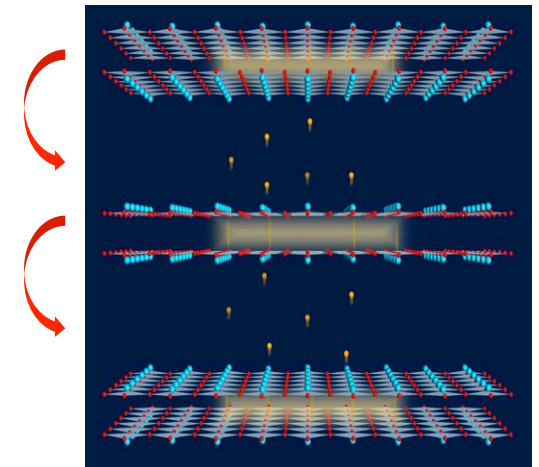
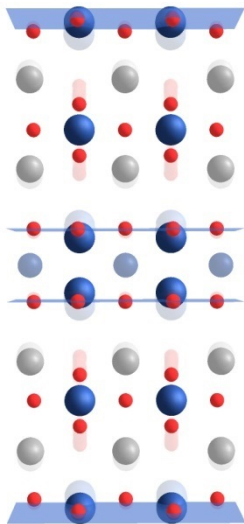
Arxiv:1406.3609

Summary



Transient Light Induced Superconductivity at Room Temperature in $\text{YBa}_2\text{Cu}_3\text{O}_{6+d}$

Coherence at low frequency appears at the expense of coupling within the bilayers



Non-linear lattice dynamics, CDW dynamics and dynamical stabilization offer hints to possible microscopic mechanisms

People

Mid-IR pump / THz probe experiments



Daniele Nicoletti



Cassi Hunt



Wanzheng Hu



Isabella Gierz

Mid-IR pump / X-Ray probe



Roman Mankowsky



Michael Först



Andrea Cavalleri

M. LeTacon

T. Loew

A. Frano

B. Keimer

MPI Stuttgart

A. Subedi

A. Georges

Paris

S. DHESI

DIAMOND

Light Source

J. Turner

D. Dakovski

M. Miniti

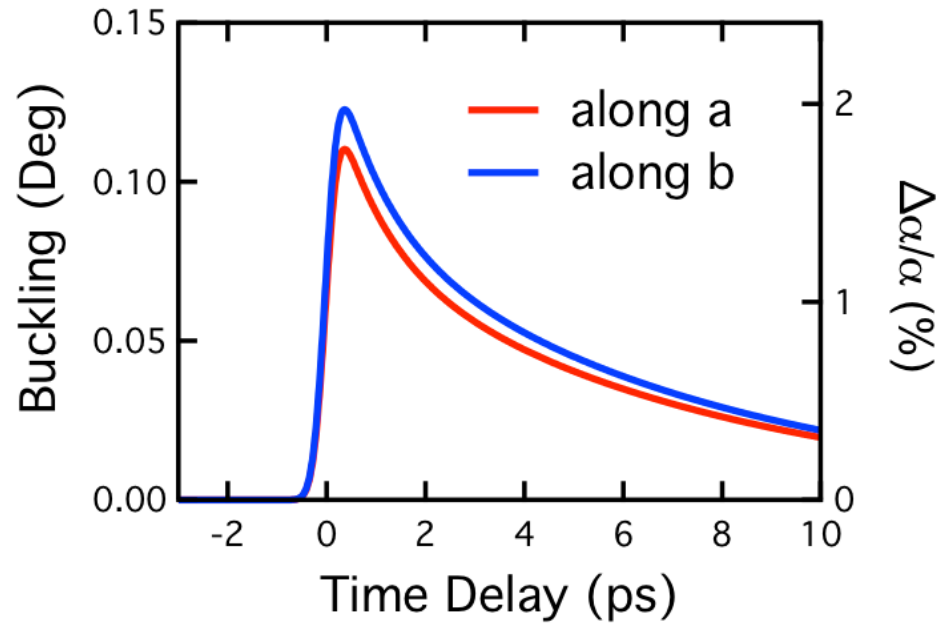
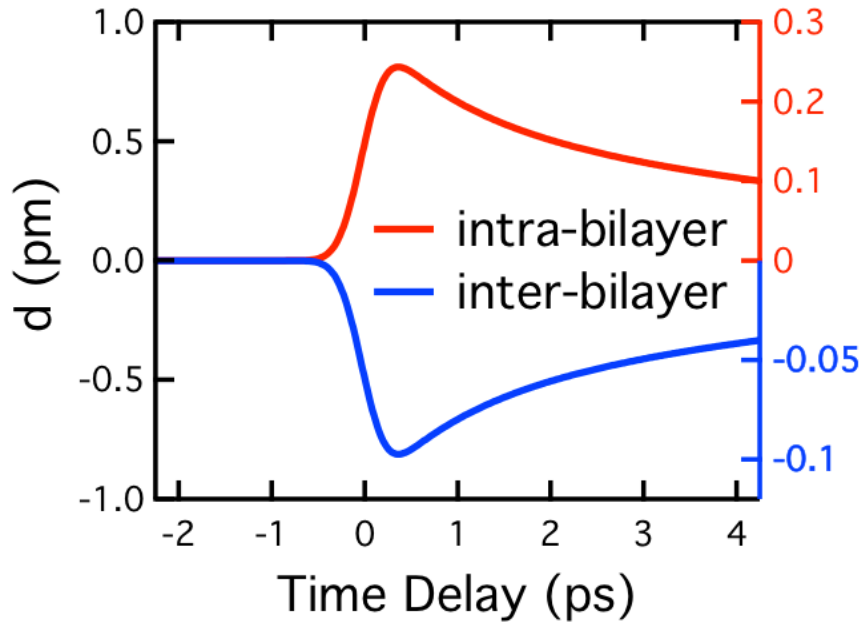
J. Robinson

Thank you for your attention!

A decorative graphic in the bottom-left corner consisting of numerous overlapping, semi-transparent, light blue shapes that resemble stylized leaves or petals, creating a layered, organic effect.

- Change in Bilayer Distance

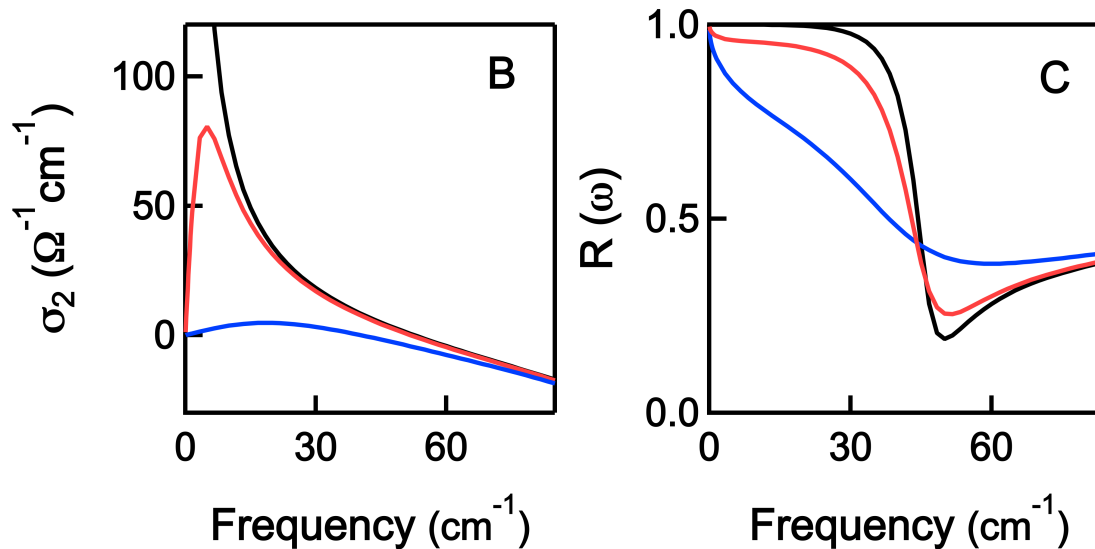
Change in Buckling



Transient Perfect Metal vs Superconductor

Transient metal:

- Finite scattering time τ (as large as the lifetime of the transient state)
- Can not be distinguished from a SC at frequencies above $1/\tau$



Induced coherence above T_c

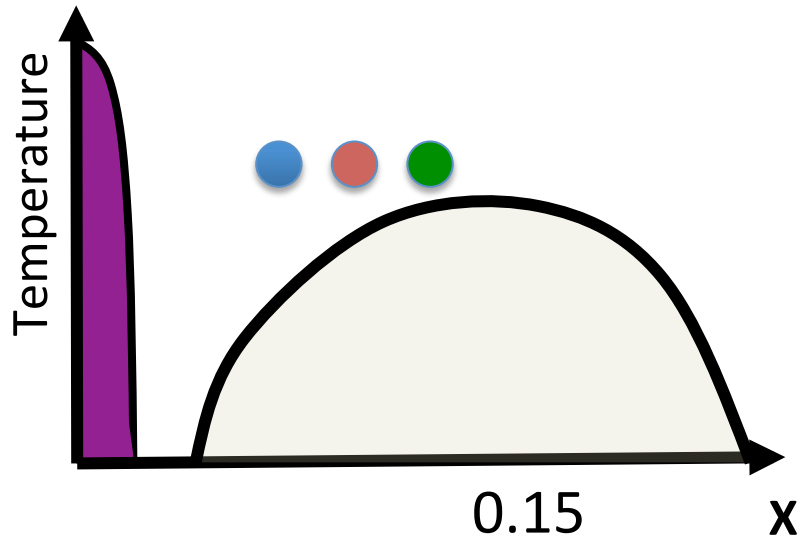
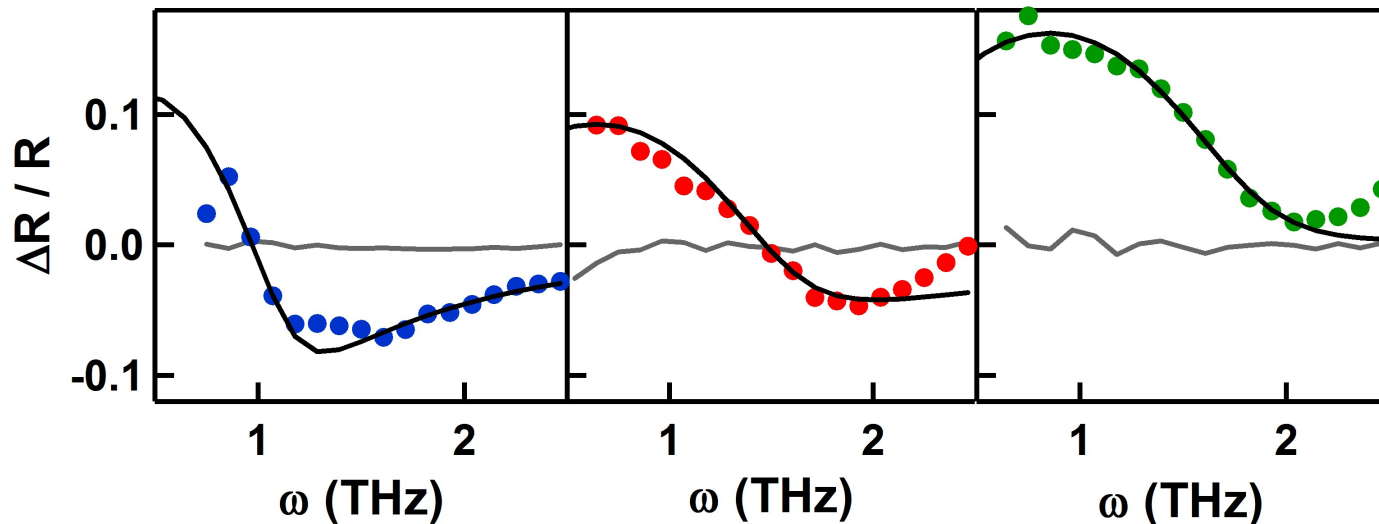


Photo-induced low frequency JPR

YBCO 6.45

YBCO 6.5

YBCO 6.6



Induced coherence above T_c

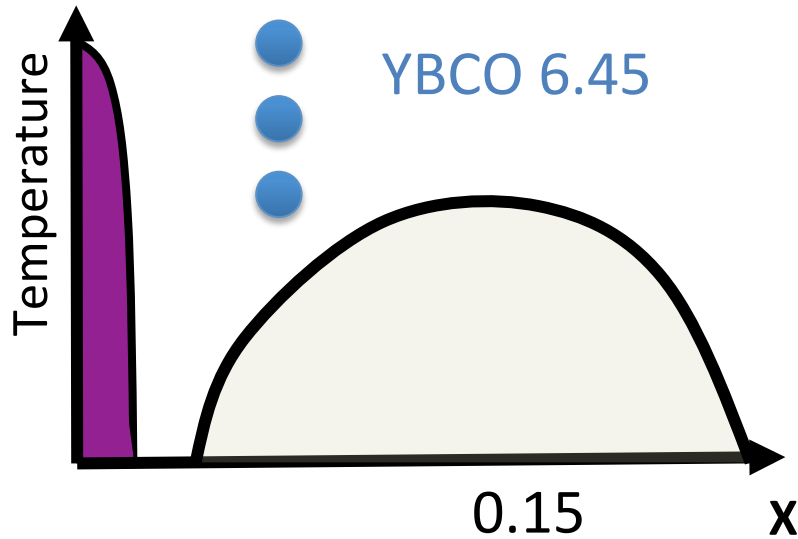
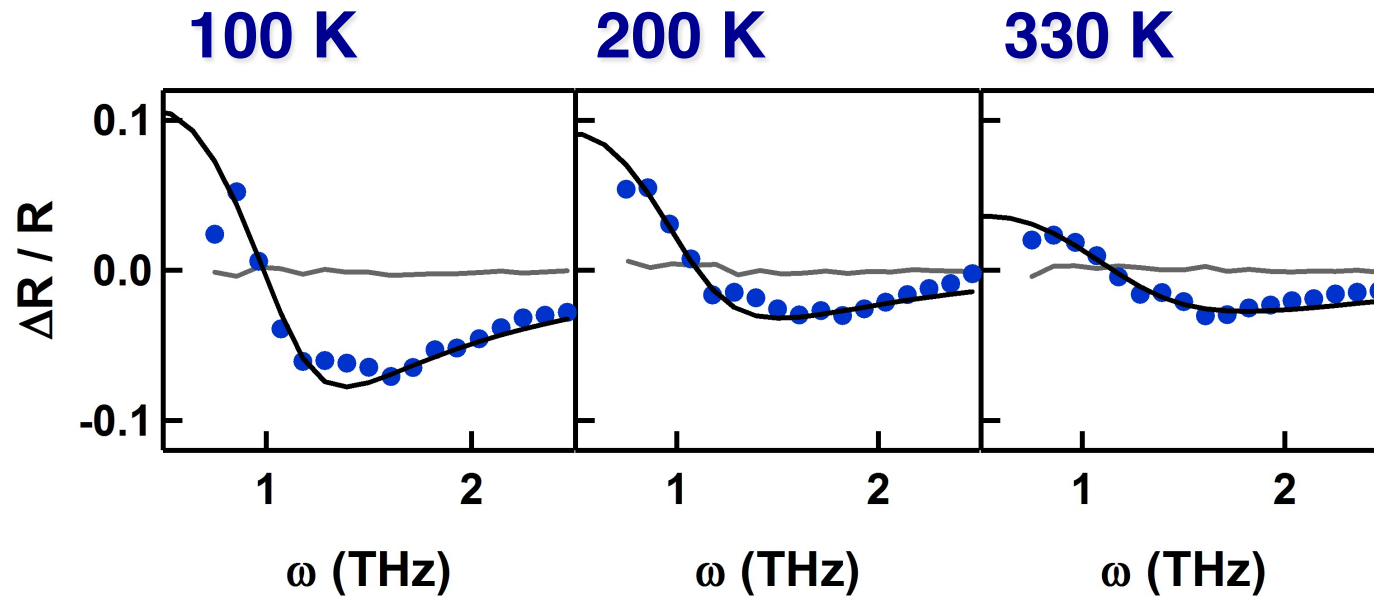
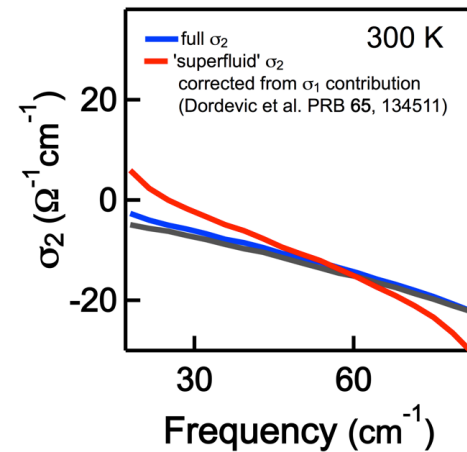
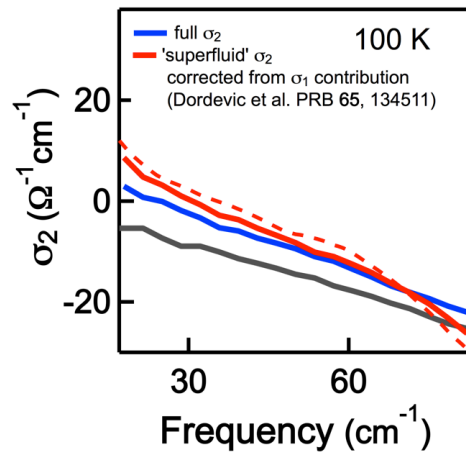
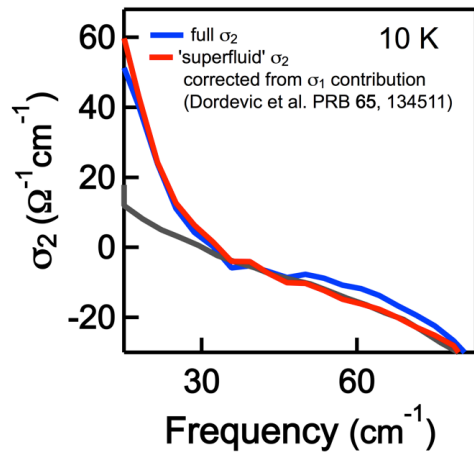
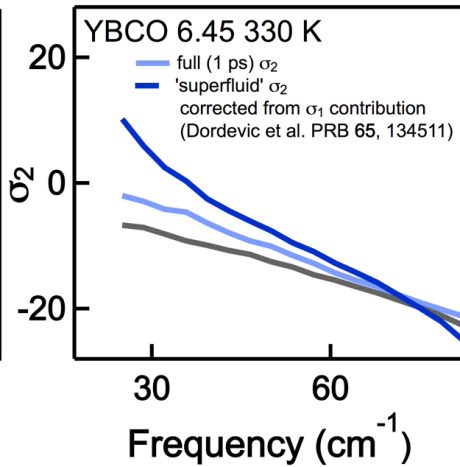
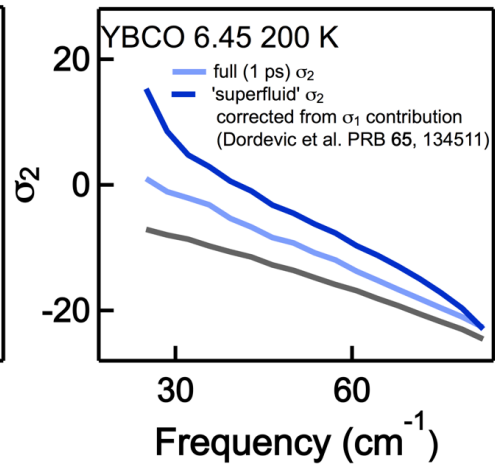
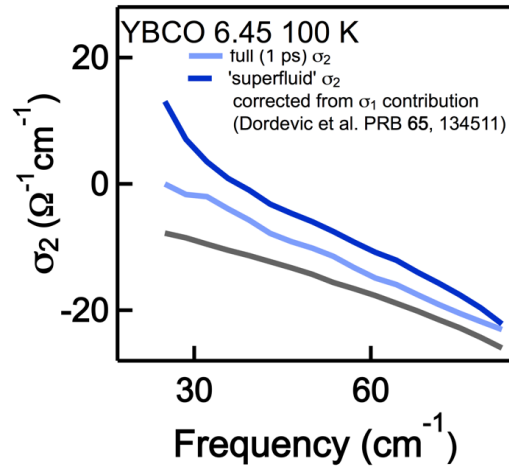


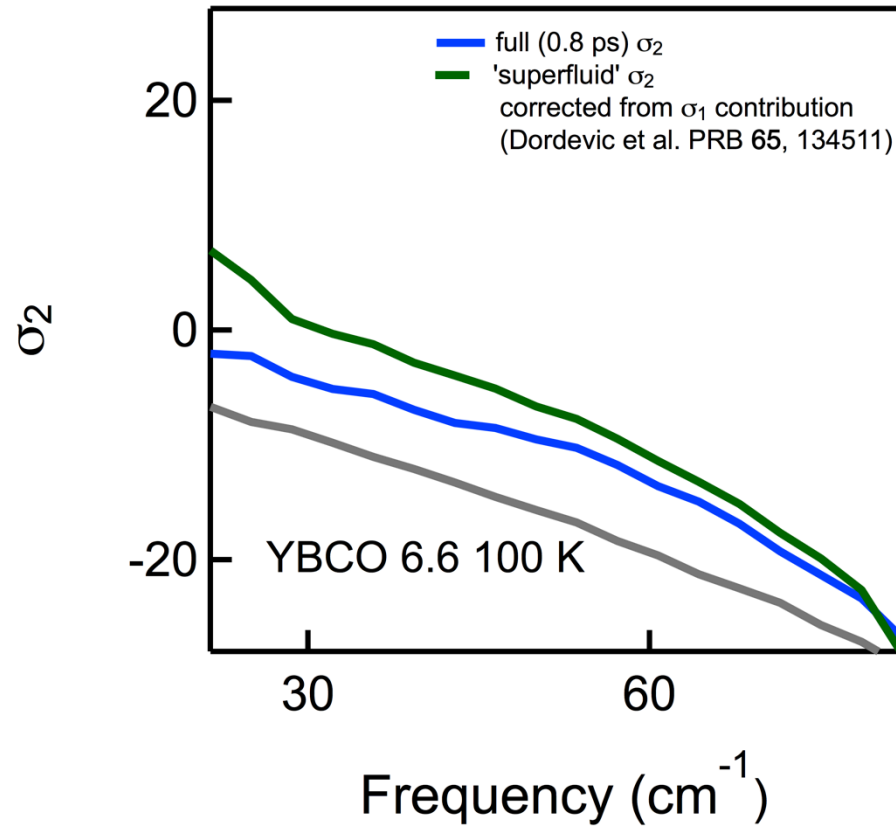
Photo-induced low frequency JPR



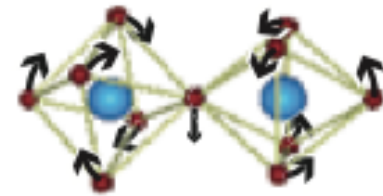
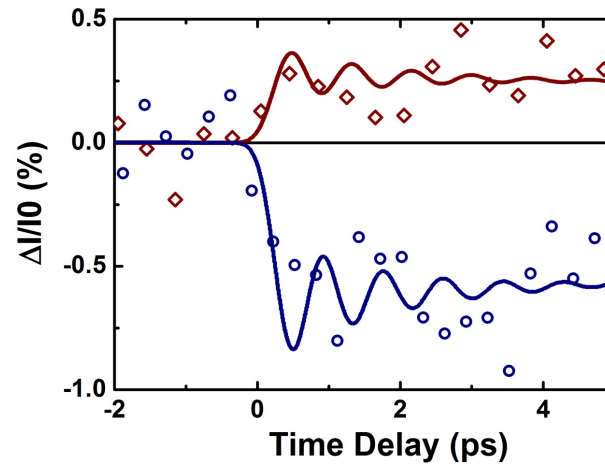
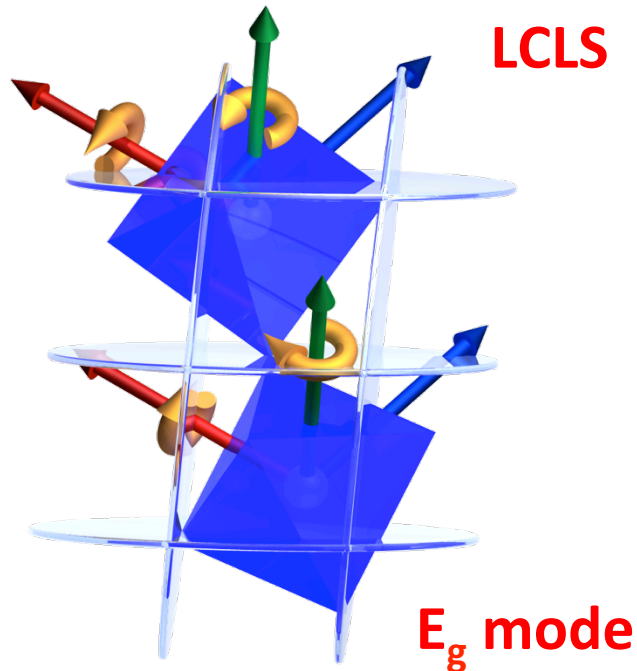
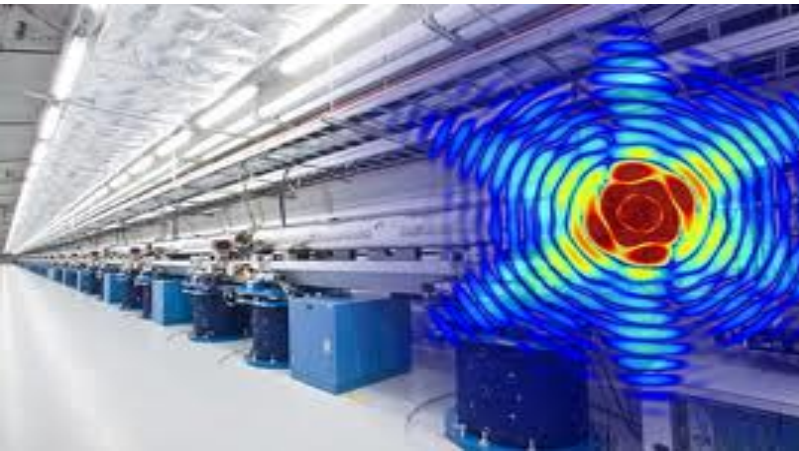
Correction of sigma2



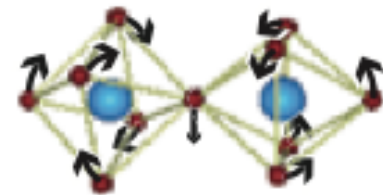
Correction of sigma2



Displacive Eg excitation: (012) up and (201) down



(012)

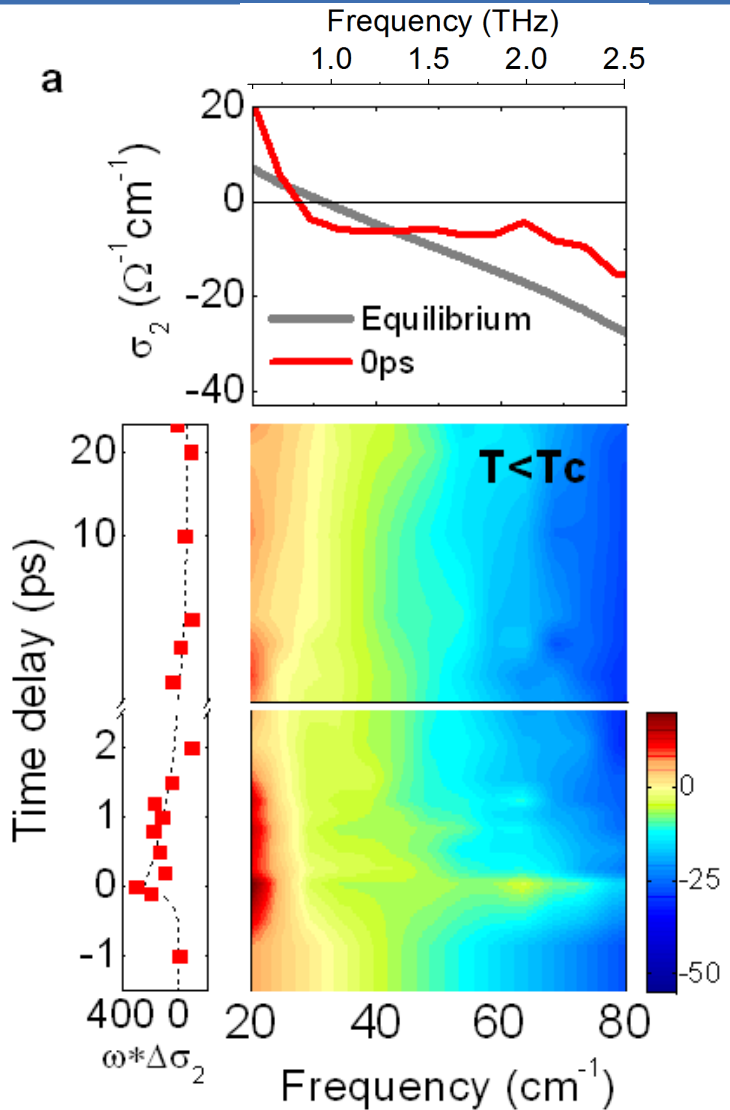


(201)

M. Foerst , R. Mankowski et al.

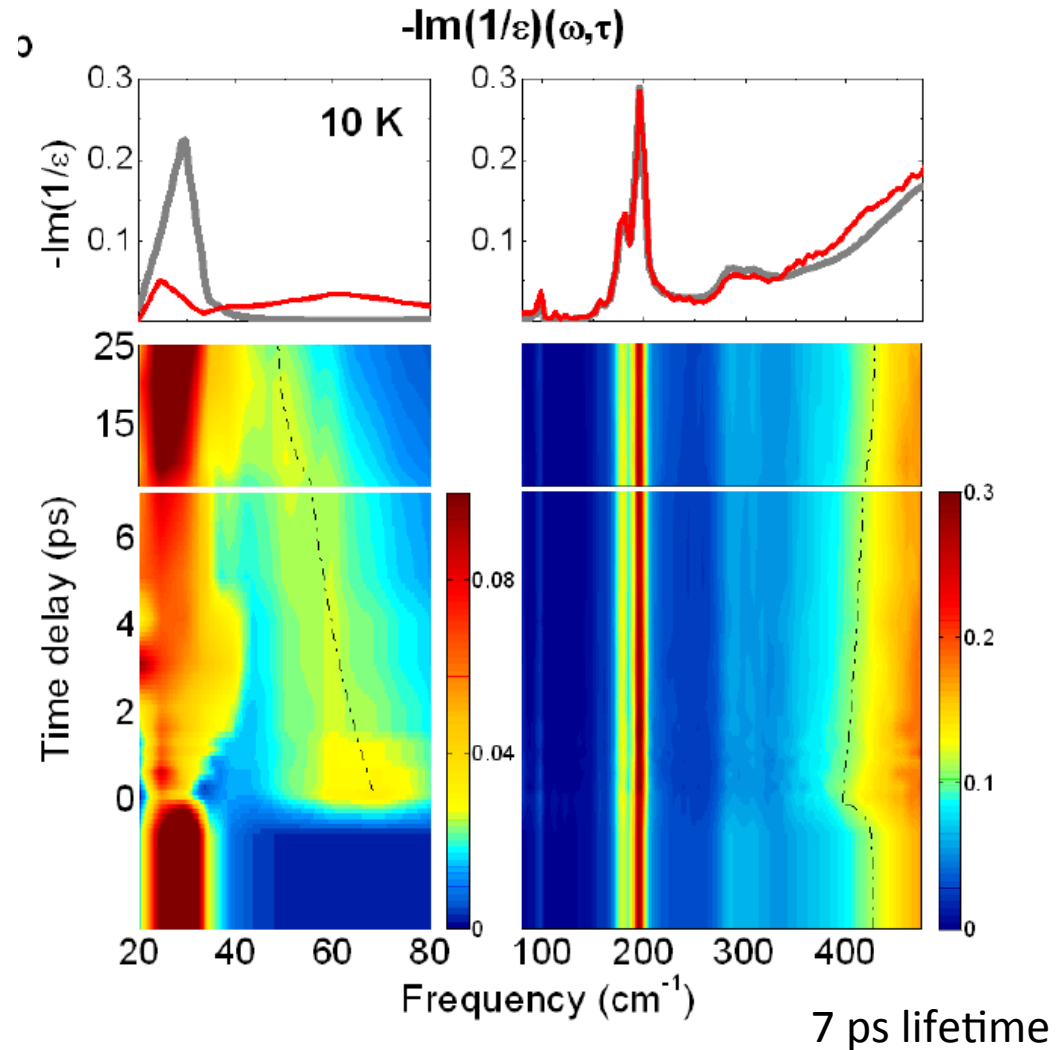
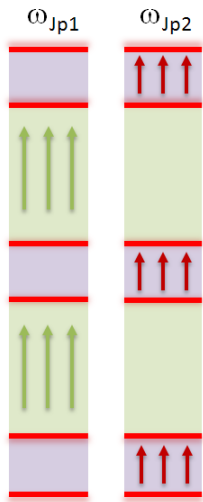
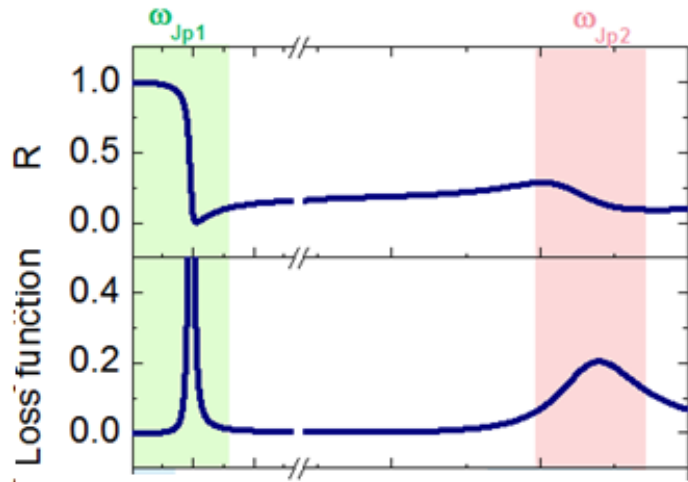
Below Tc: Enhancement of Superconductivity

Enhancement of divergent σ_2

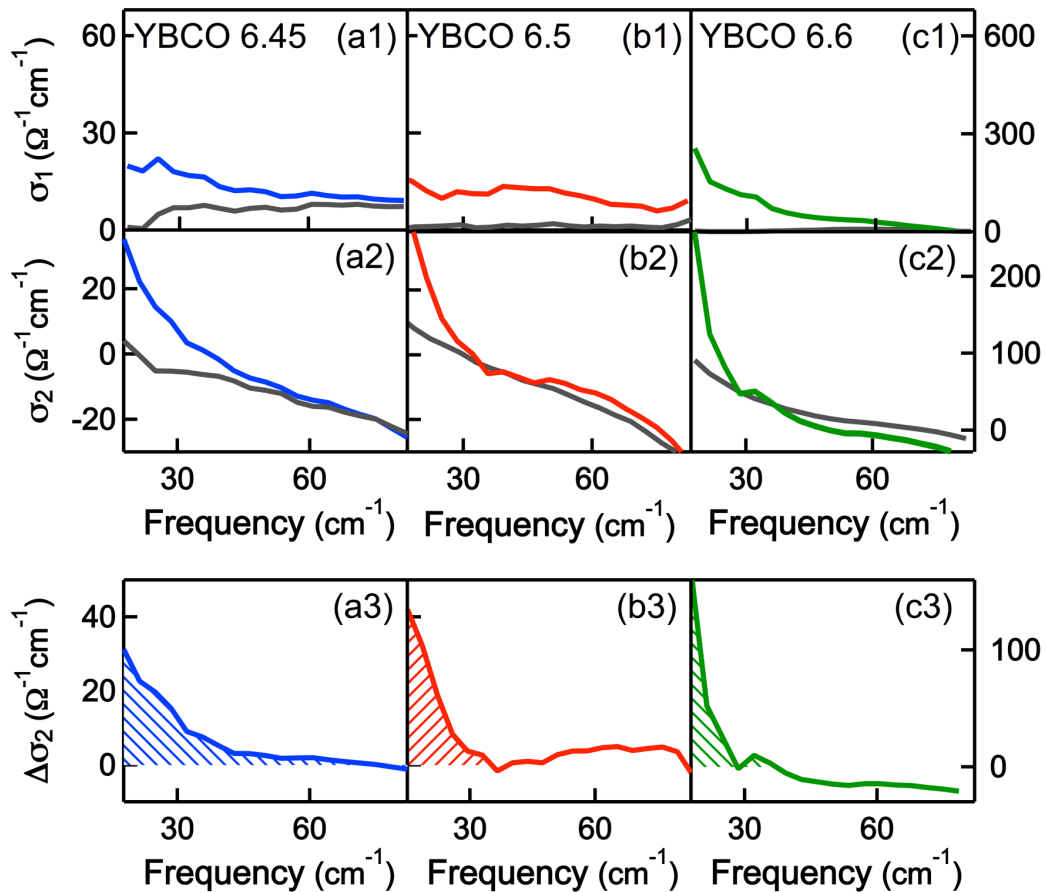


1 THz = 33 cm^{-1}

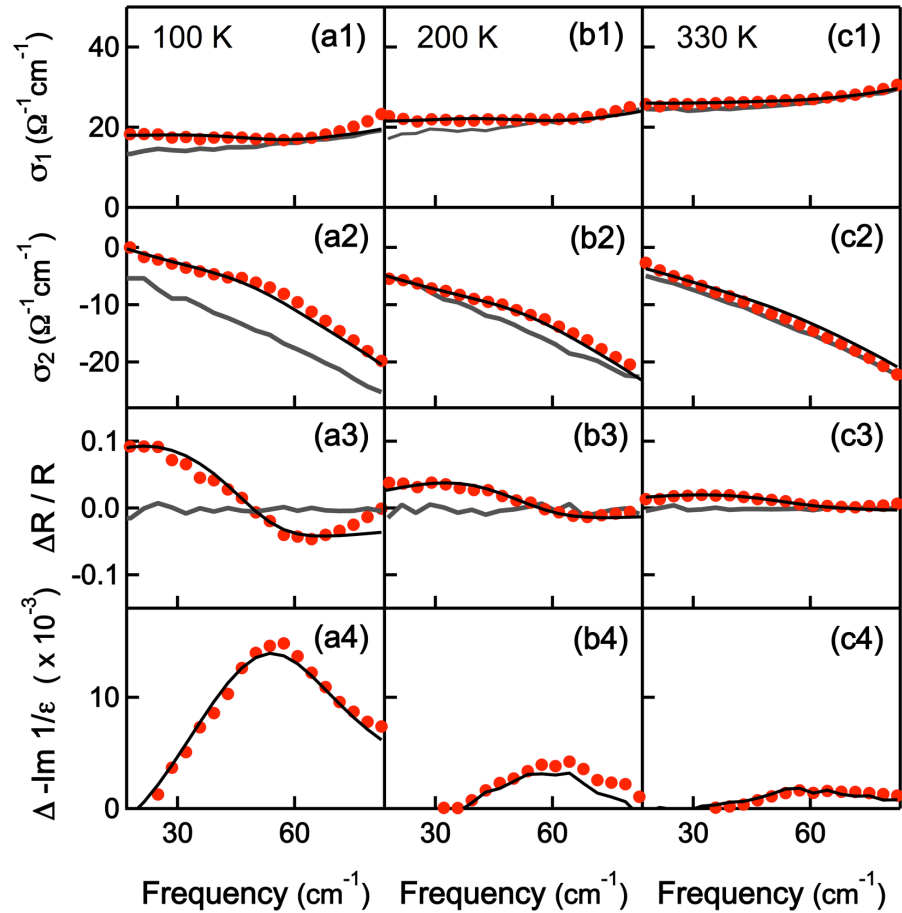
Two longitudinal Josephson plasma modes



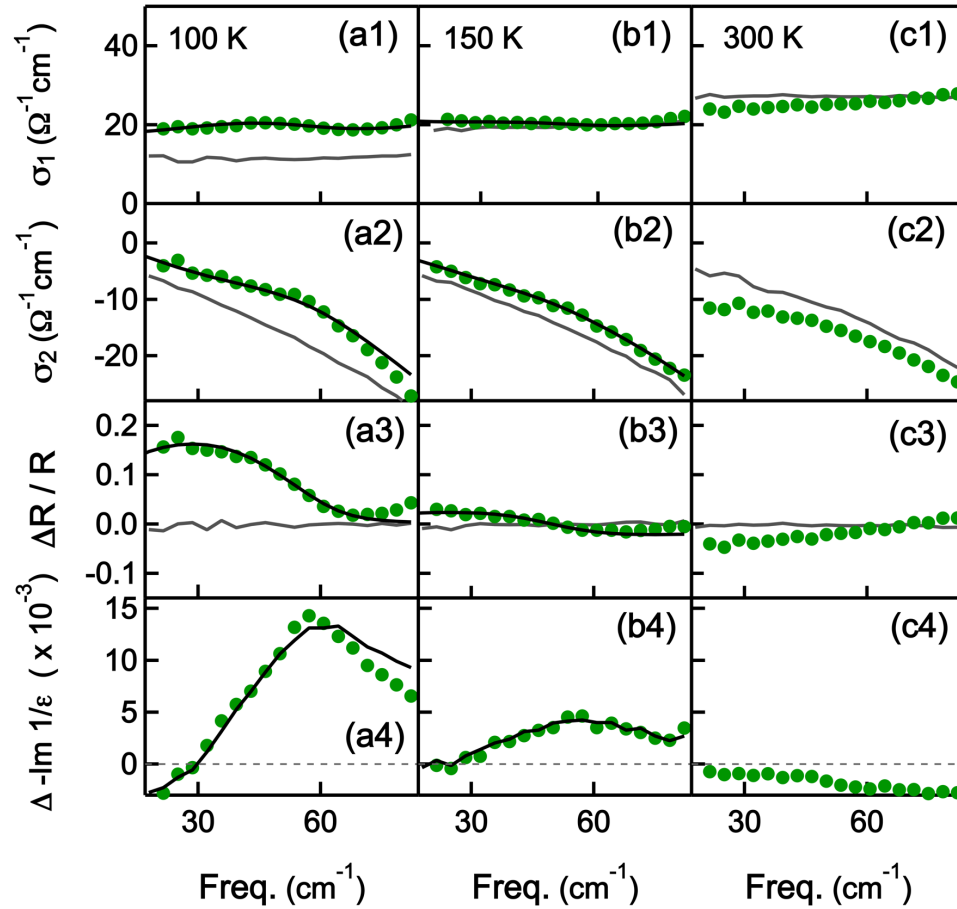
YBCO below T_c



YBCO 6.5

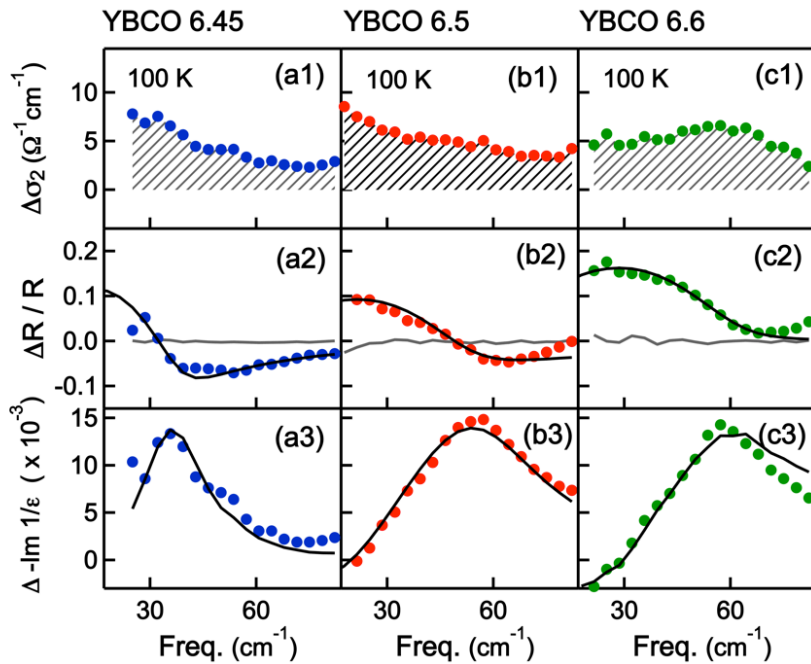


YBCO 6.6

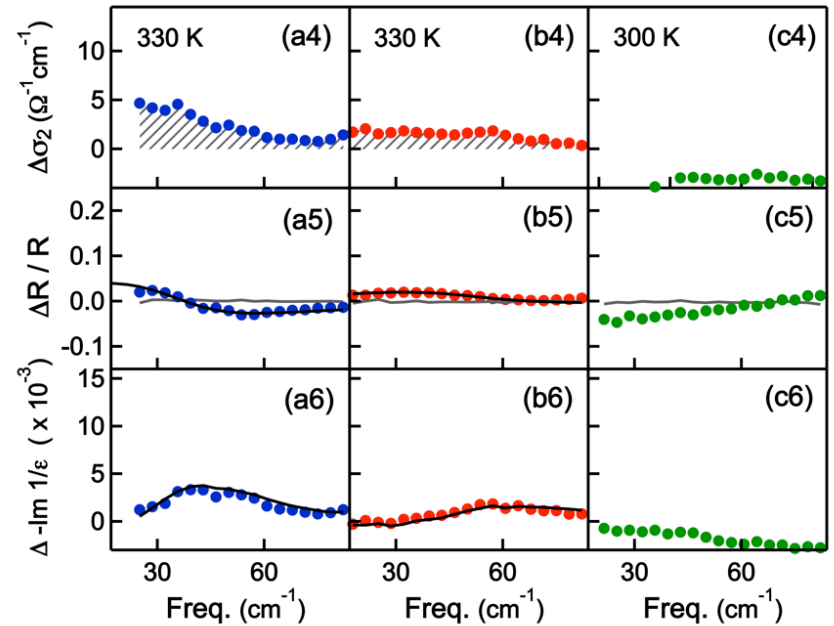


YBCO above T_c

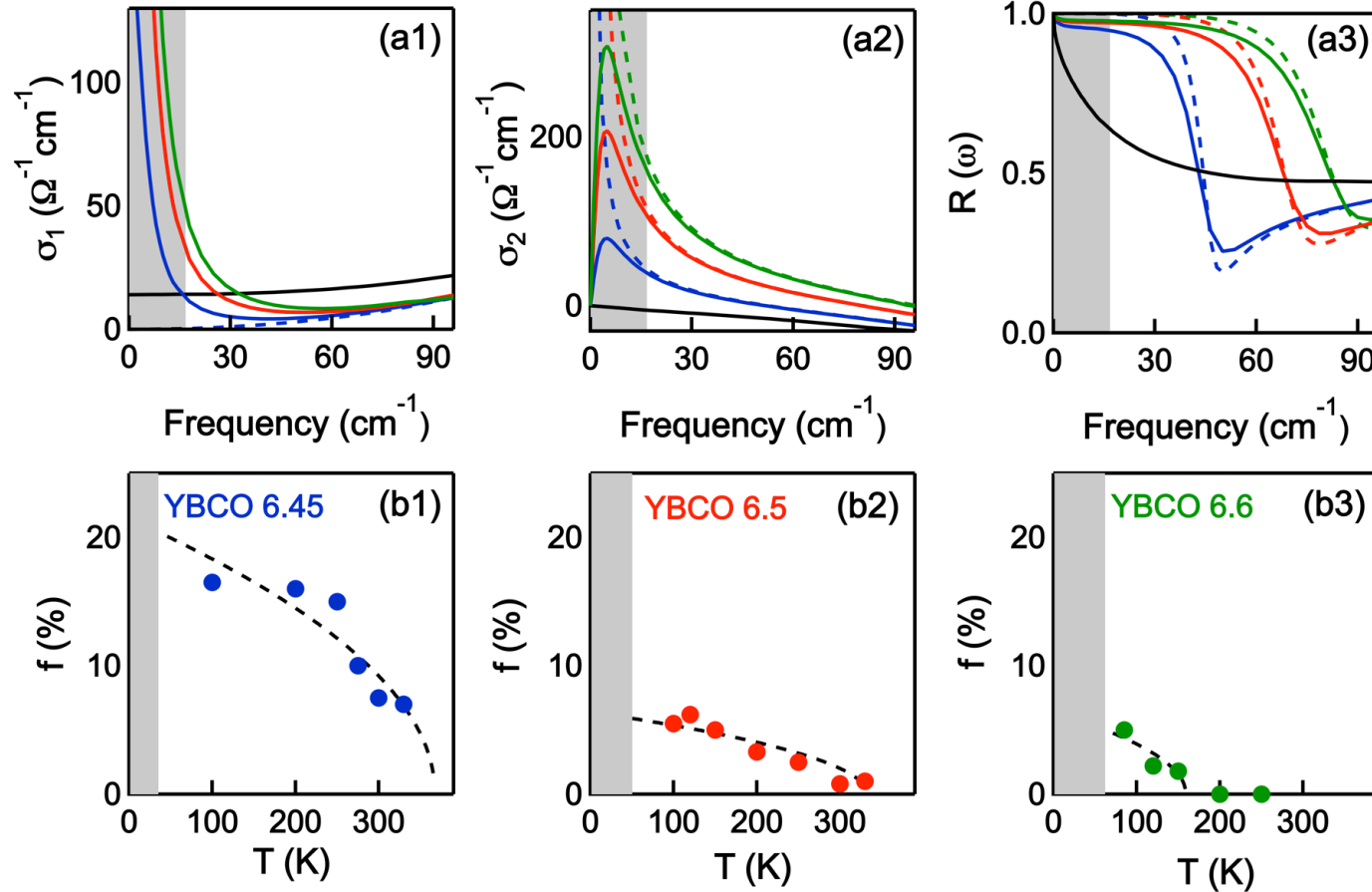
100 K



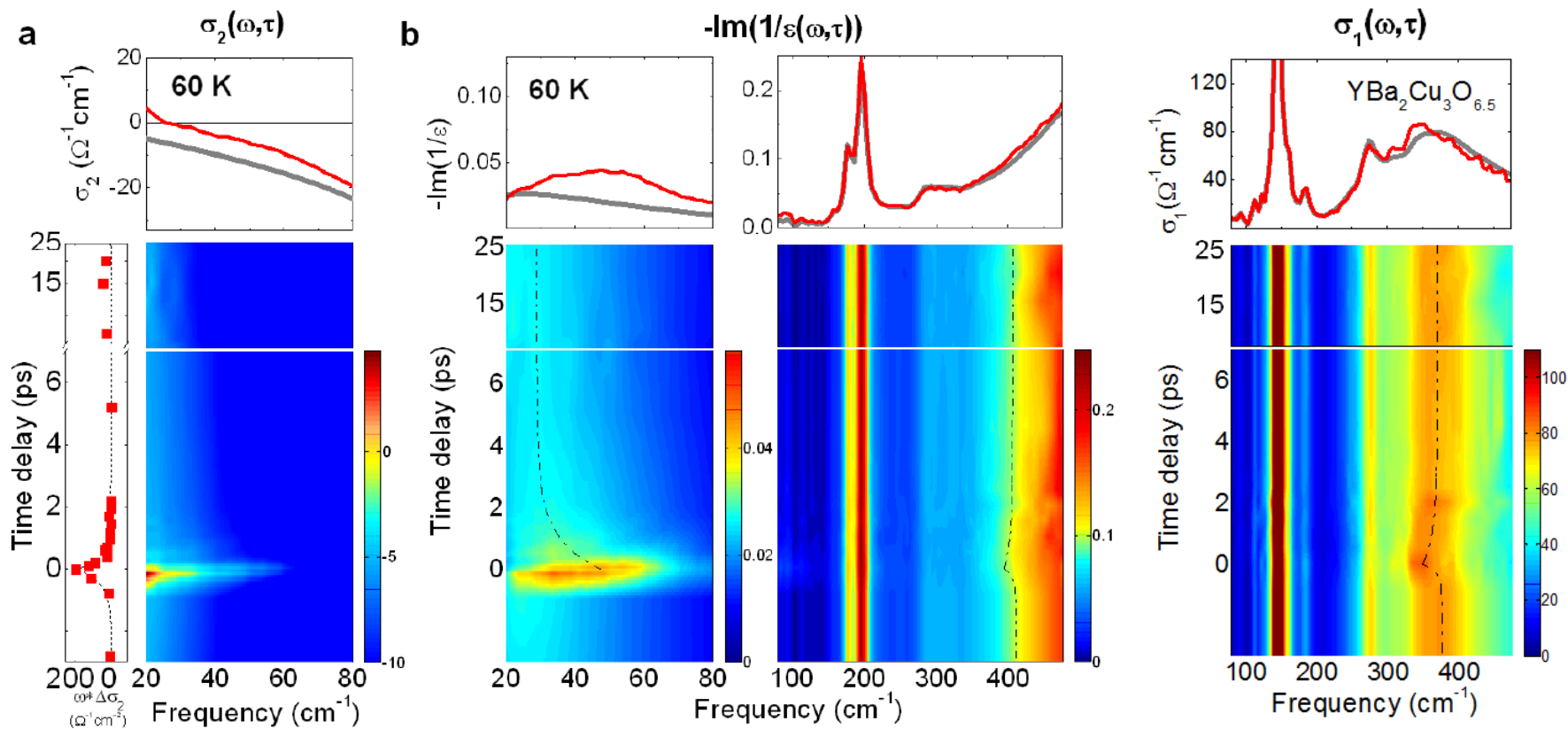
330 K / 300 K



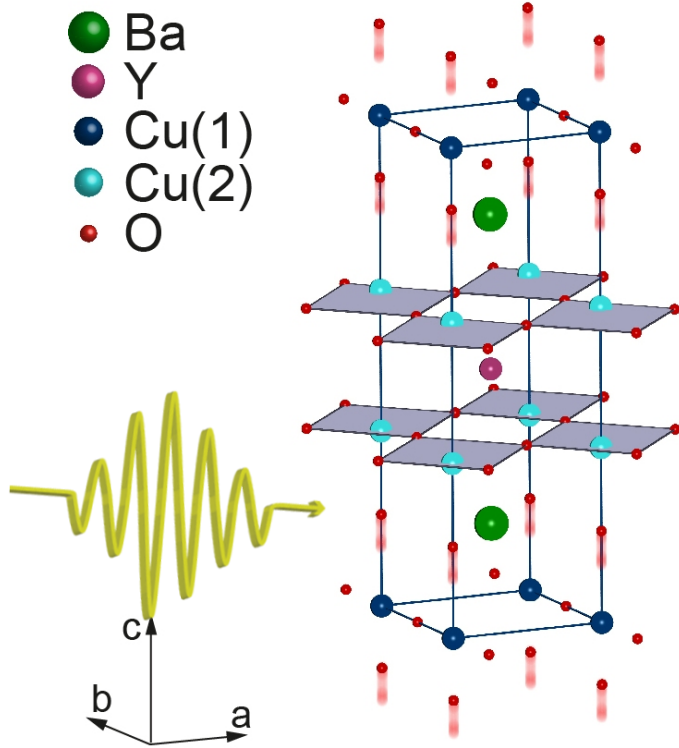
Extracted SC properties



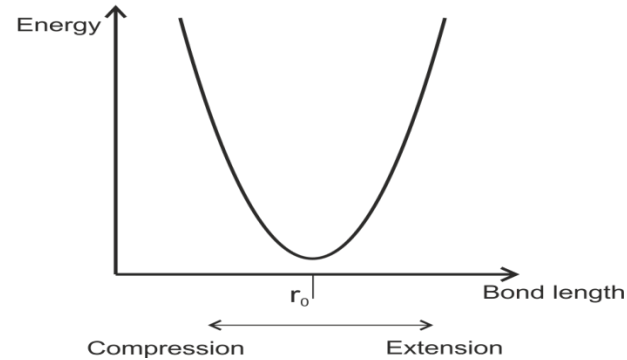
Above T_c



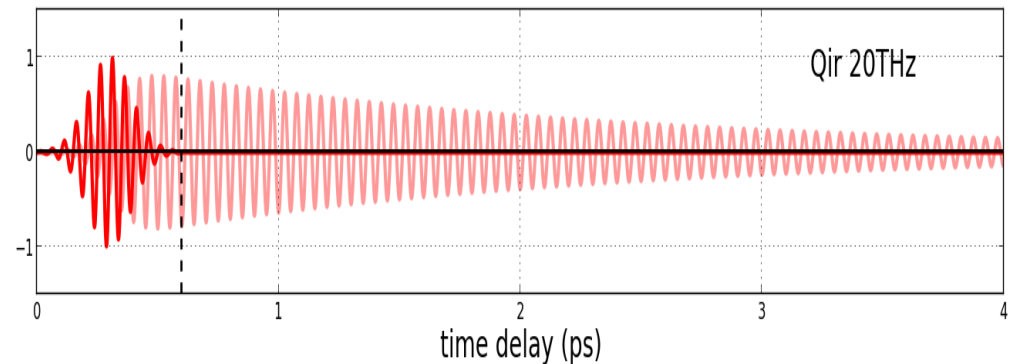
Linear response: no average displacement



$$V = \frac{1}{2} \mu_{IR} \omega_{IR}^2 Q_{IR}^2$$



$$\ddot{Q}_{IR} + 2\gamma\dot{Q}_{IR} + \omega_{IR}^2 Q_{IR} = f(t)$$

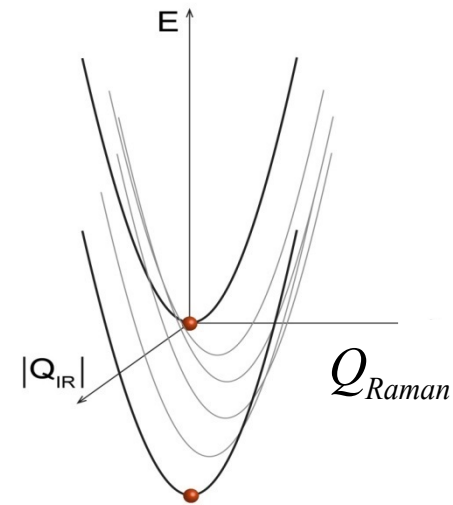


Anharmonic coupling

$$V = \frac{1}{2} \mu_{IR} \omega_{IR}^2 Q_{IR}^2 + NAQ_{IR}^2 Q_2$$

$$\mu_{IR} (\ddot{Q}_{IR} + 2\gamma\dot{Q}_{IR} + \omega_{IR}^2 Q_{IR}) = Q_{IR} Q_2 + f(t)$$

$$\mu_2 (\ddot{Q}_2 + 2\gamma\dot{Q}_2 + \omega_2^2 Q_2) = A Q_{IR}^2$$

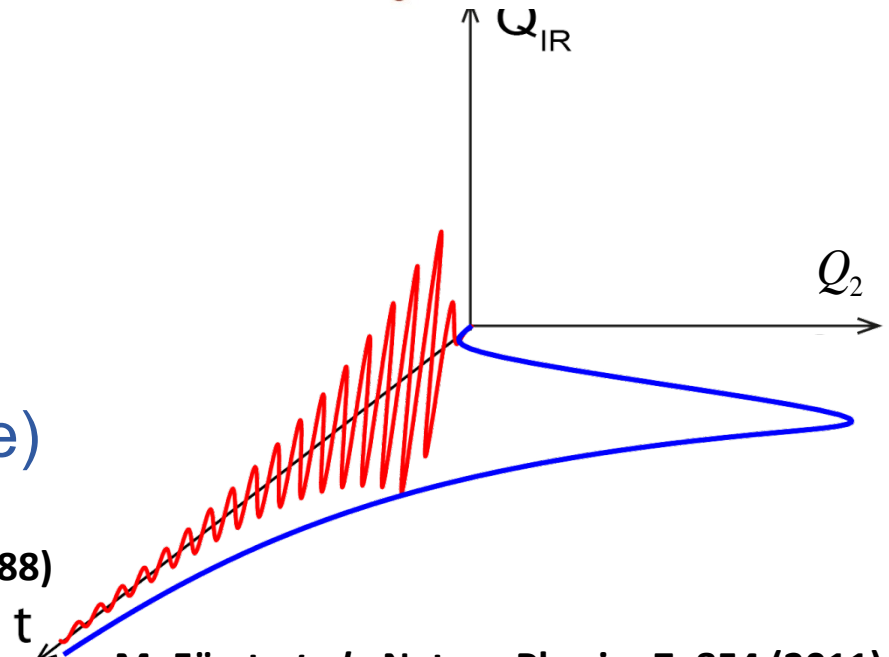


Q_{IR} (B_{1u})

$Q_{IR}^2 Q_2$ not zero

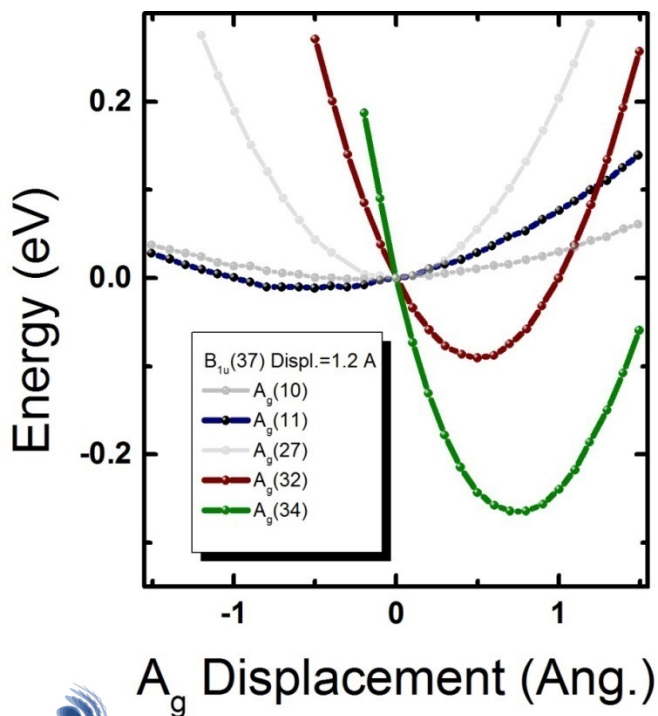
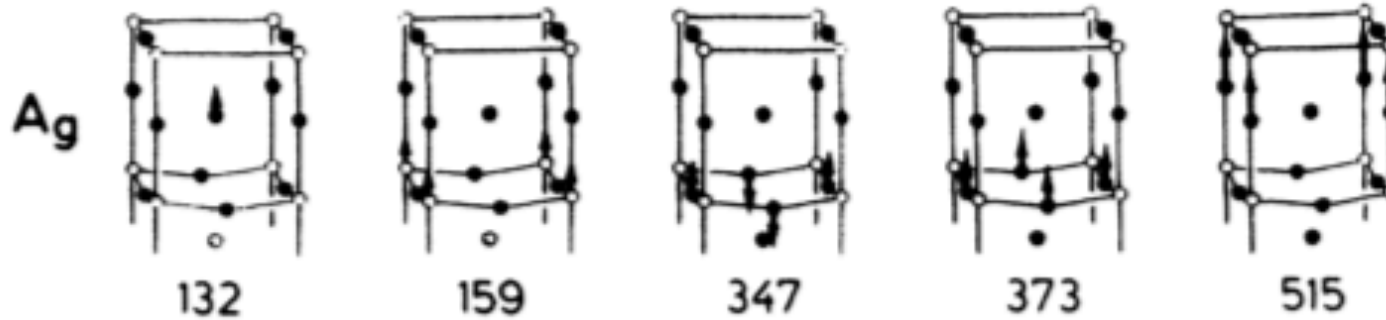
only for Q_2 (A_g , Raman active)

R. Liu et al., *Phys. Rev. B* 37, 7971 (1988)



M. Först et al., *Nature Physics* 7, 854 (2011)

3 coupled Raman modes



DFT Calculations

Frozen B_{1u} distortion (1 Å)

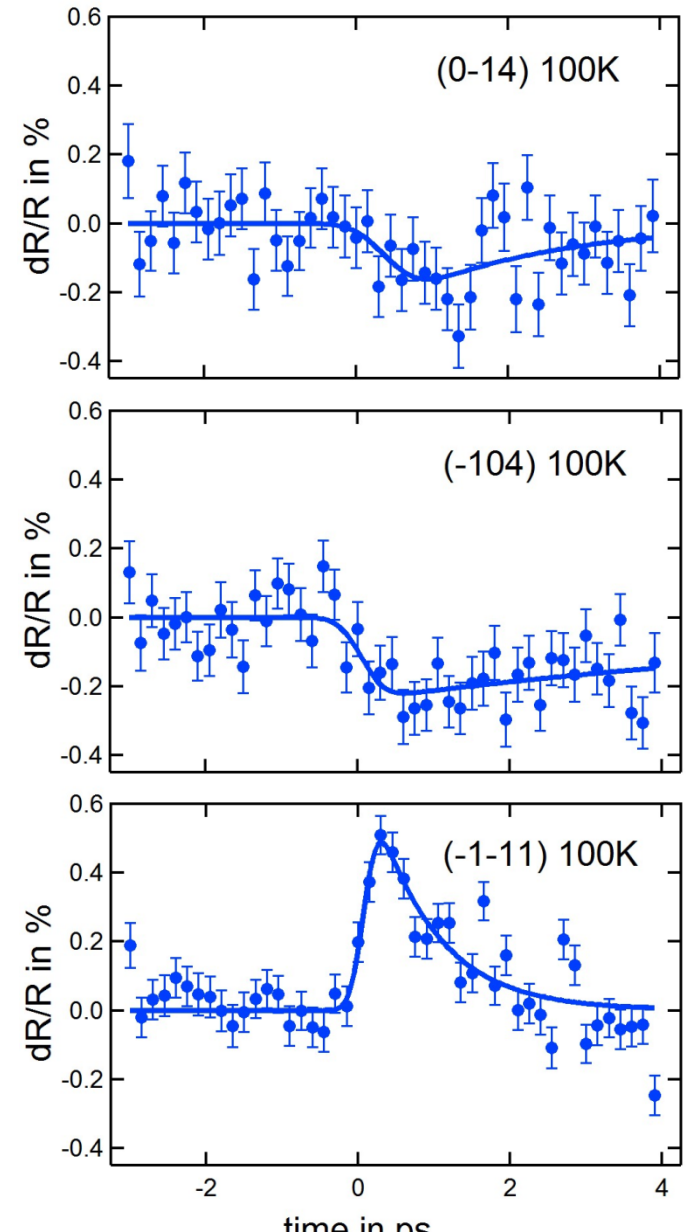
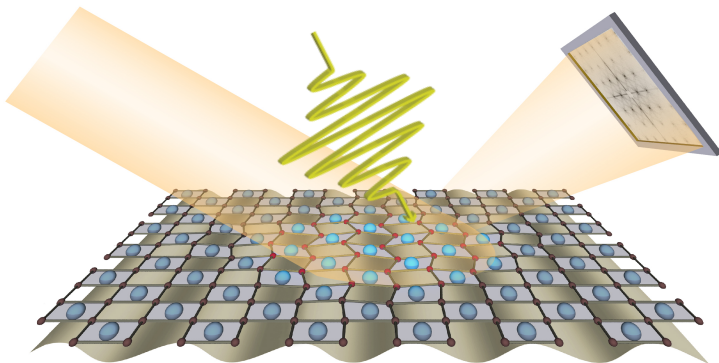
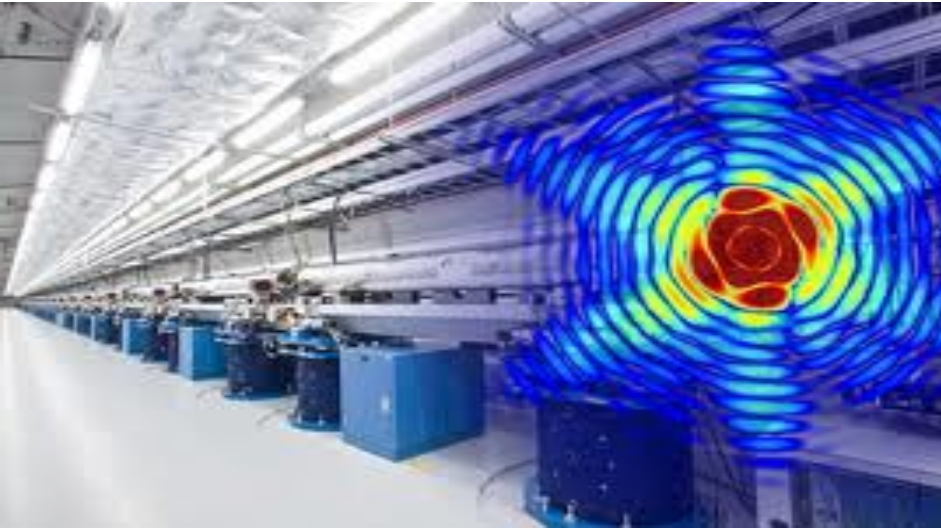
3 coupled modes (out of 5)

Alaska Subedi – Antoine Georges (Ecole Polytechnique)

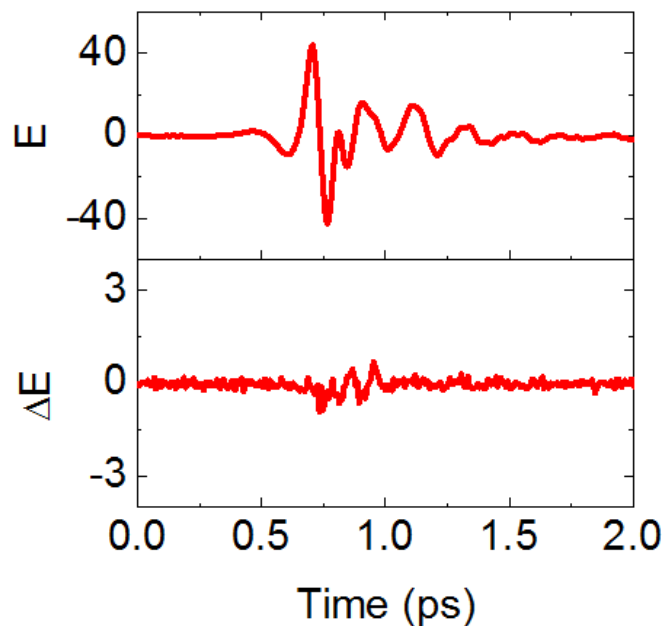
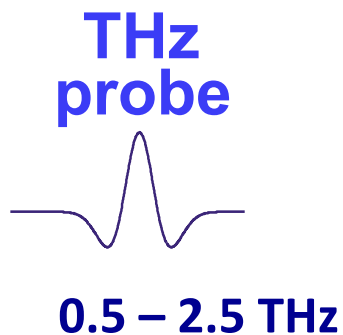
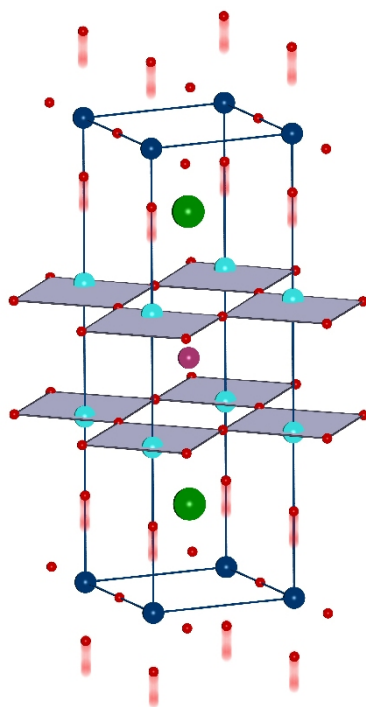
Michael Fechner - Nicola Spaldin (ETH)

How much is each mode moving ?

LCLS: Femtosecond Crystallography

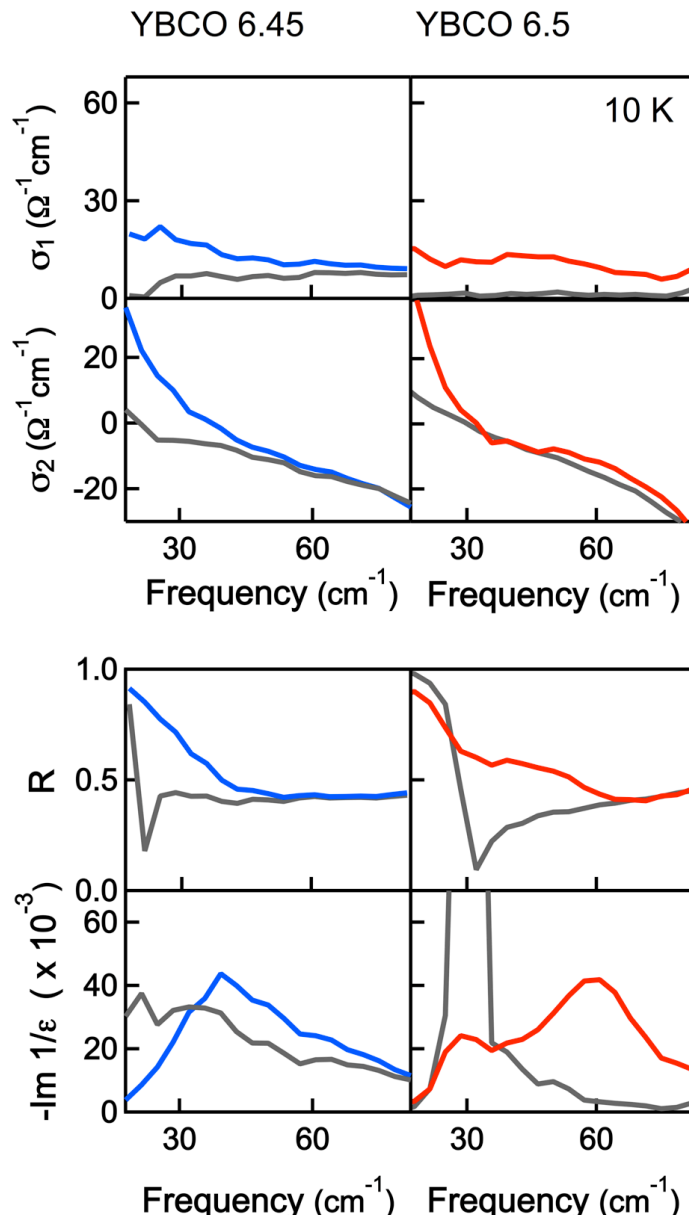


Direct Excitation of the Lattice in YBCO



$$\frac{\Delta \tilde{E}_R(\omega, \tau)}{\tilde{E}_R(\omega)} = \frac{\tilde{r}(\omega, \tau) - \tilde{r}_0(\omega)}{\tilde{r}_0(\omega)}$$

Enhanced coherence below T_c



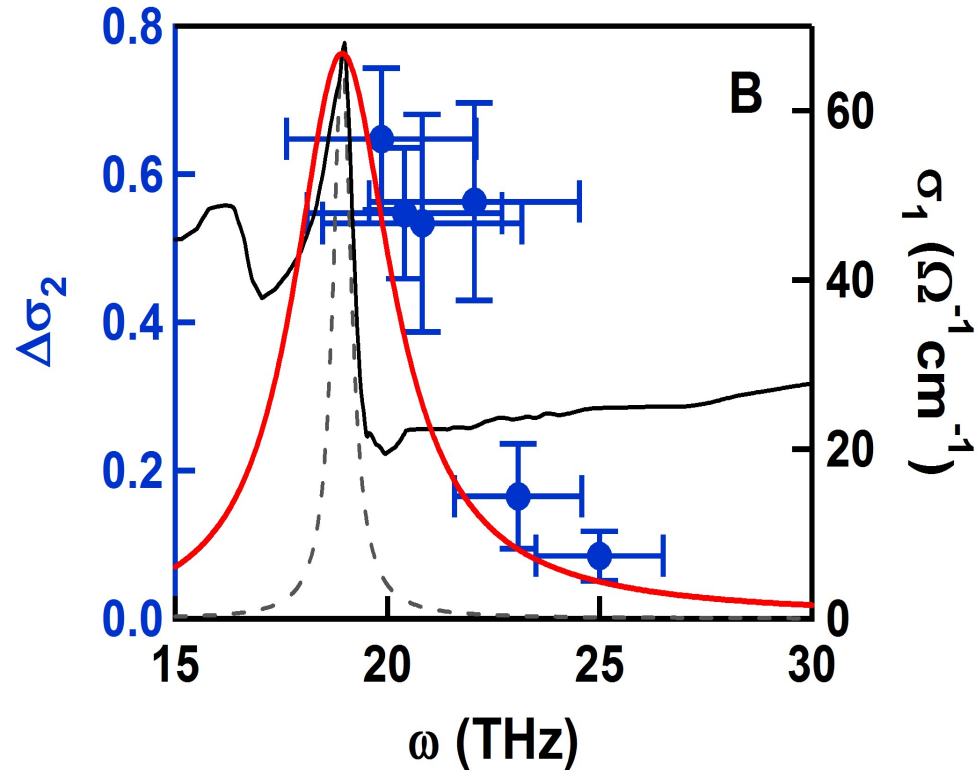
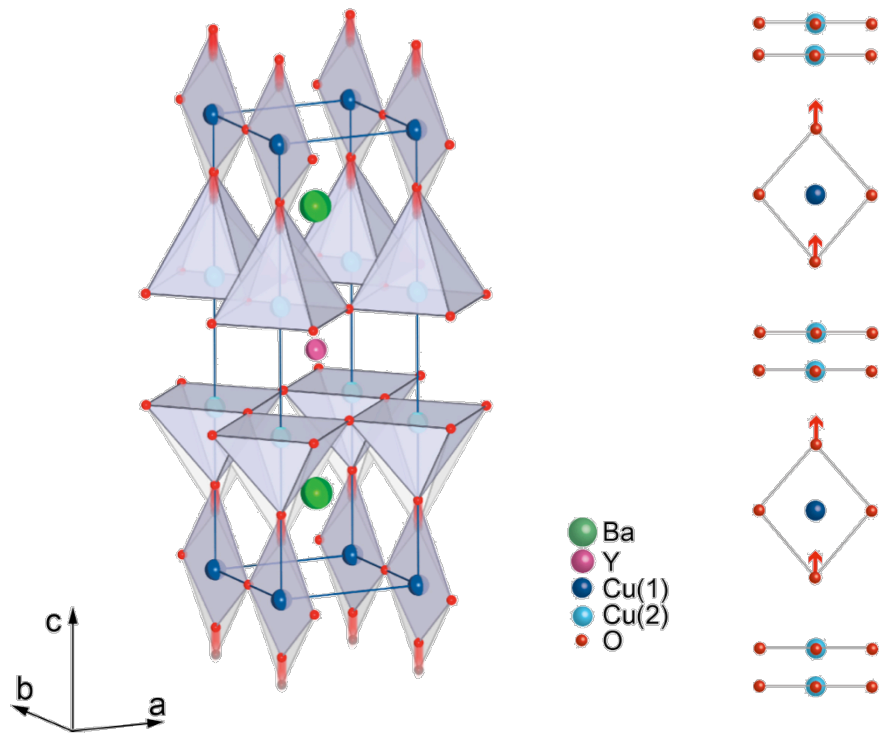
Increase of superconductivity

enhanced σ_2

Blue shift of the

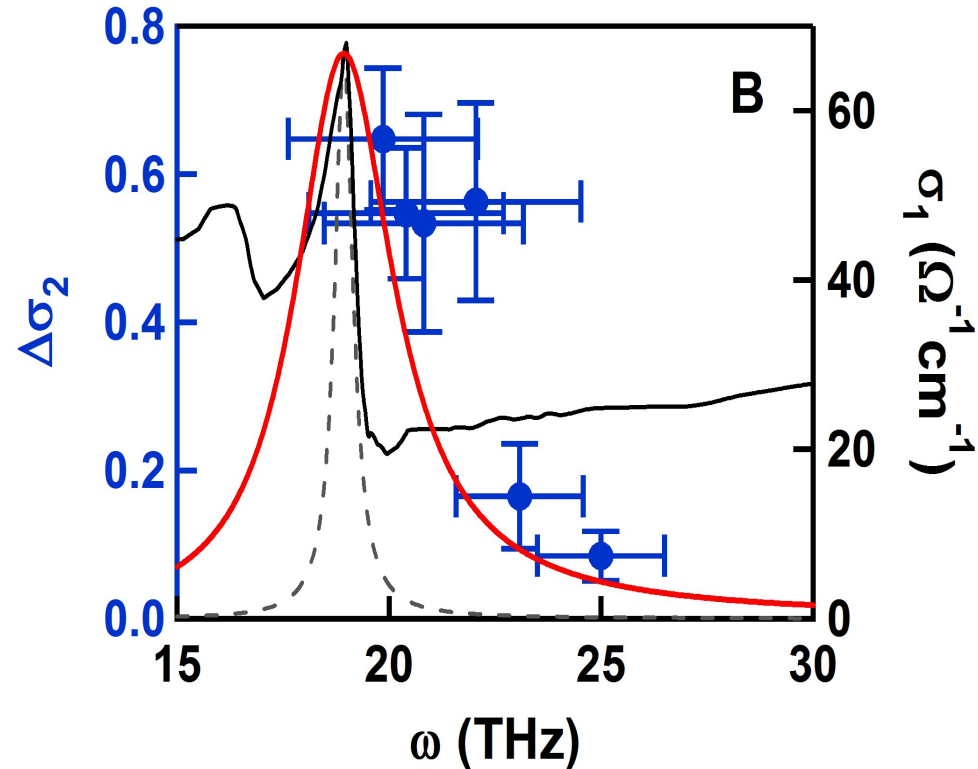
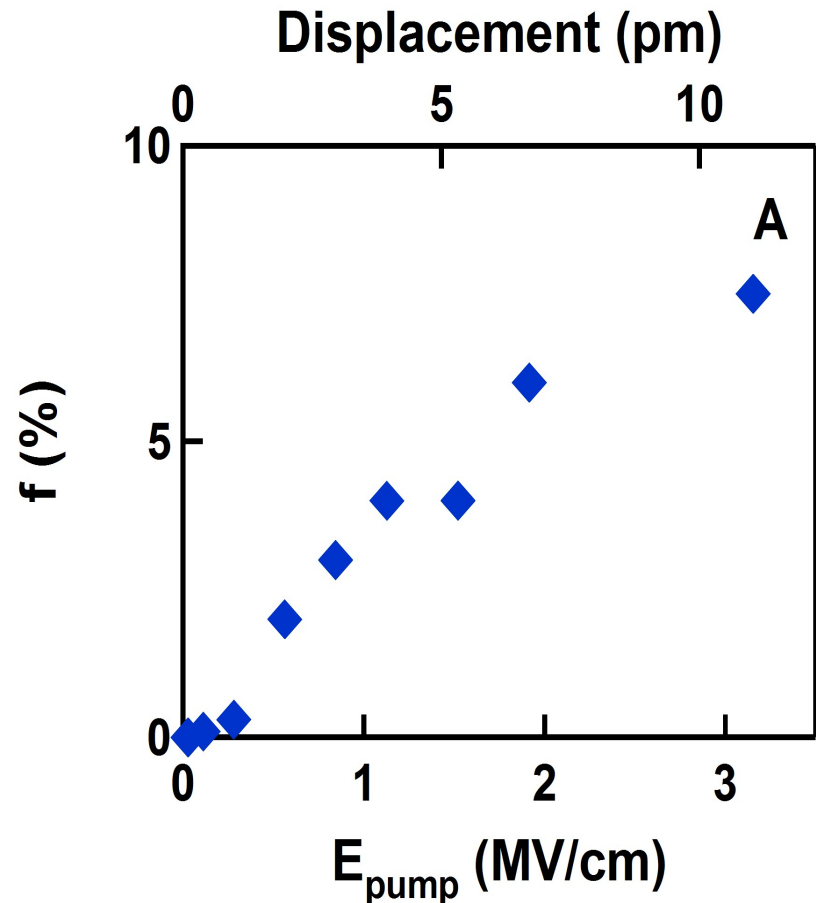
- reflectivity edge
- Loss function

Fluence and wavelength dependence



Resonant to apical oxygen phonon

Fluence and wavelength dependence

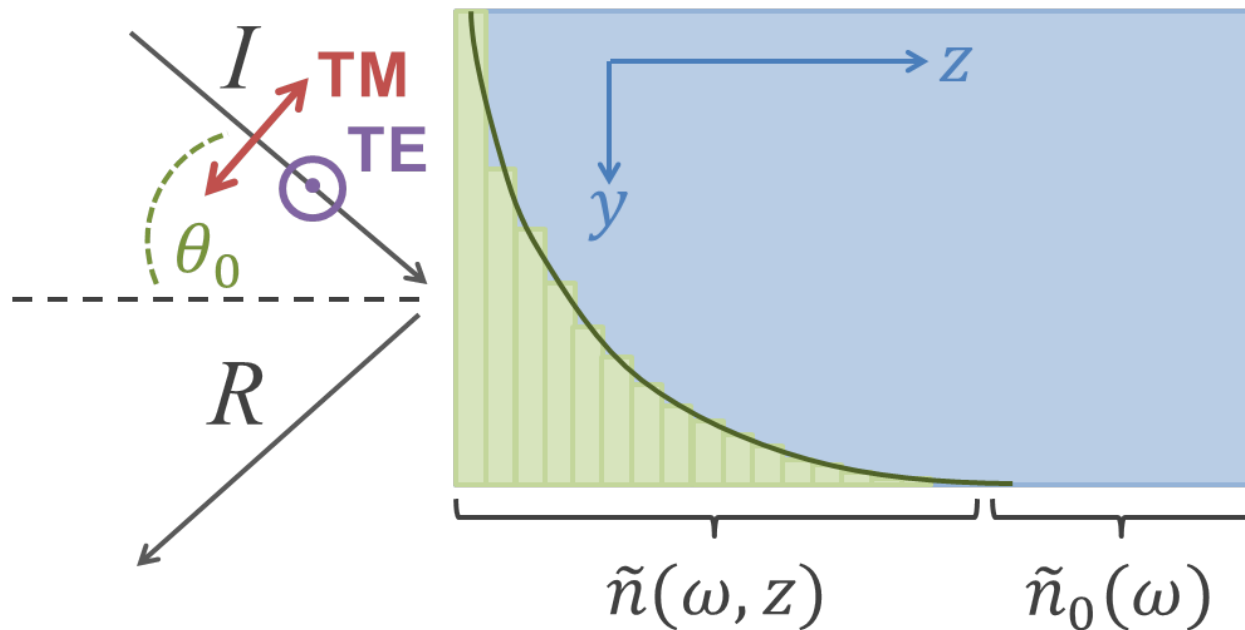


Linear field dependence
No percolation behavior found

Resonant to apical oxygen phonon

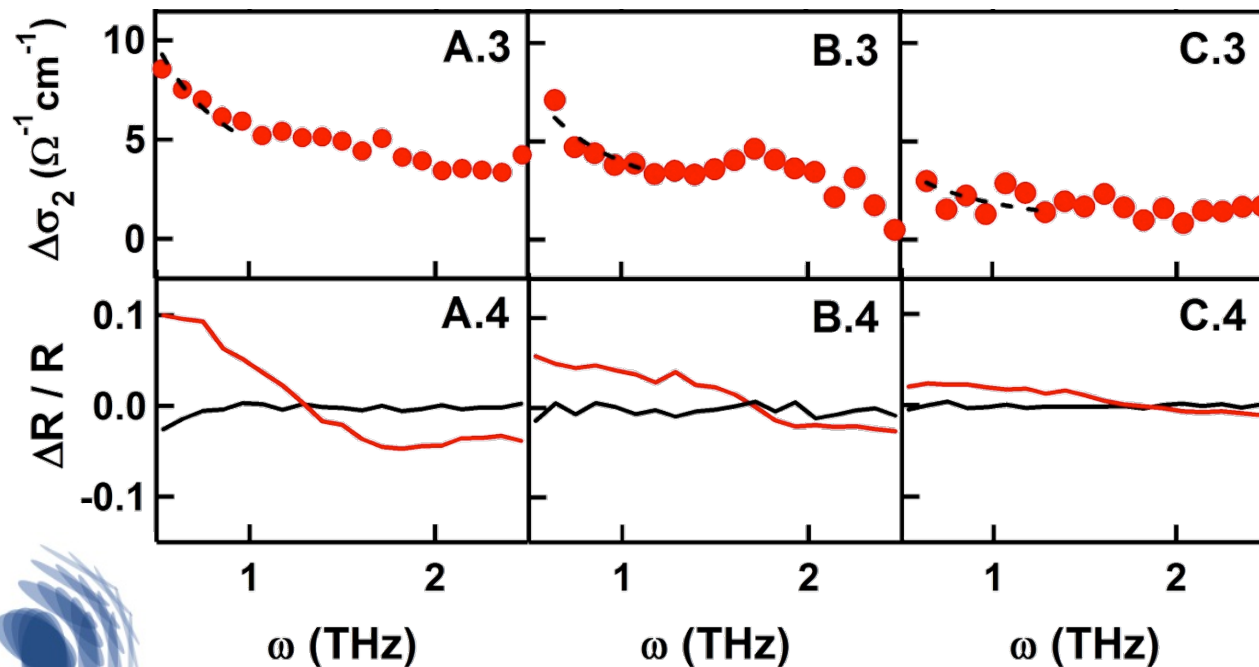
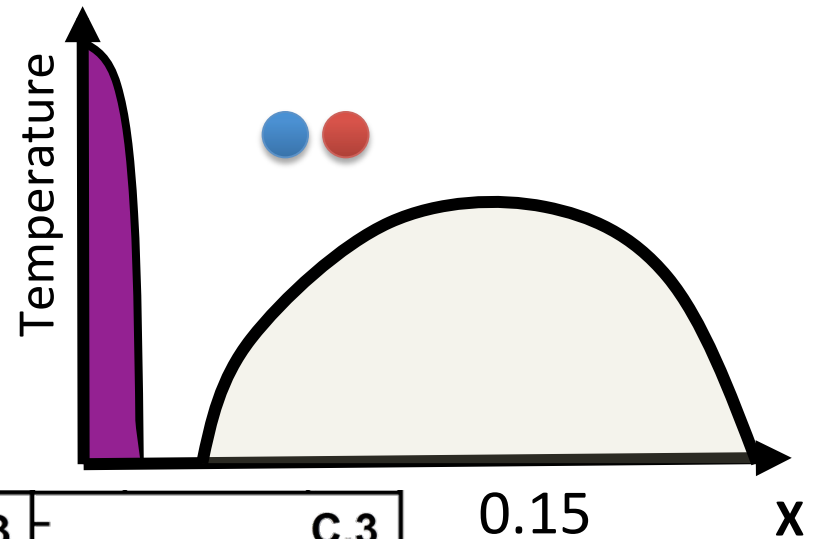
Penetration depth mismatch

$$\tilde{n}(\omega, z) = \tilde{n}_0(\omega) + \Delta\tilde{n}(\omega)e^{-\alpha z}$$

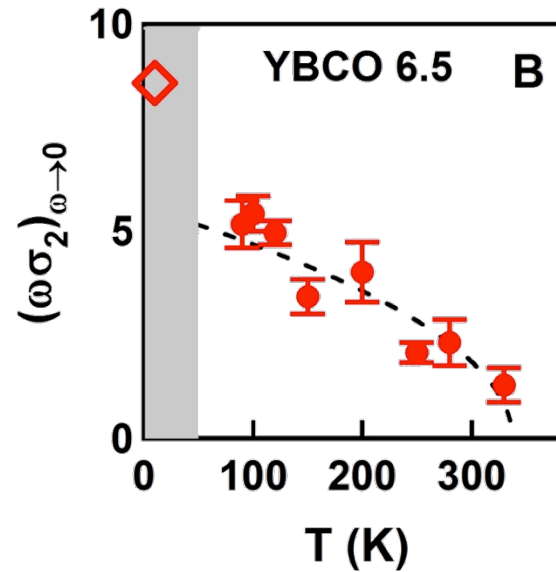
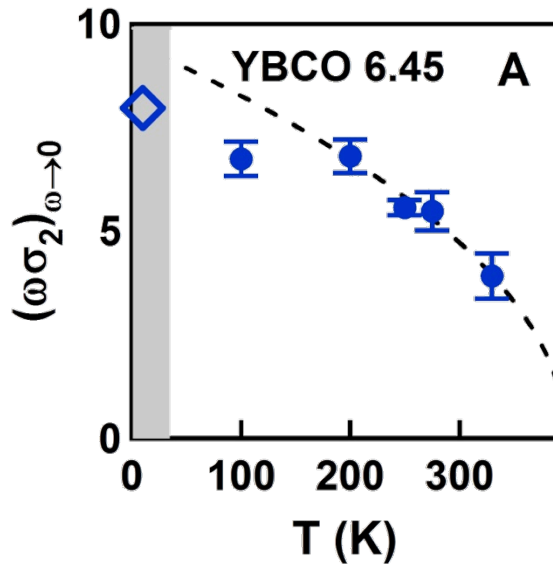


Multilayer model to calculate the optical properties

Temperature dependence YBCO 6.5 ($T_c=50$ K)

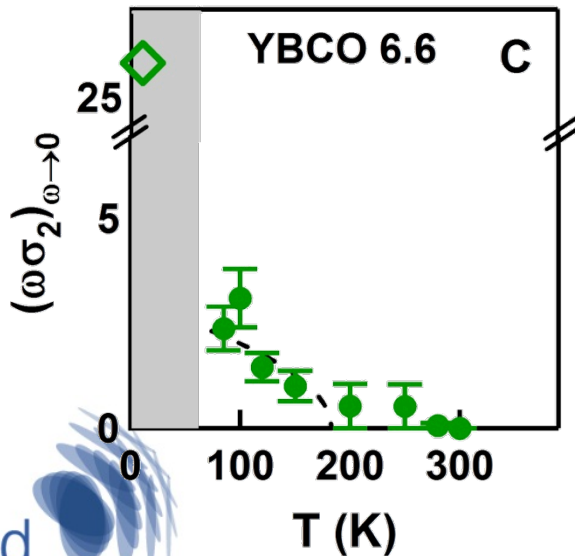


Temperature dependence of superfluid density

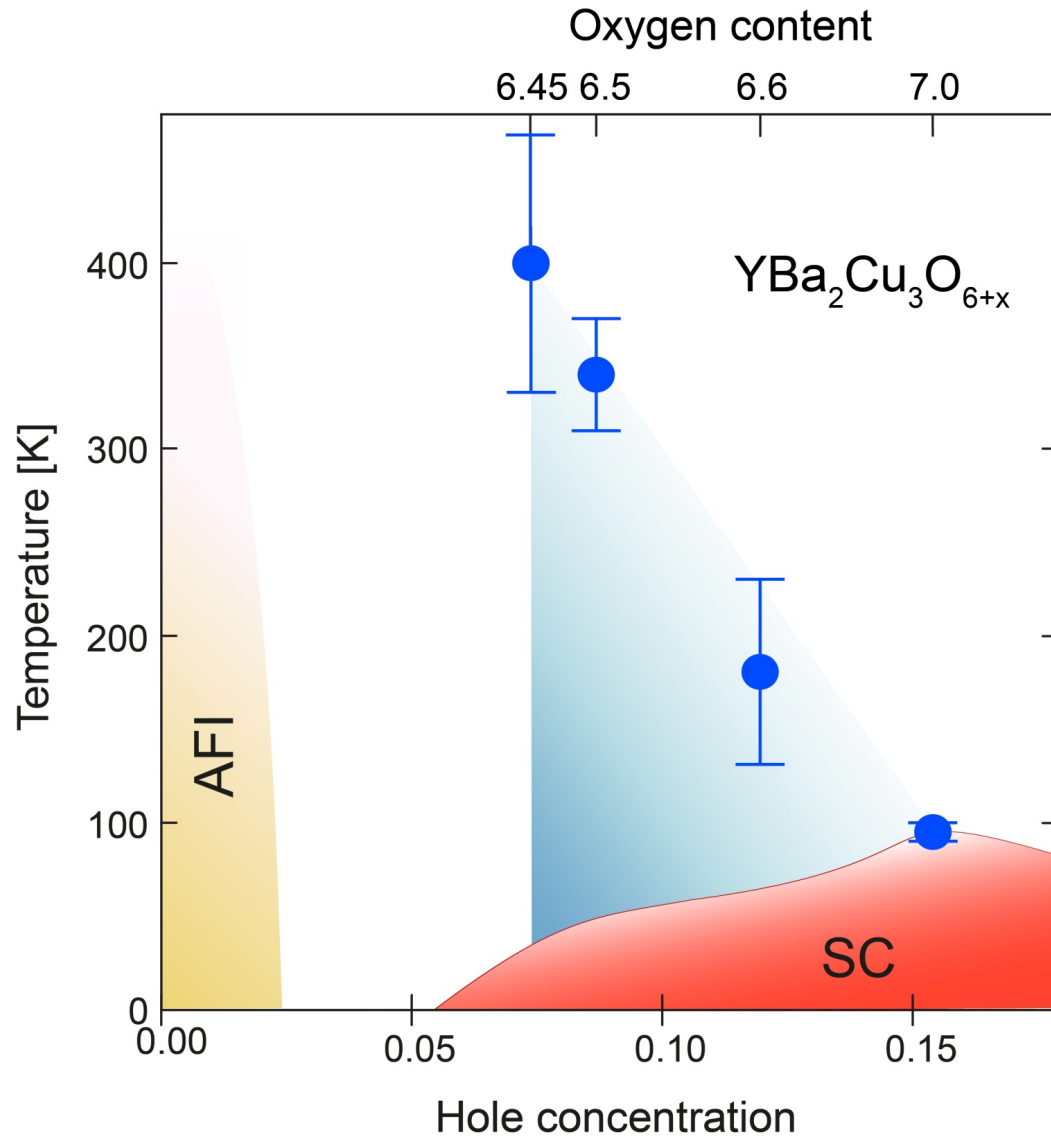


$$D^{SF} \propto (\omega\sigma_2)_{\omega \rightarrow 0}$$

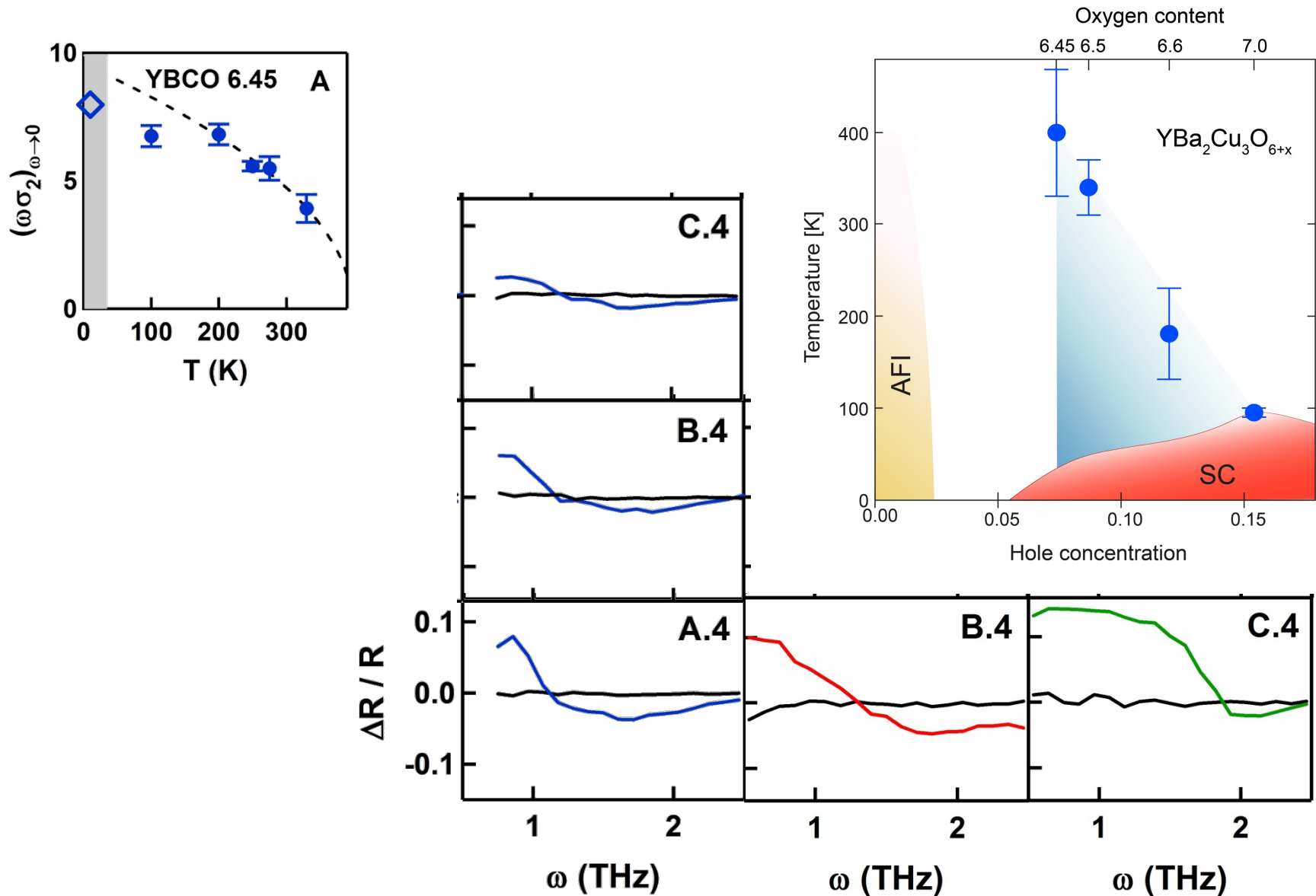
$$D^{SF} \propto \sqrt{1 - \frac{T}{T'}}$$



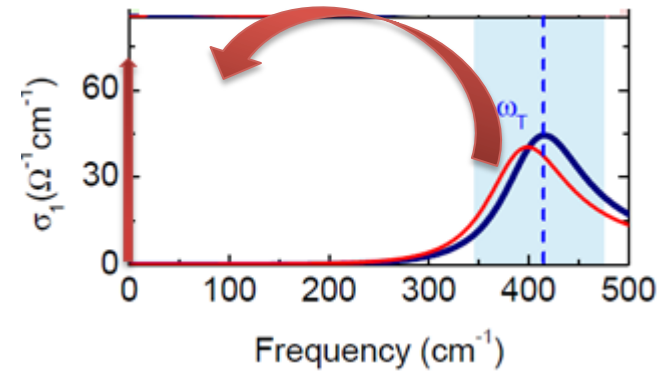
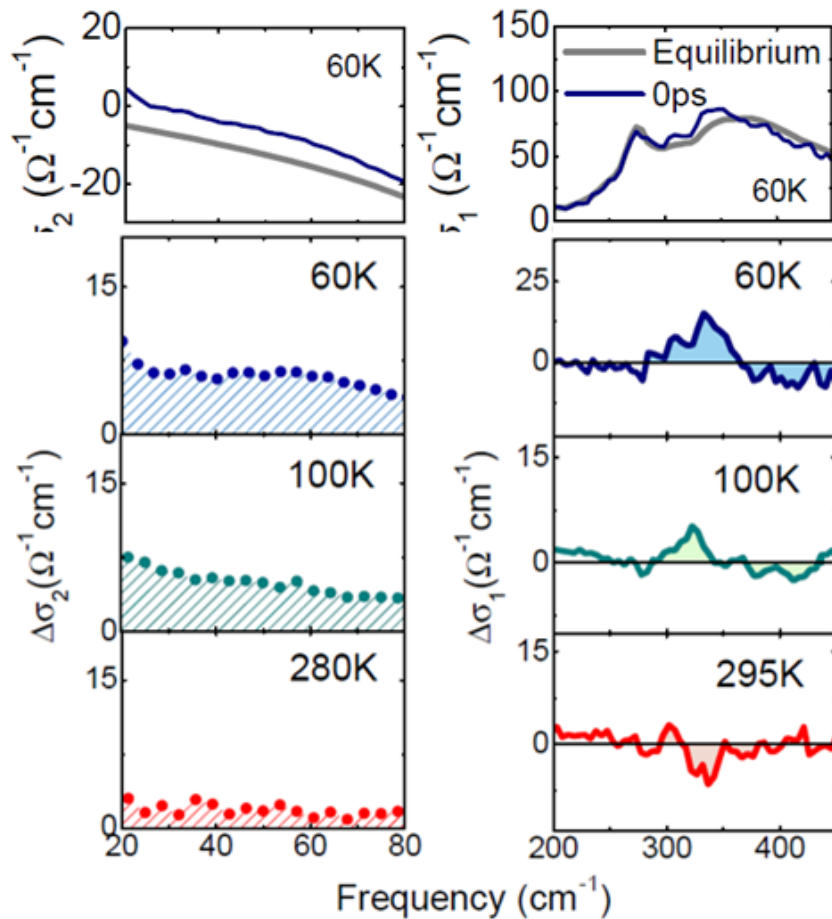
Induced Inductive Coupling far above T_c



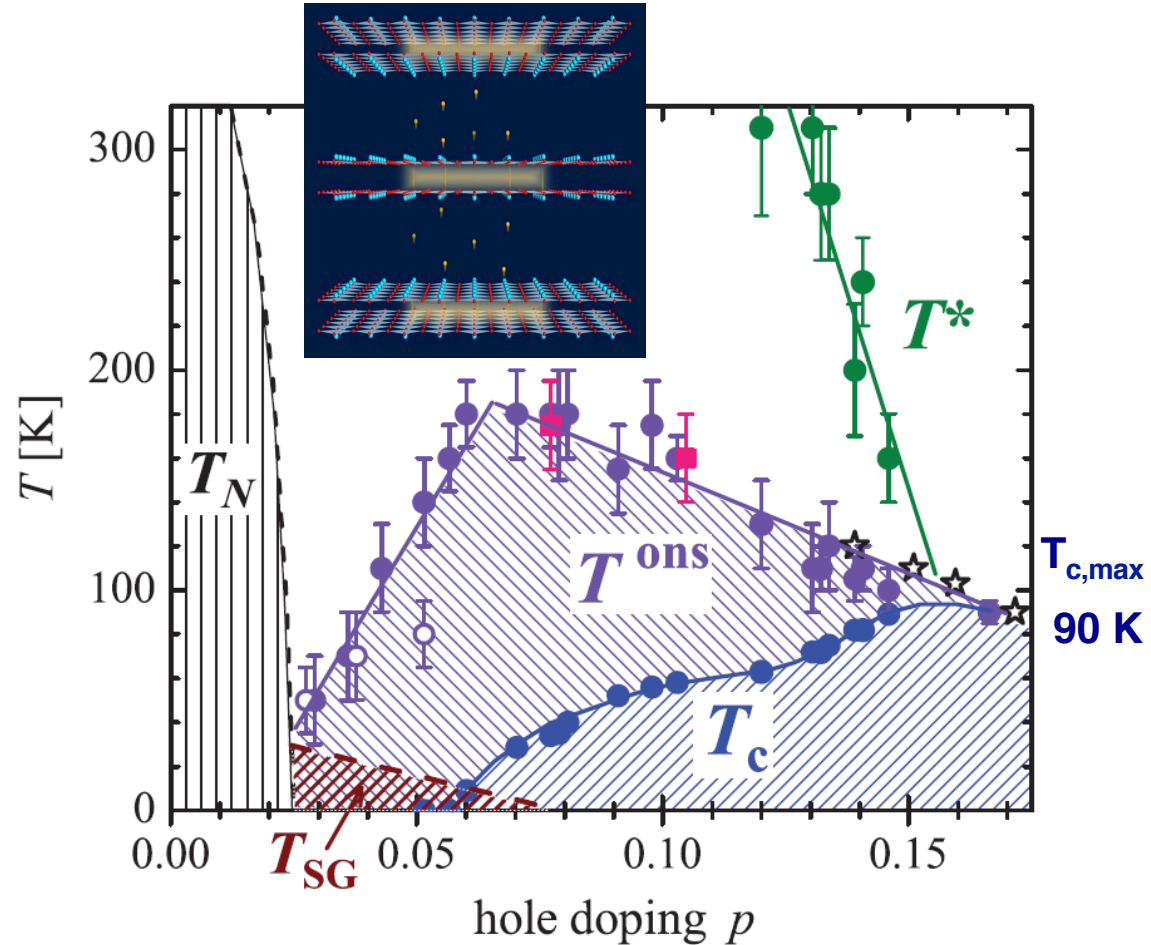
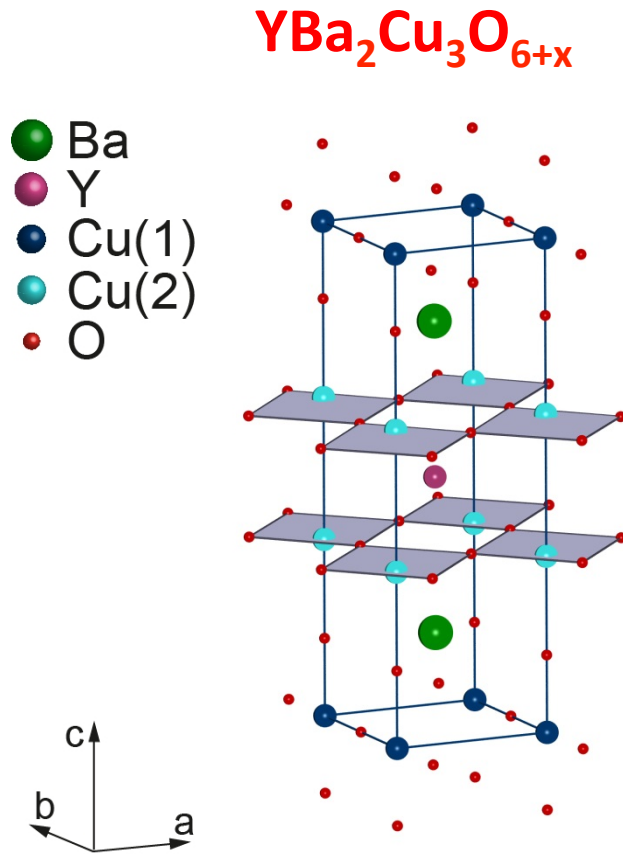
Inhomogeneous nature of the SC state



Temperature dependence

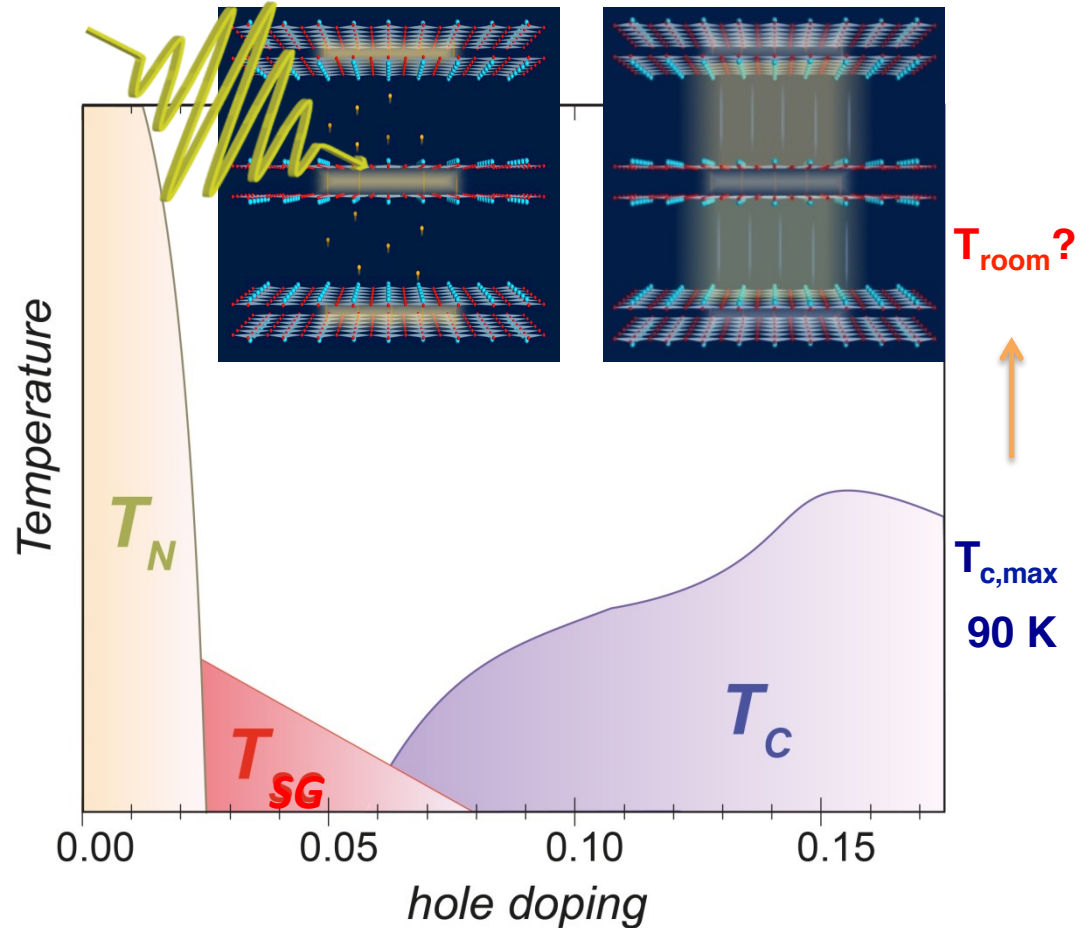
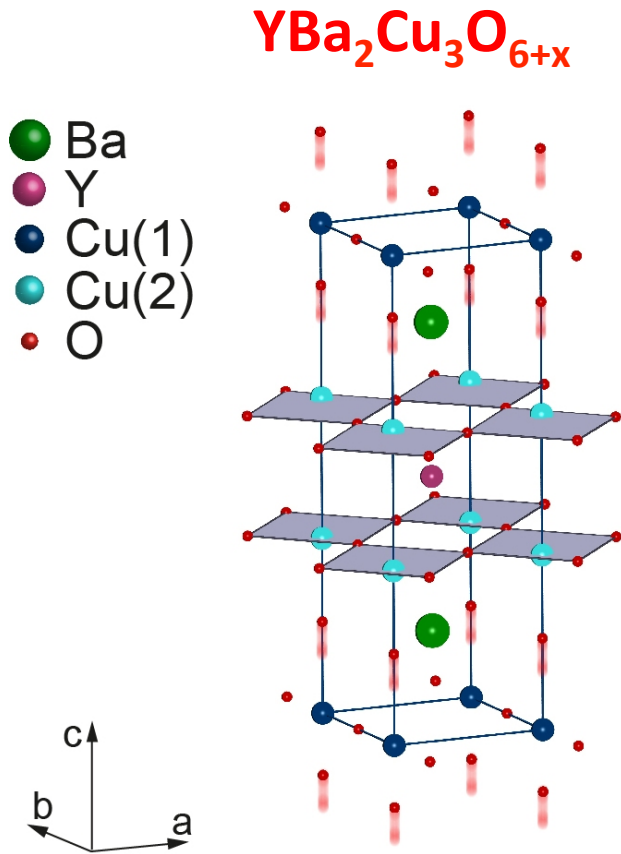


Precursor superconducting state above T_c



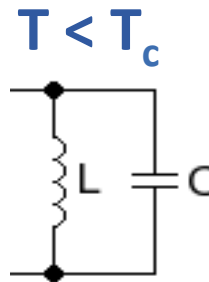
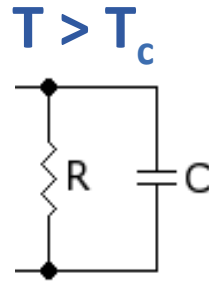
A. Dubroka *et al.*,
Phys. Rev. Lett. 107, 047006 (2011)




Coherent control of superconductivity



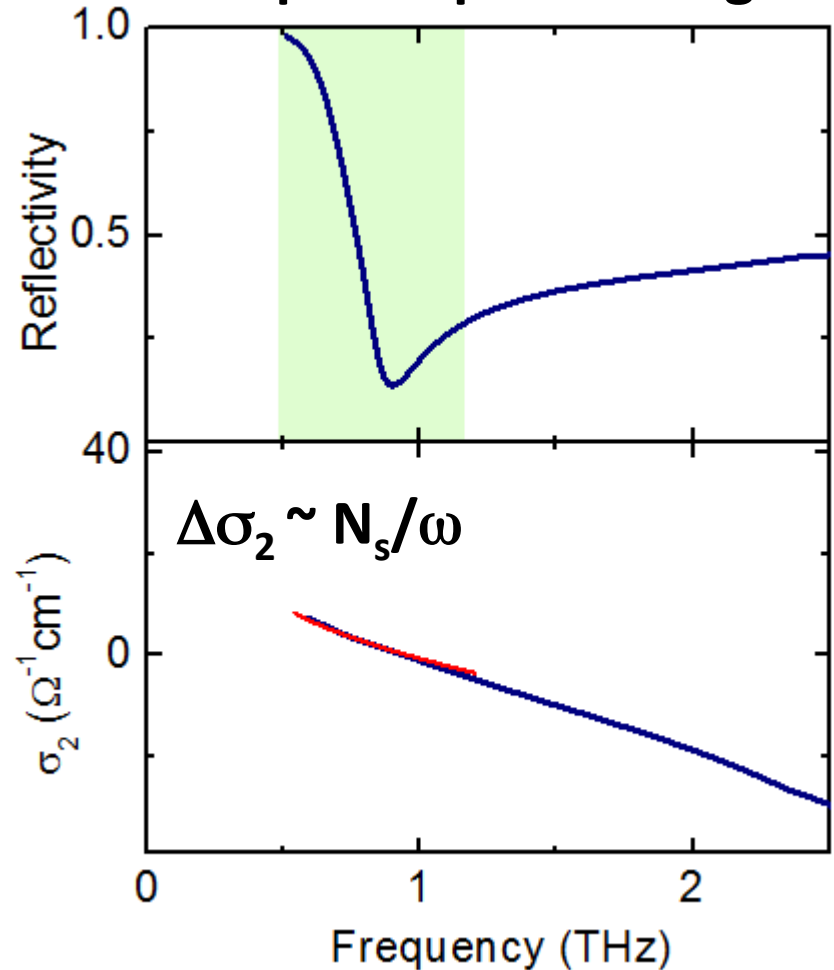
Signature of superconductivity

YBa₂Cu₃O_{6+x} c-axis

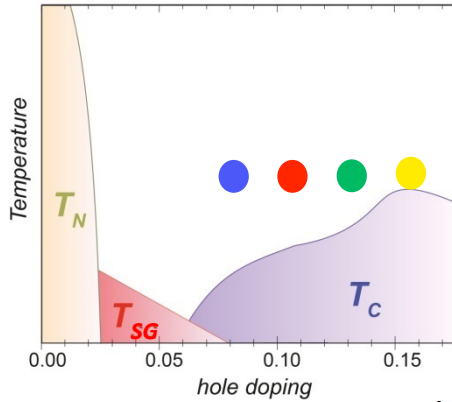


- Insulating layer 1 
- Superconducting layer 
- Insulating layer 2 

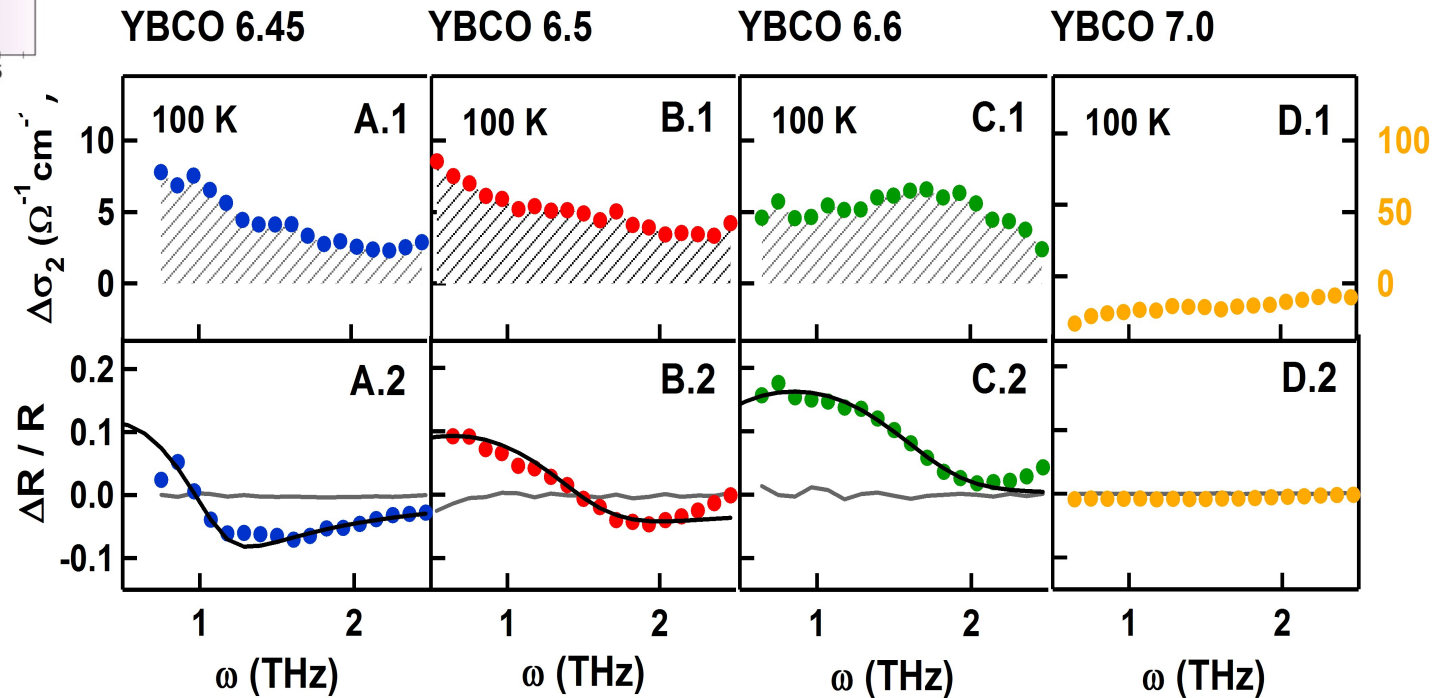
Josephson plasma edge



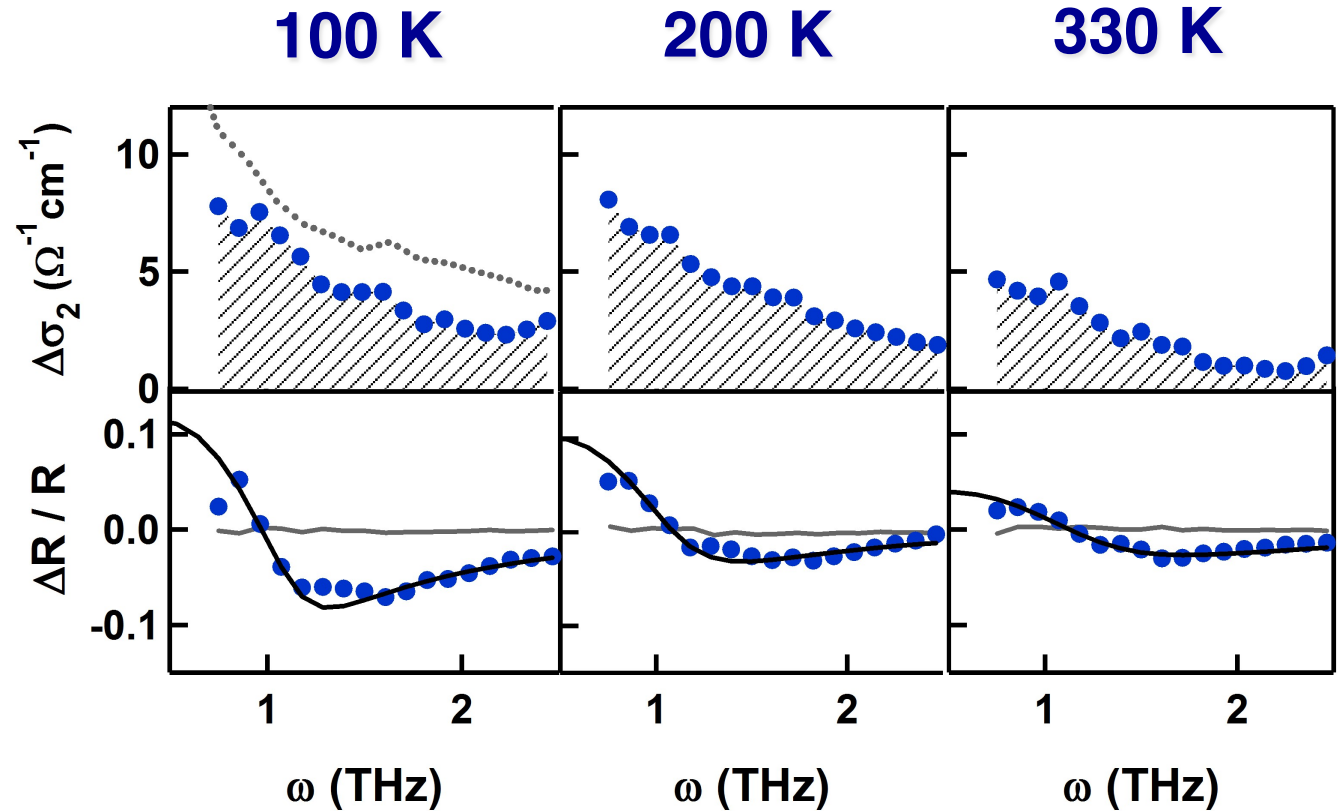
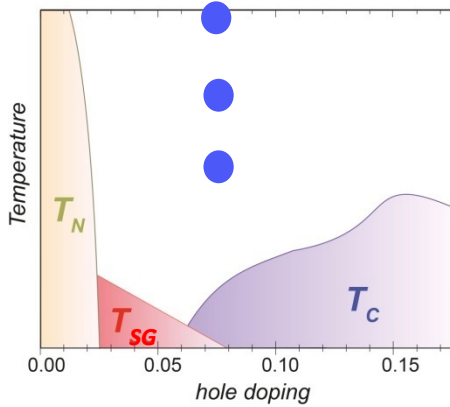
Transient Superconductivity above T_c



T = 100 K

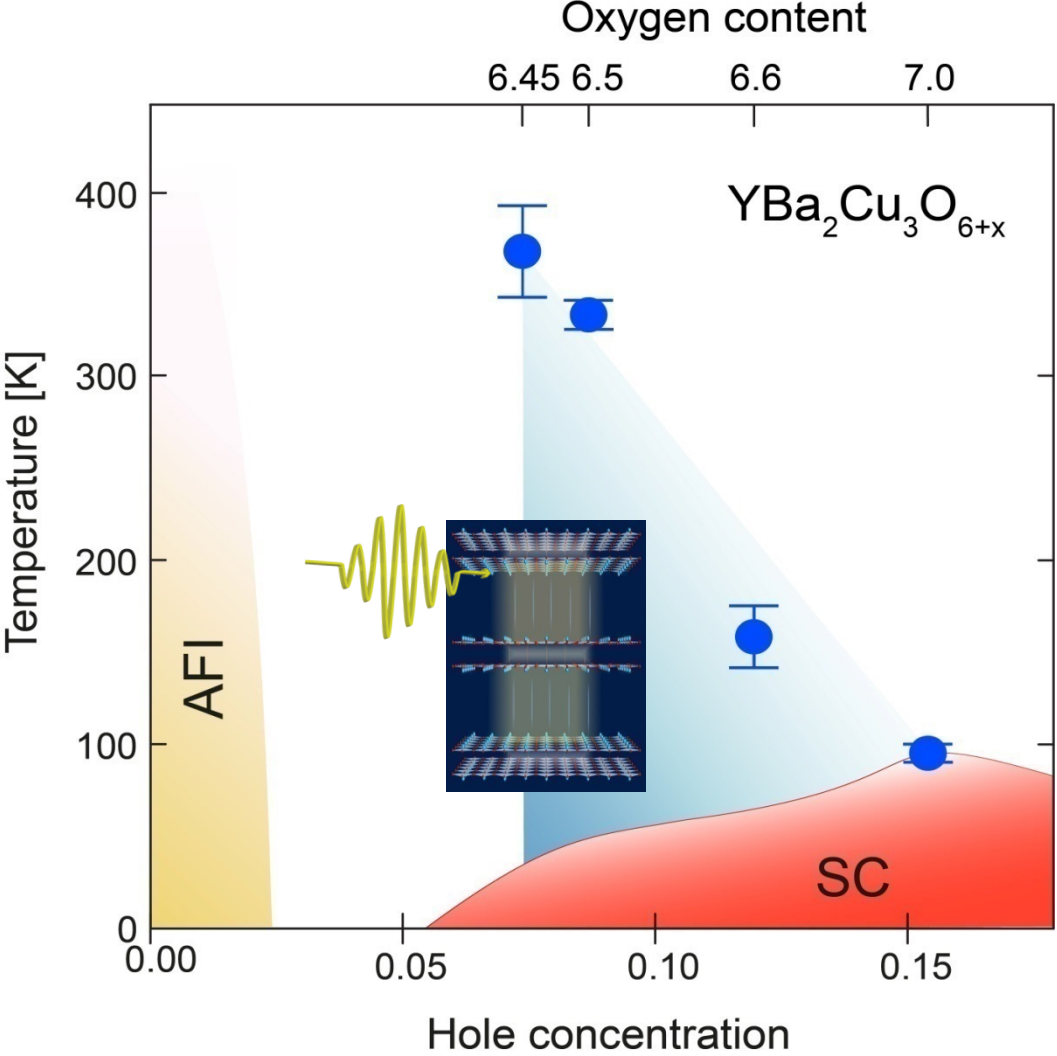


YBCO_{6.45}: Temperature Dependence

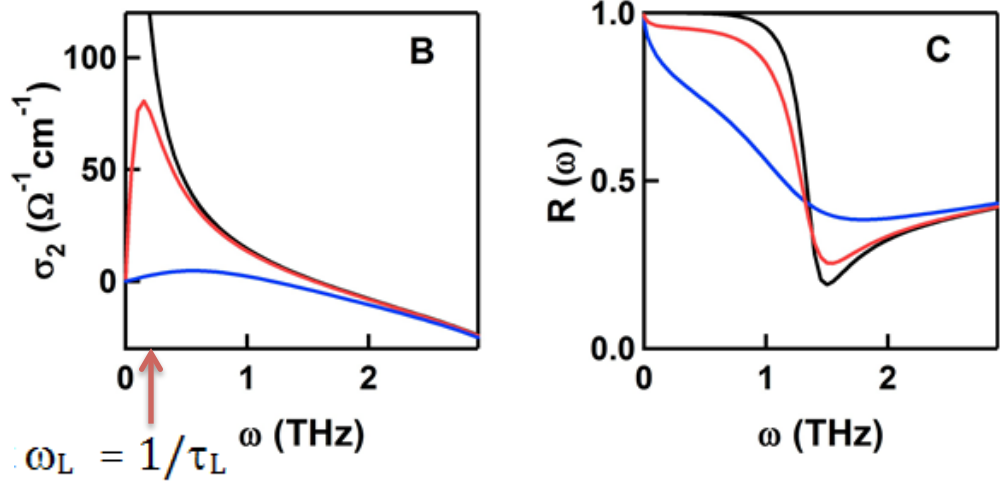


Light-induced Superconductivity up to Room Temperature

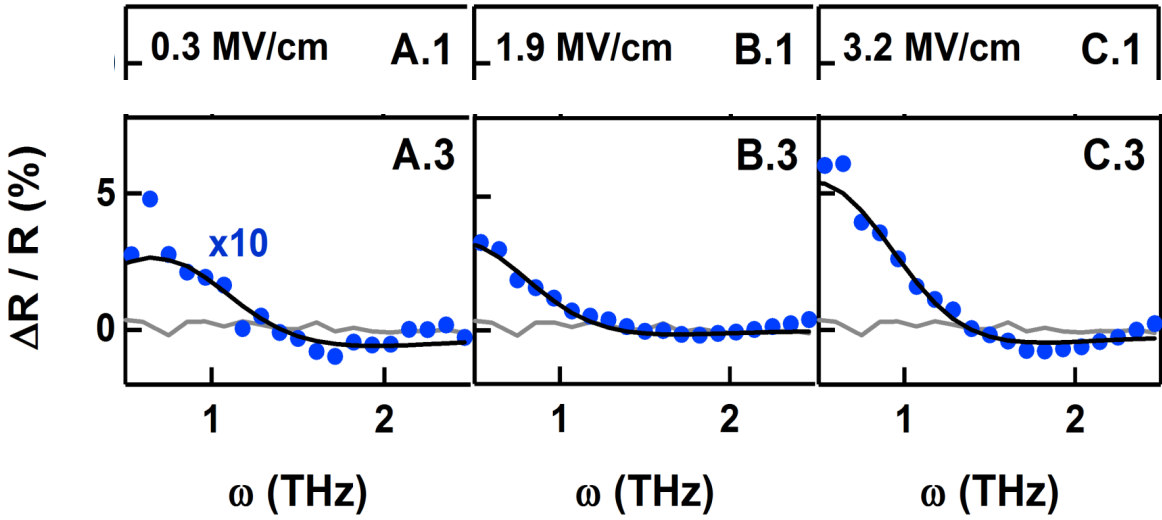
Phase Diagram



Other possible explanation



high-mobility
Drude
conductor ?
Photo-carriers
with a long τ ?



The position
of the edge
exhibited
no fluence
dependence